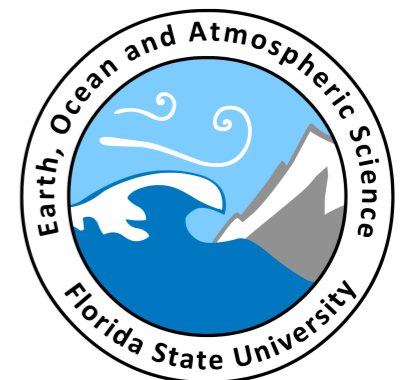
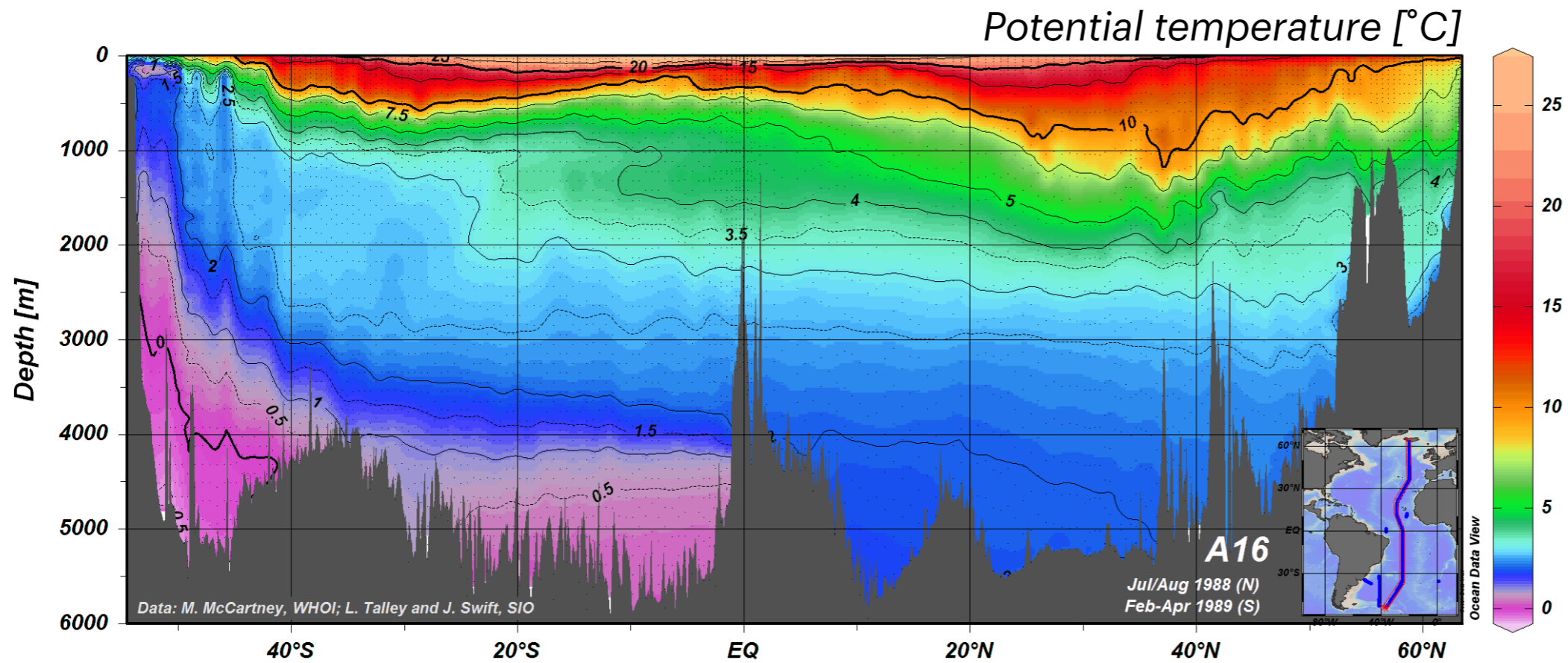


Towards a potential-vorticity based mesoscale closure scheme

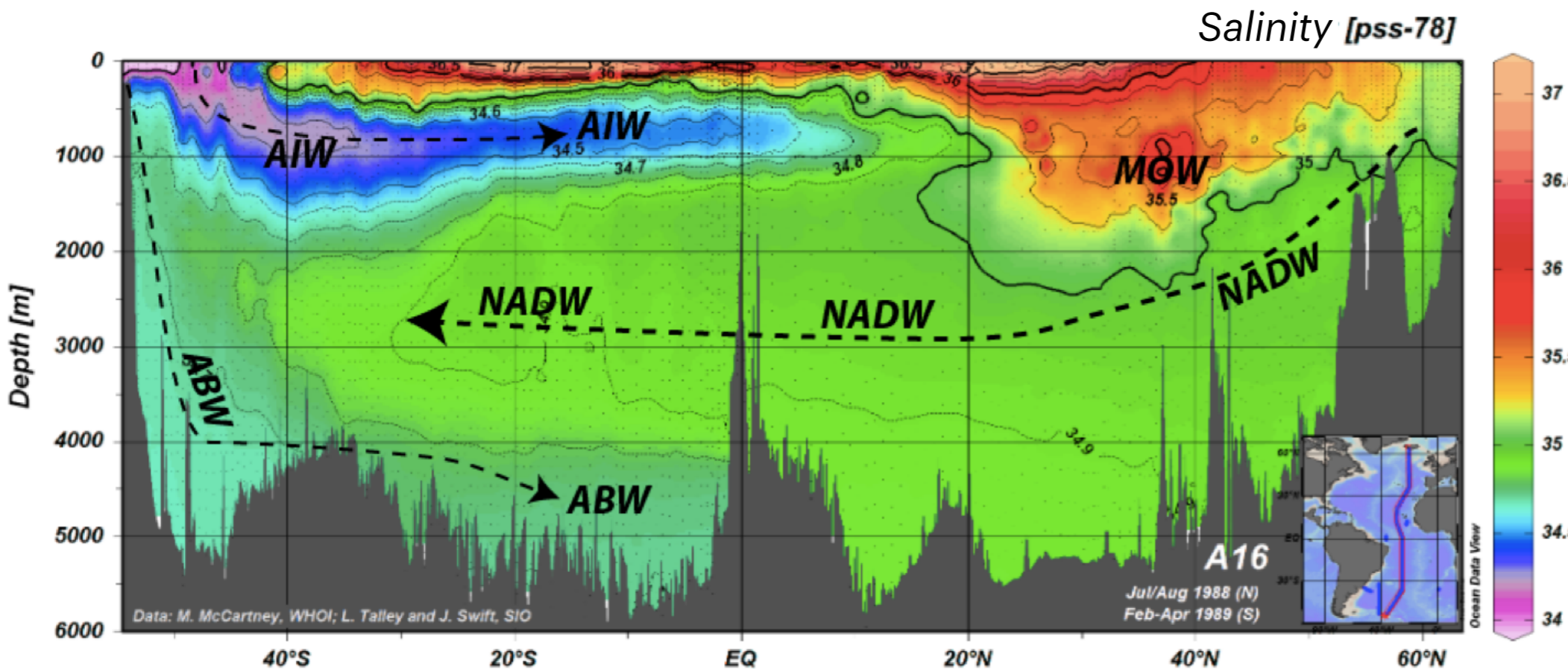
T. Uchida, Q. Jamet, W. Dewar, D. Balwada, J. Le Sommer & T. Penduff.
CESM Ocean Model Working Group Meeting, Feb. 3, 2021



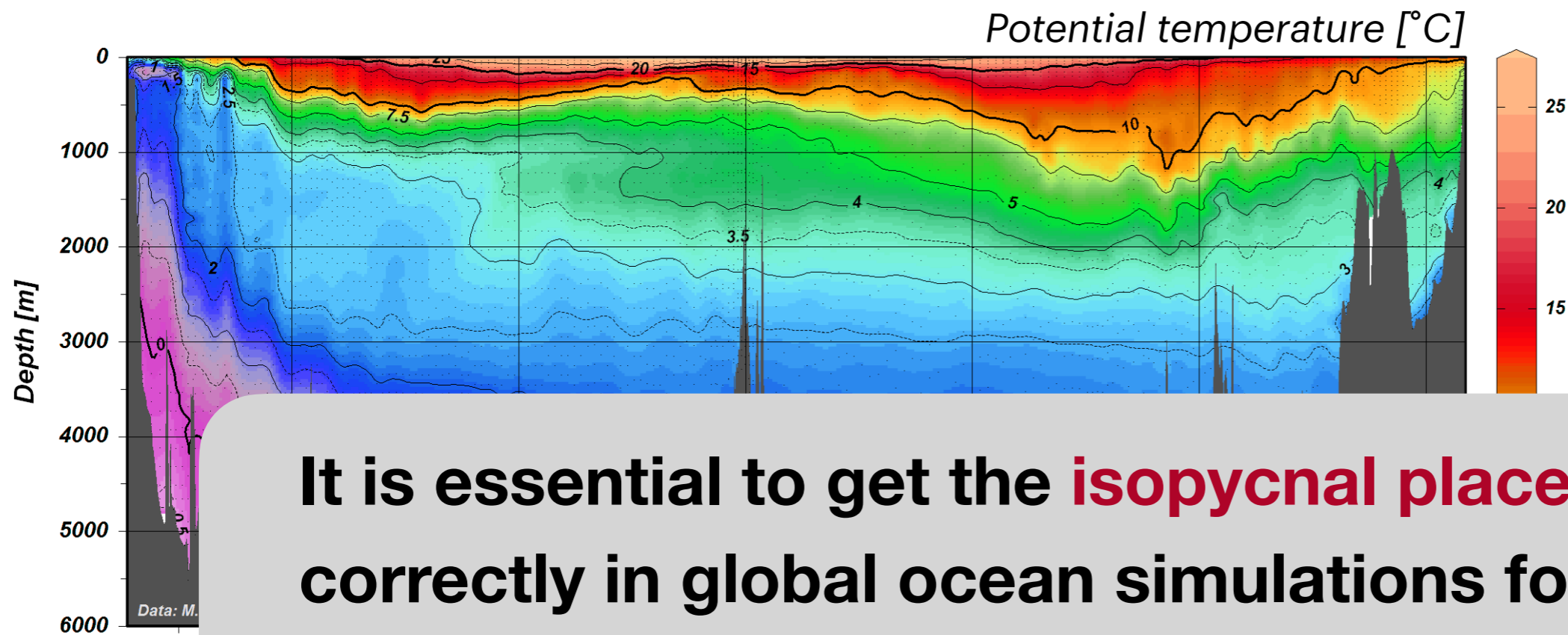
The stratified ocean



The ocean interior is stratified and quasi-adiabatic, so much so that we infer the global ocean circulation from tracer distribution along isopycnals.

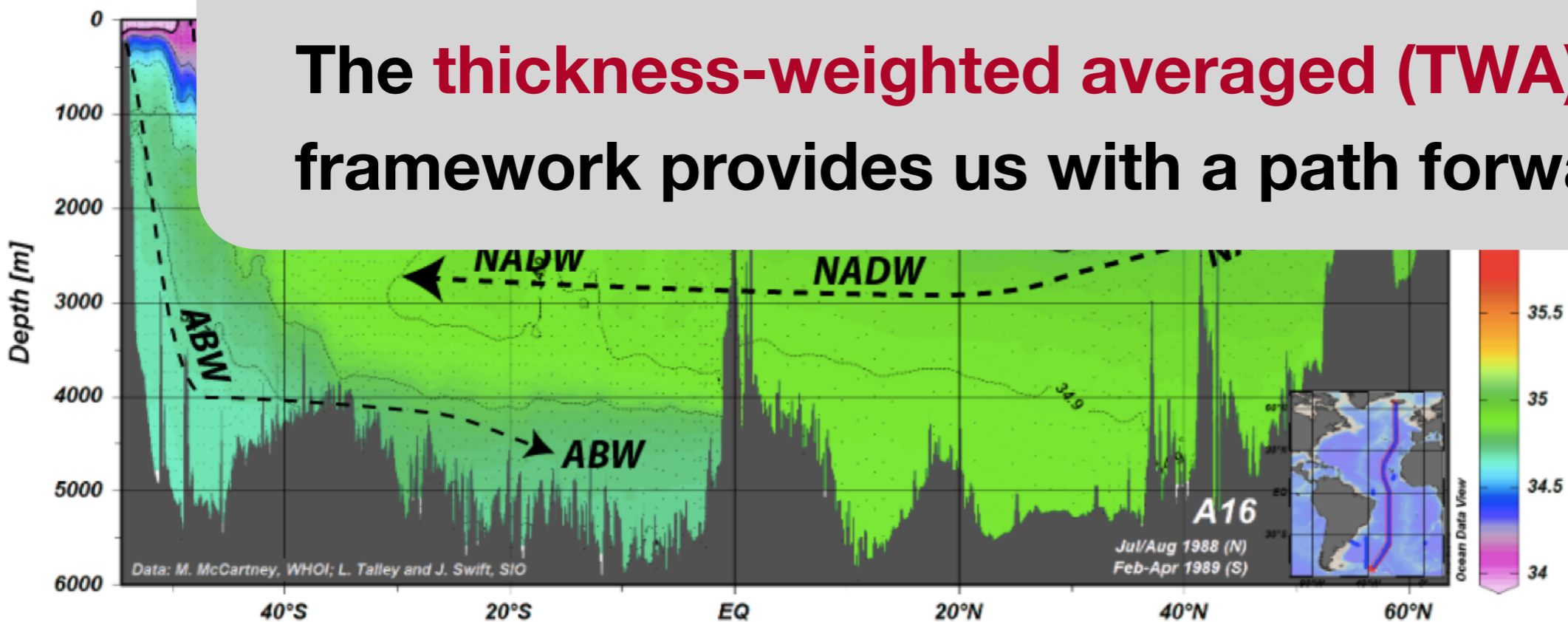


Figures taken from the World Ocean Circulation Experiment (WOCE)

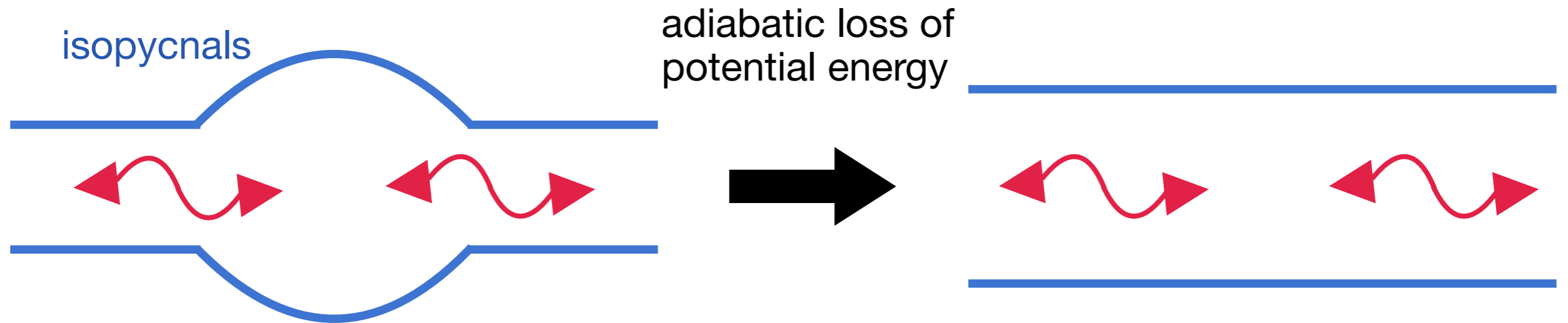


It is essential to get the **isopycnal placement** correctly in global ocean simulations for estimating global heat and tracer transport.

The **thickness-weighted averaged (TWA)** framework provides us with a path forward.



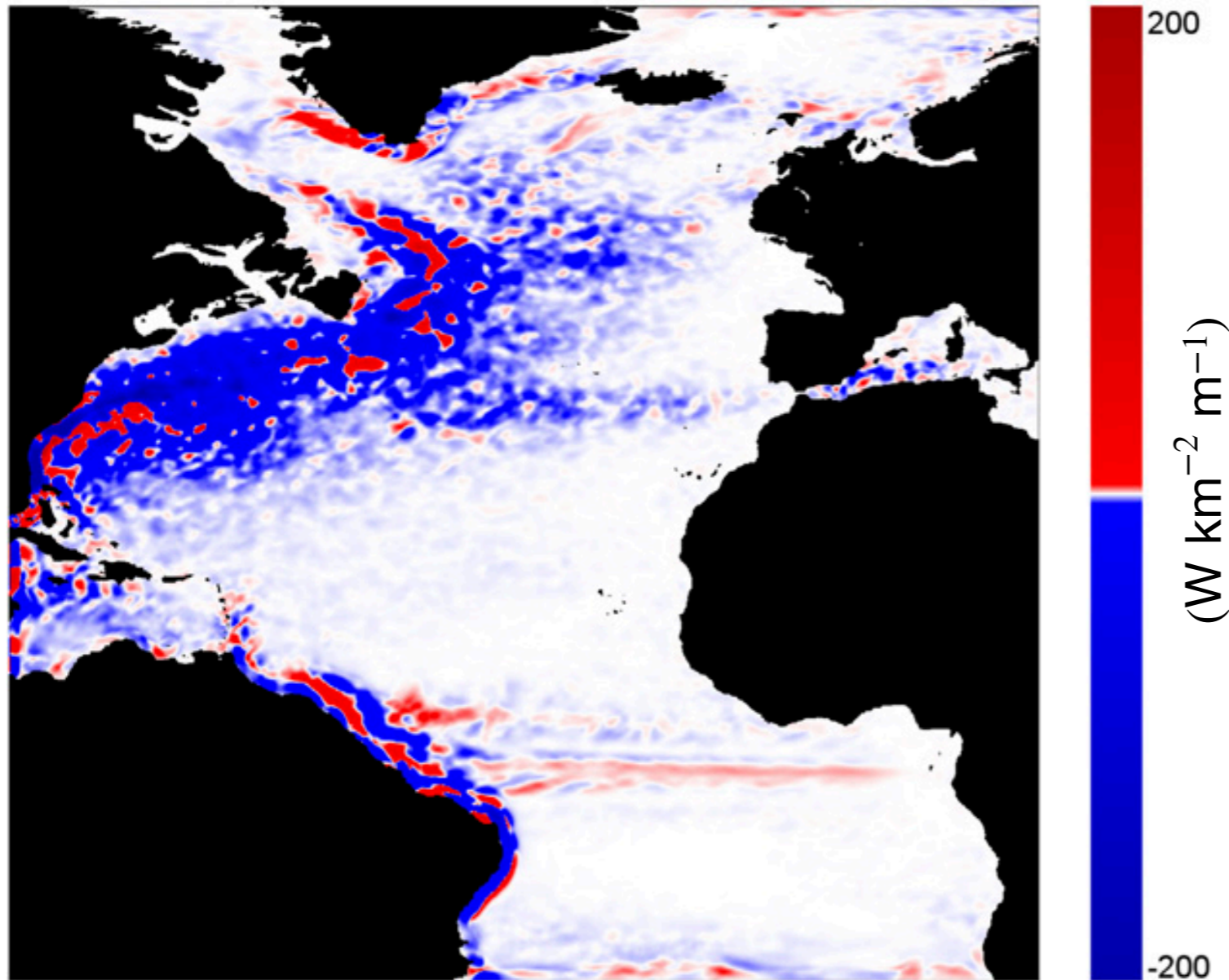
Figures taken from the World Ocean Circulation Experiment (WOCE)



- The **Gent-McWilliams (GM) skew diffusivity** diffuses the isopycnal thickness in a similar manner to how baroclinic instability would if resolved.
- The **REDI diffusivity** represents the enhanced tracer stirring along isopycnals due to eddies.

- **GM and REDI should be related to one another.**
- **Can we capture the full eddy feedback and not just the release of available potential energy?**

The eddy momentum feedback



Aluie et al. (2018)

- Employing a coarse-graining method, Aluie et al. (2018) examined the direction of kinetic energy cascade from a 0.1° model simulation.
- **Blue** shadings indicate the eddies fluxing kinetic energy back into the mean flow.

**Mesoscale eddies energize the Gulf Stream.
Can we say more on **how** and **where**?**

$$\hat{u}_{\tilde{t}} + \hat{u}\hat{u}_{\tilde{x}} + \hat{v}\hat{u}_{\tilde{y}} + \hat{\omega}\hat{u}_{\tilde{b}} - f\hat{v} + \overline{m}_{\tilde{x}} = -\bar{\mathbf{e}}_1 \cdot (\tilde{\nabla} \cdot \mathbf{E}) + \hat{\mathcal{X}}$$

$$\hat{v}_{\tilde{t}} + \hat{u}\hat{v}_{\tilde{x}} + \hat{v}\hat{v}_{\tilde{y}} + \hat{\omega}\hat{v}_{\tilde{b}} + f\hat{u} + \overline{m}_{\tilde{y}} = -\bar{\mathbf{e}}_2 \cdot (\tilde{\nabla} \cdot \mathbf{E}) + \hat{\mathcal{Y}}$$

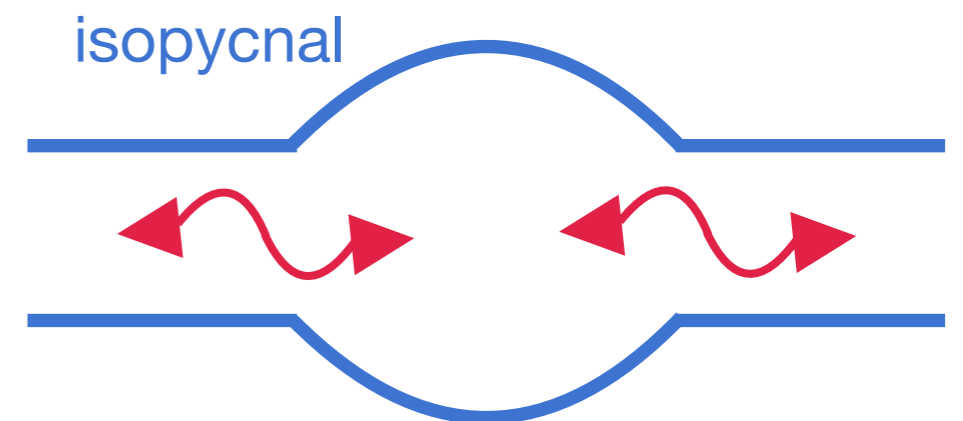
- $\hat{\mathbf{u}}$ ($= \frac{\overline{\sigma\mathbf{u}}}{\bar{\sigma}}$): the thickness-weighted averaged (TWA) velocity.

- $\overline{(\cdot)}$ is the ensemble mean.

- σ ($= \zeta_{\tilde{b}}$): the isopycnal thickness.

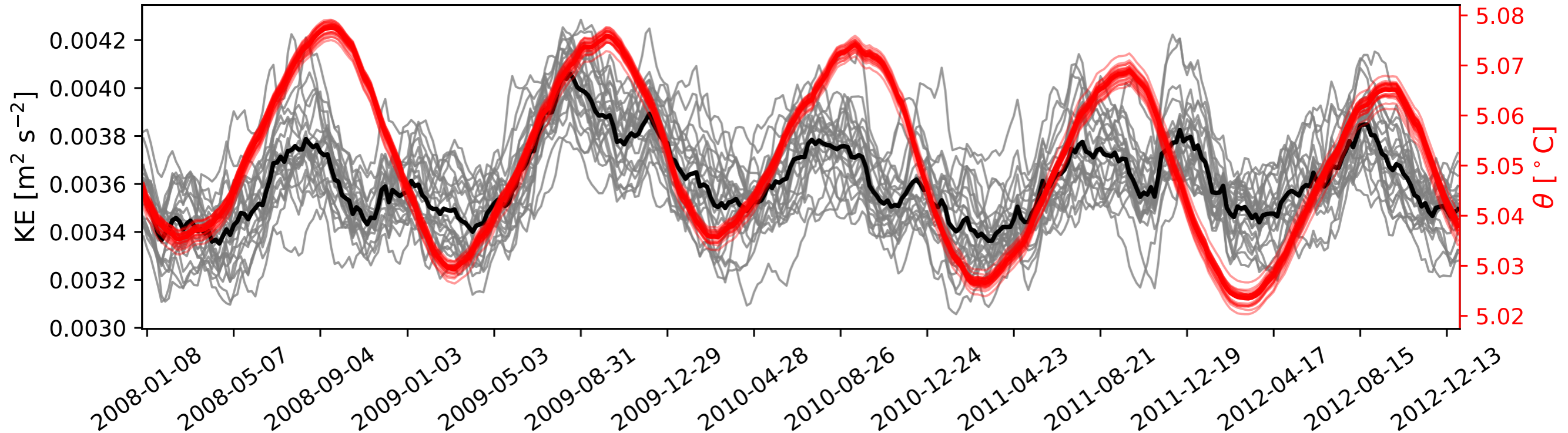
- ω : the diapycnal velocity.

- m ($= \phi - b\zeta$): the Montgomery potential.



A 24-member ensemble simulation

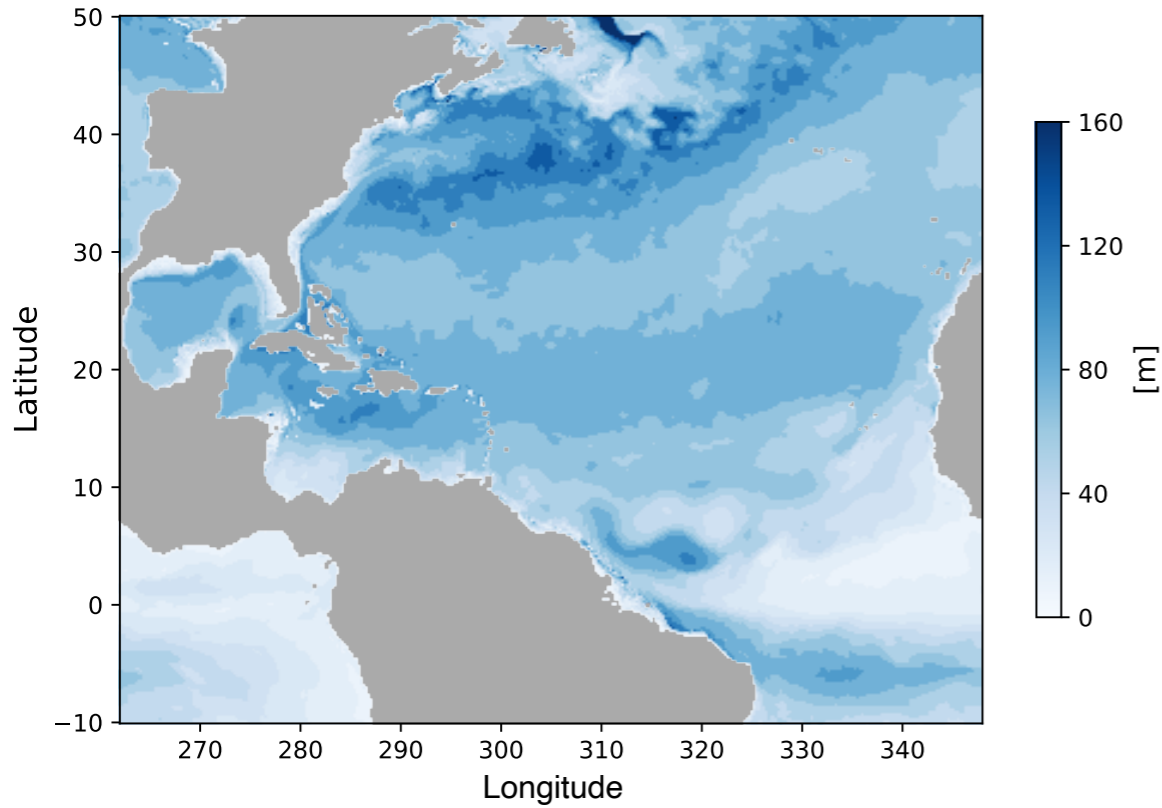
Domain-averaged kinetic energy and potential temperature



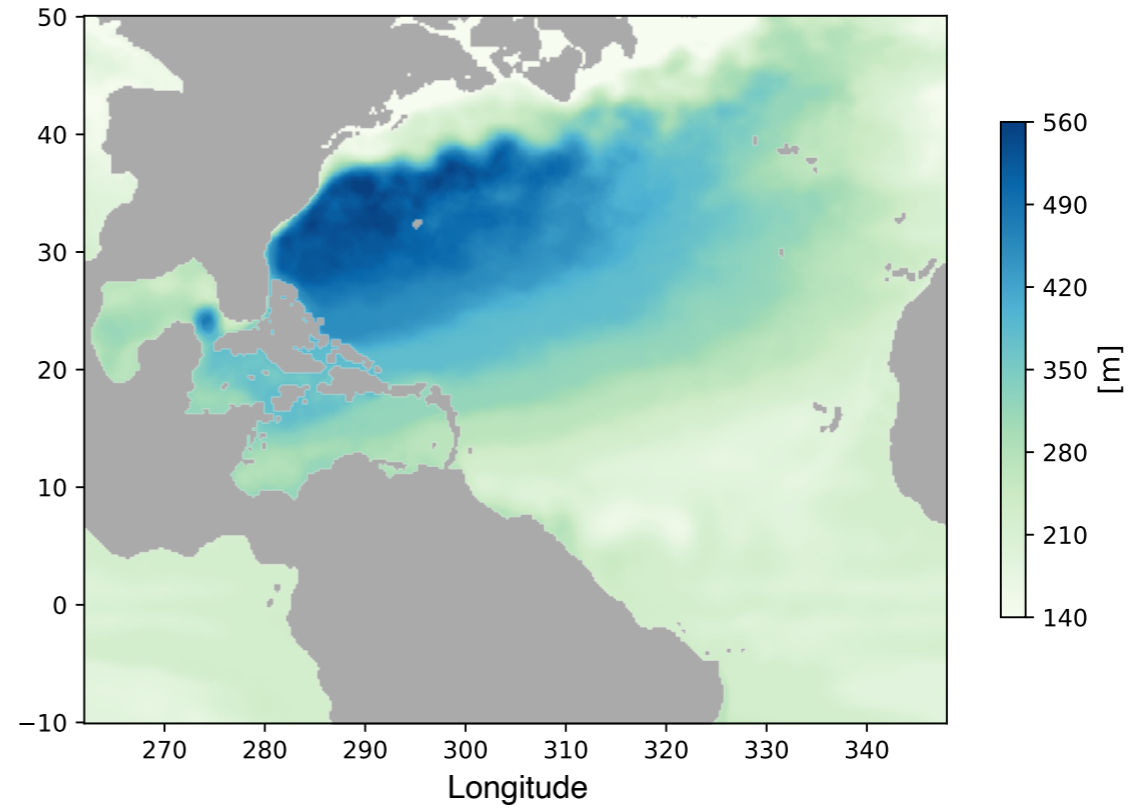
Thin lines: each ensemble member
Thick lines: ensemble mean

- **No. of ensemble members: 24.**
- **Resolution: 1/12°; Duration: 50 years (1963-2012).**
- **Model: MITgcm; Basin: North Atlantic.**
- **Surface boundary condition: partially air-sea coupled.**
- **Lateral boundary condition: relaxation and radiation conditions.**

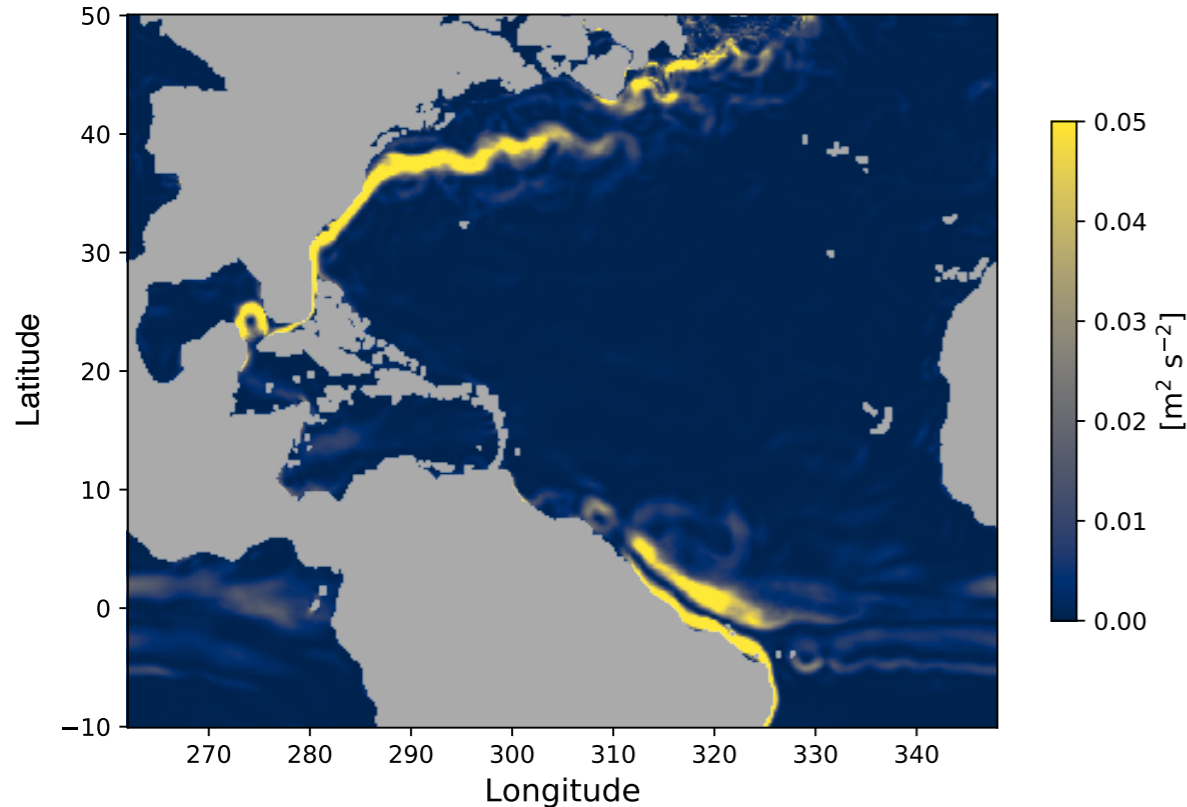
$\overline{\text{MLD}}$



$\overline{\zeta(\tilde{b})}$



TWA kinetic energy $((\hat{u}^2 + \hat{v}^2)/2)$



- Focus on an isopycnal whose ensemble-mean depth ($\overline{\zeta}$) is below the ensemble-mean mixed-layer depth ($\overline{\text{MLD}}$).
- The isopycnal shoals drastically across the separated Gulf Stream.

$$\hat{u}_{\tilde{t}} + \hat{u}\hat{u}_{\tilde{x}} + \hat{v}\hat{u}_{\tilde{y}} + \hat{\omega}\hat{u}_{\tilde{b}} - f\hat{v} + \overline{m}_{\tilde{x}} = -\bar{\mathbf{e}}_1 \cdot (\tilde{\nabla} \cdot \mathbf{E}) + \hat{\mathcal{X}}$$

$$\hat{v}_{\tilde{t}} + \hat{u}\hat{v}_{\tilde{x}} + \hat{v}\hat{v}_{\tilde{y}} + \hat{\omega}\hat{v}_{\tilde{b}} + f\hat{u} + \overline{m}_{\tilde{y}} = -\bar{\mathbf{e}}_2 \cdot (\tilde{\nabla} \cdot \mathbf{E}) + \hat{\mathcal{Y}}$$

- The net eddy feedback onto the (TWA) mean fields are encapsulated in the Eliassen-Palm flux (\mathbf{E}) divergence.
- For an eddy closure, it shifts the focus from the buoyancy equation (GM) to the momentum equations.

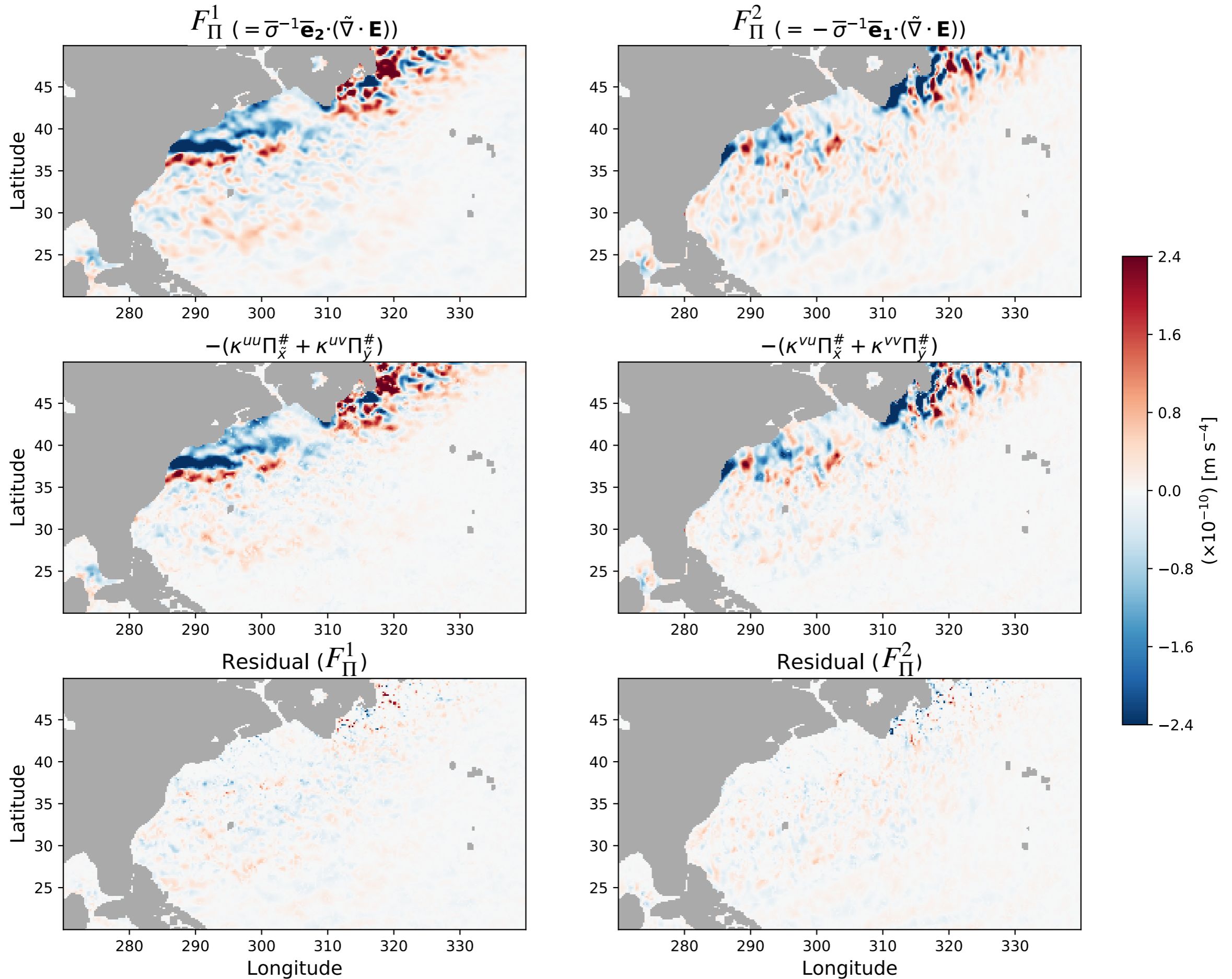
If we can parametrize $\bar{\mathbf{e}} \cdot (\nabla \cdot \mathbf{E})$, we have a physically consistent eddy closure scheme which represents the eddy buoyancy & momentum fluxes.

Can we parametrize the Eliassen-Palm flux divergence?

- The Eliassen-Palm flux divergence is directly related to the eddy Ertel potential vorticity (PV) flux (\mathbf{F}^Π).
- This primes us to connect the Eliassen-Palm flux divergence to the large-scale Ertel PV.
- We relate the eddy Ertel PV flux to the local-gradient flux of the mean Ertel PV ($\Pi^\#$) via an anisotropic eddy diffusivity tensor (\mathbf{K}).

$$\begin{array}{c}
 \text{Redi} \\
 \text{GM} \\
 + \\
 \text{Eddy momentum flux}
 \end{array}
 \underbrace{\left(\begin{array}{cc}
 \widehat{u''\theta''} & \widehat{v''\theta''} \\
 \widehat{u''s''} & \widehat{v''s''} \\
 F^{\Pi 1} & F^{\Pi 2}
 \end{array} \right)}_{\mathbf{F}}
 = -
 \underbrace{\left(\begin{array}{cc}
 \hat{\theta}_{\tilde{x}} & \hat{\theta}_{\tilde{y}} \\
 \hat{s}_{\tilde{x}} & \hat{s}_{\tilde{y}} \\
 \Pi^\#_{\tilde{x}} & \Pi^\#_{\tilde{y}}
 \end{array} \right)}_{\mathbf{G}}
 \cdot
 \underbrace{\left(\begin{array}{cc}
 \kappa^{uu} & \kappa^{vu} \\
 \kappa^{uv} & \kappa^{vv}
 \end{array} \right)}_{\mathbf{K}}$$

Can we parametrize the Eliassen-Palm flux divergence?



- The Eliassen-Palm flux divergence, which is directly related to the eddy PV flux, encapsulates the net eddy feedback onto the mean flow.
- The eddy PV flux can be related to the TWA field as an active tracer.
- The 2×2 diffusivity tensor \mathbf{K} , which provides a closure for the eddy PV flux, single-handedly includes the information of eddy momentum fluxes in addition to bringing the GM and Redi diffusivities together.
- For a prognostic closure scheme, we would need to inform the “ κ ”s via the (resolved?) TWA field.

(Extra slide) Defining the buoyancy coordinate for a realistic EOS

With a realistic equation of state (EOS) the vertical coordinate cannot “naively” be defined by potential density (ρ_θ) as the

pressure anomaly ($\phi = \int -g\rho_0^{-1}(\rho_\theta - \rho_0)d\zeta$) does not

translate to a body force in buoyancy coordinates, i.e.

$\nabla_h \phi \neq \tilde{\nabla}_h m$. We, therefore, argue for the use of in-situ density anomaly δ ($= \rho - \rho(\zeta)$) where ρ is the in-situ density and ρ is

a function of only depth. The buoyancy coordinate can then be

defined as $\tilde{b} = -g\rho_0^{-1}\delta$ which removes a large portion of

compressibility; the iso-surfaces of \tilde{b} become close to neutral

surfaces. The formulation of \tilde{b} is analogous to where the

buoyancy reduces to $\tilde{b} = -g\rho_0^{-1}(\rho - \rho_0)$ for a linear EOS.

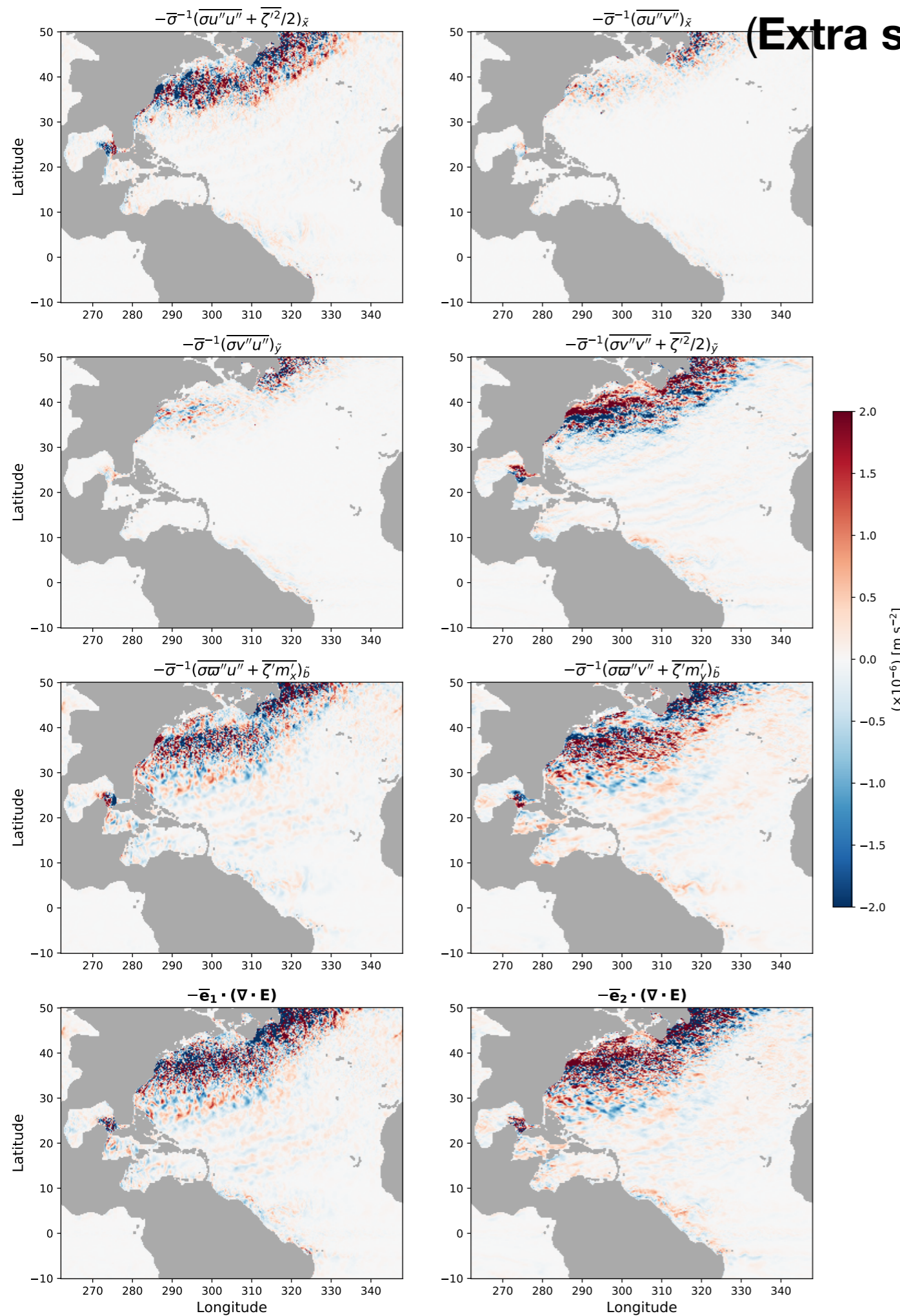
$$\hat{u}_{\tilde{t}} + \hat{u}\hat{u}_{\tilde{x}} + \hat{v}\hat{u}_{\tilde{y}} + \hat{\omega}\hat{u}_{\tilde{b}} - f\hat{v} + \overline{m}_{\tilde{x}} = -\bar{\mathbf{e}}_1 \cdot (\tilde{\nabla} \cdot \mathbf{E}) + \hat{\mathcal{X}}$$

$$\hat{v}_{\tilde{t}} + \hat{u}\hat{v}_{\tilde{x}} + \hat{v}\hat{v}_{\tilde{y}} + \hat{\omega}\hat{v}_{\tilde{b}} + f\hat{u} + \overline{m}_{\tilde{y}} = -\bar{\mathbf{e}}_2 \cdot (\tilde{\nabla} \cdot \mathbf{E}) + \hat{\mathcal{Y}}$$

$$\mathbf{E} = \begin{pmatrix} \widehat{u''u''} + \frac{1}{2\overline{\sigma}} \overline{\zeta'^2} & \widehat{u''v''} & 0 \\ \widehat{v''u''} & \widehat{v''v''} + \frac{1}{2\overline{\sigma}} \overline{\zeta'^2} & 0 \\ \widehat{\omega''u''} + \frac{1}{\overline{\sigma}} \overline{\zeta' m'_{\tilde{x}}} & \widehat{\omega''v''} + \frac{1}{\overline{\sigma}} \overline{\zeta' m'_{\tilde{y}}} & 0 \end{pmatrix}$$

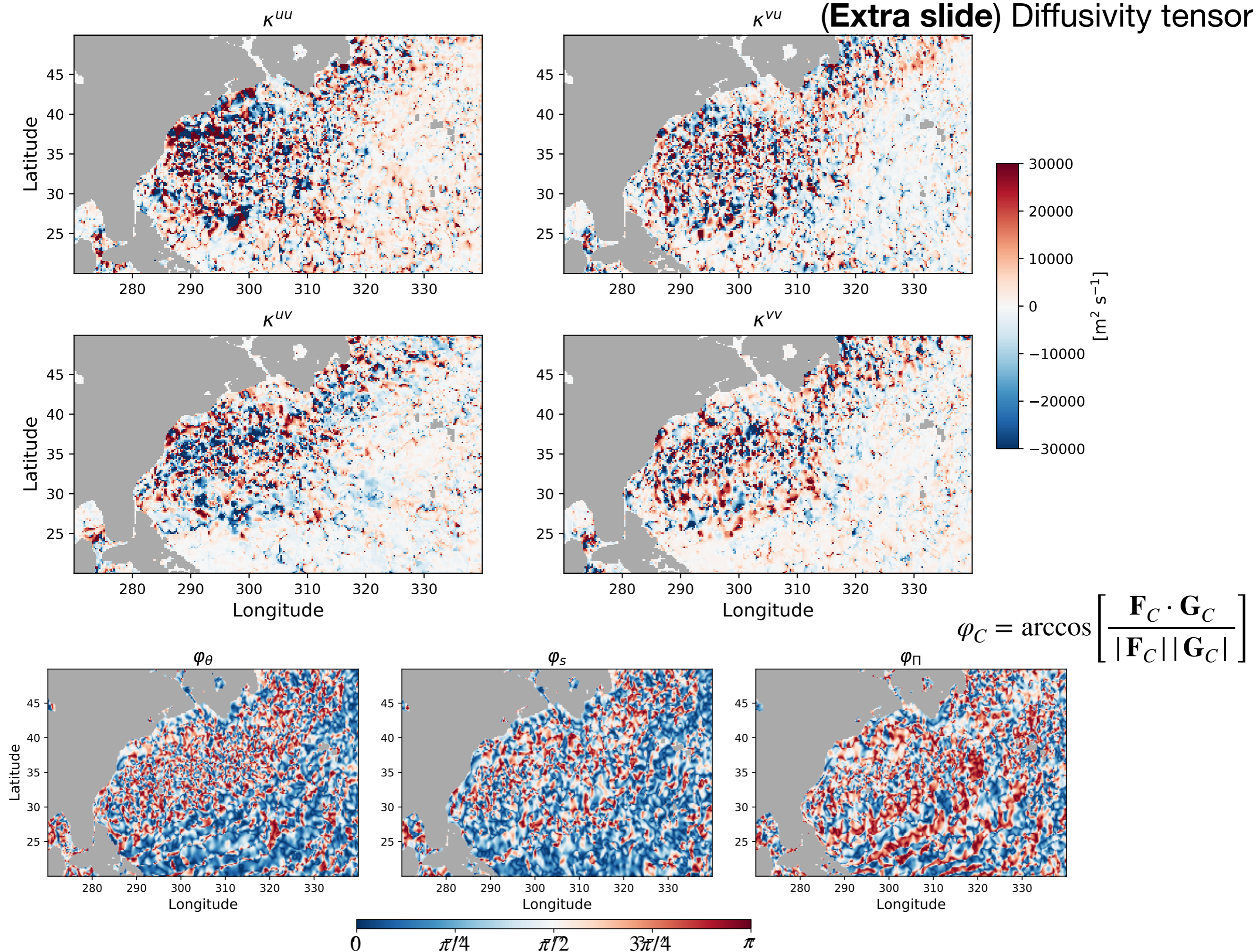
- \mathbf{u}'' ($= \mathbf{u} - \hat{\mathbf{u}}$): the eddy velocity.
- $(\cdot)'$ ($= (\cdot) - \overline{(\cdot)}$): the residual from the ensemble mean.

(Extra slide) E-P flux divergence on Jan. 3, 2008

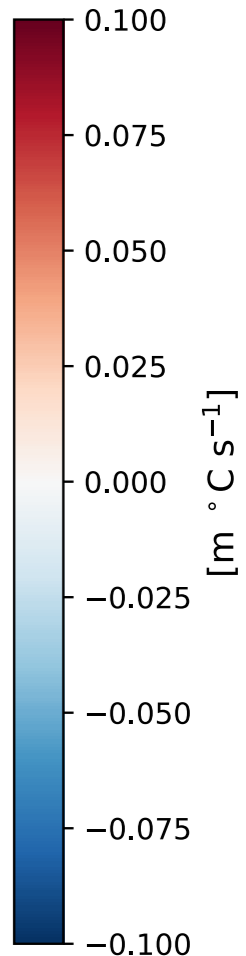
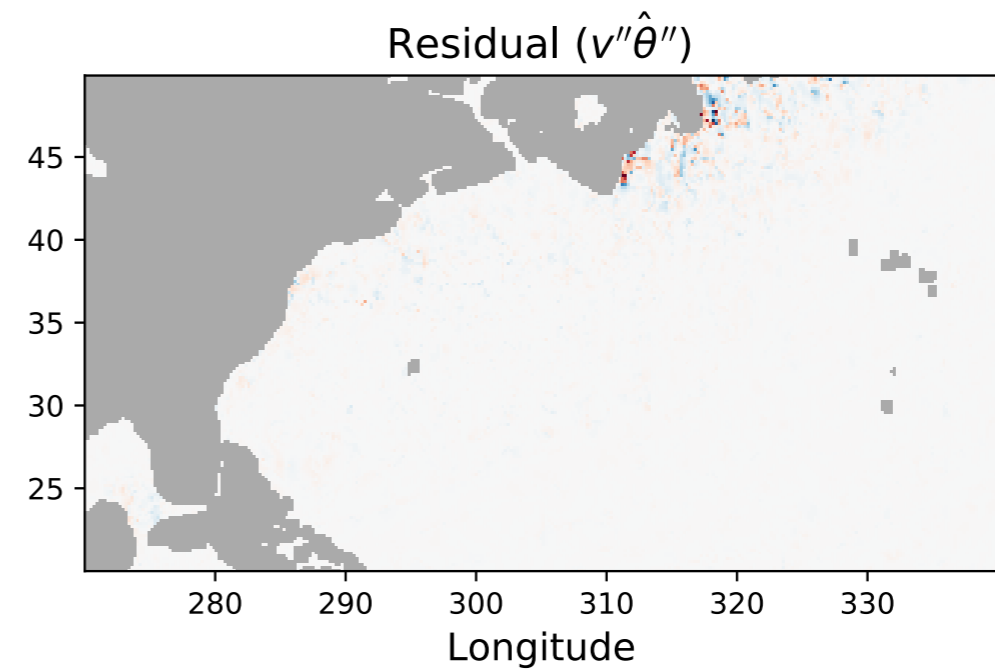
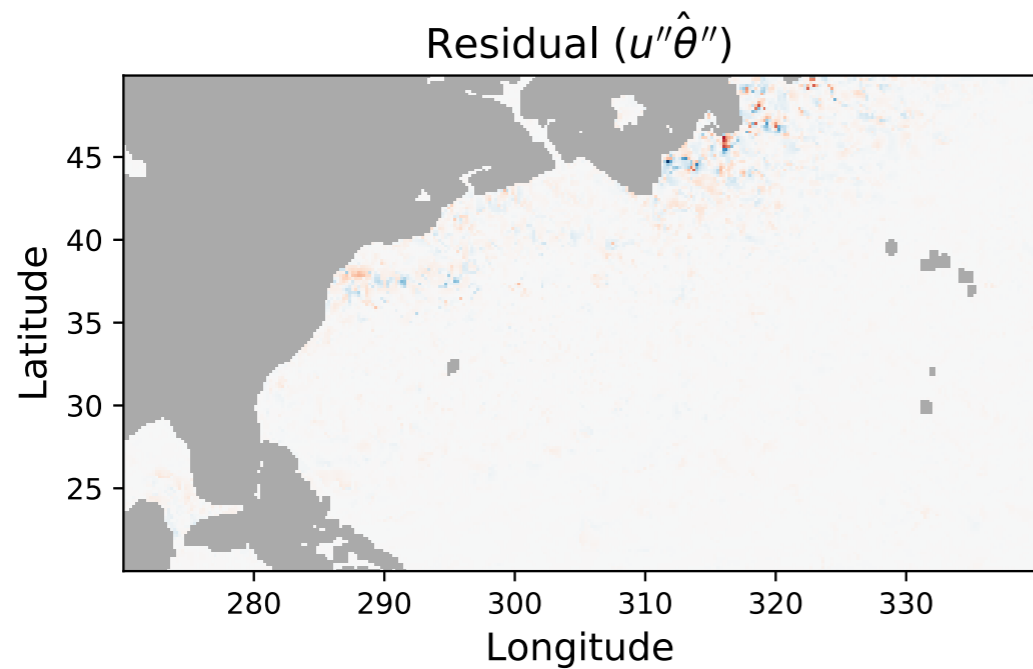
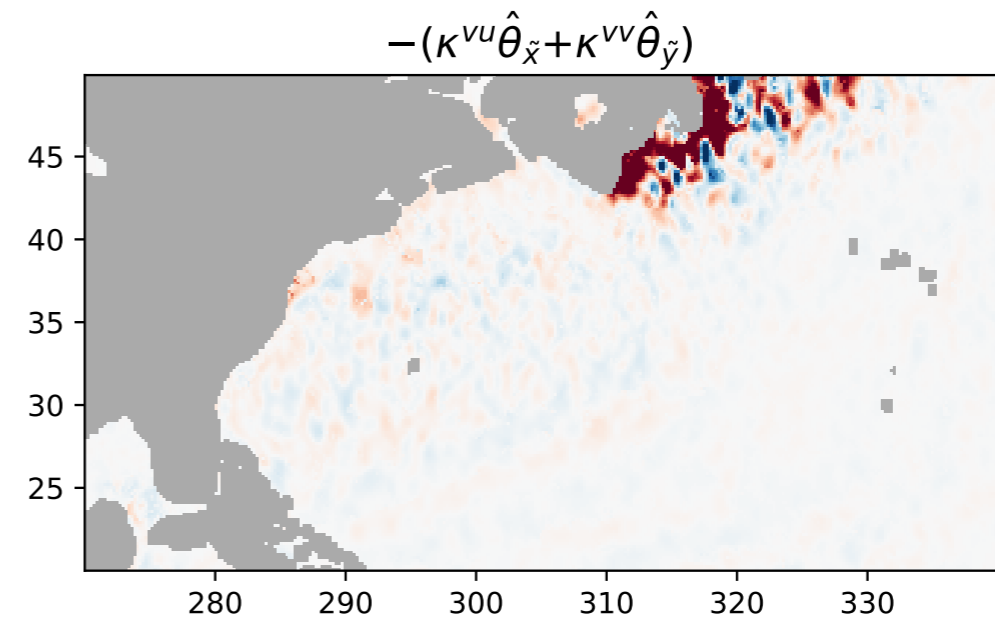
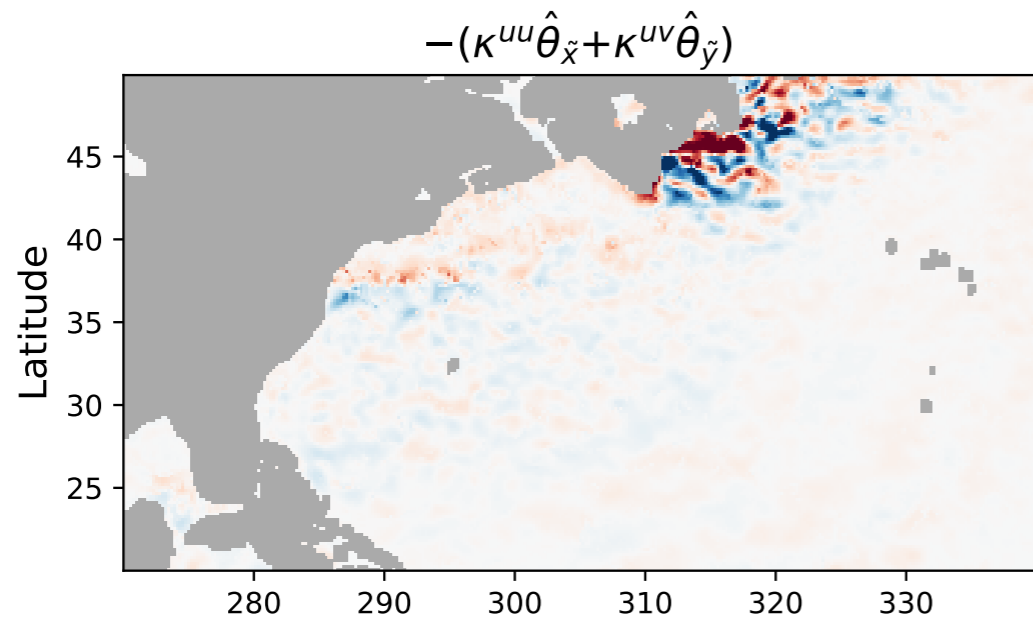
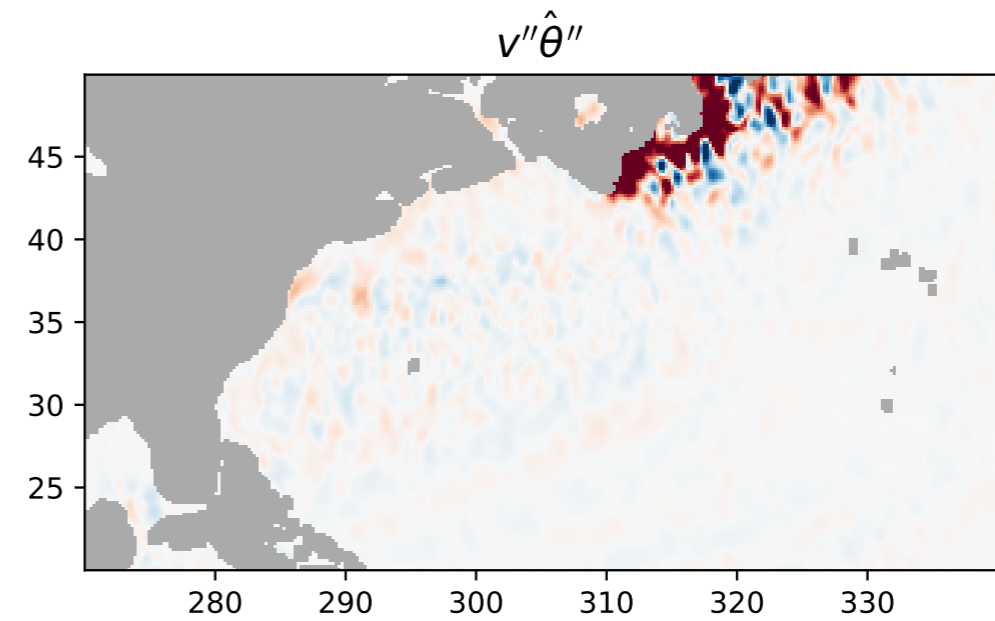
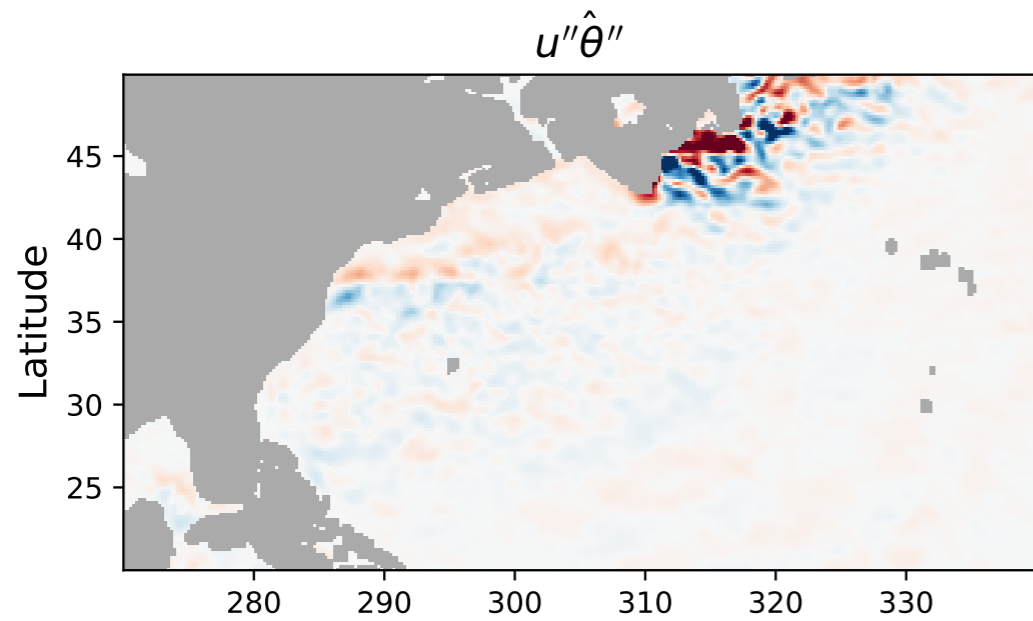


- Each column is laid out so that the sum of the first three rows sum up to the E-P flux divergence shown in the bottom row.
- The terms associated with eddy momentum flux and baroclinic instability tend to cancel each other out.

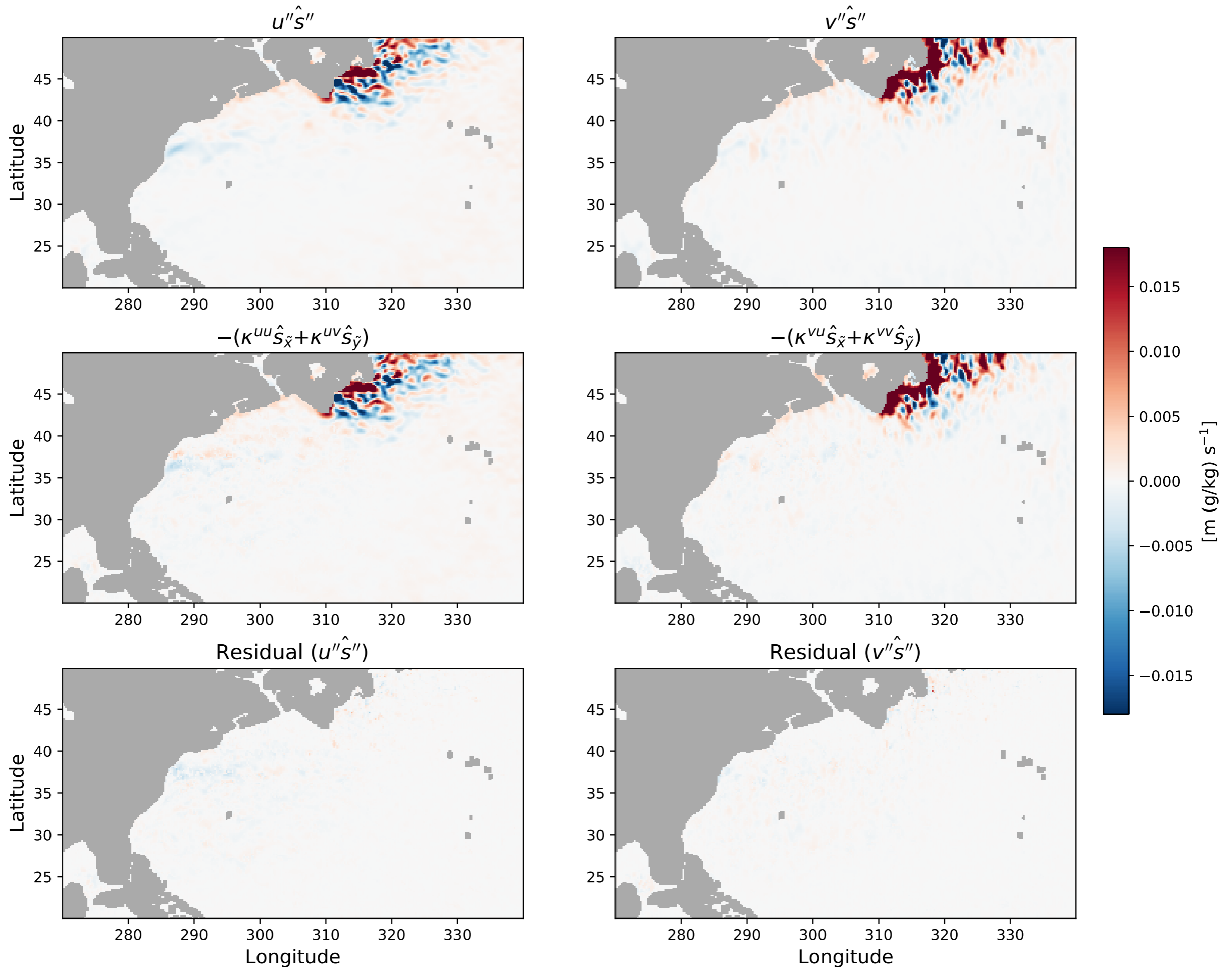
(Extra slide) Diffusivity tensor



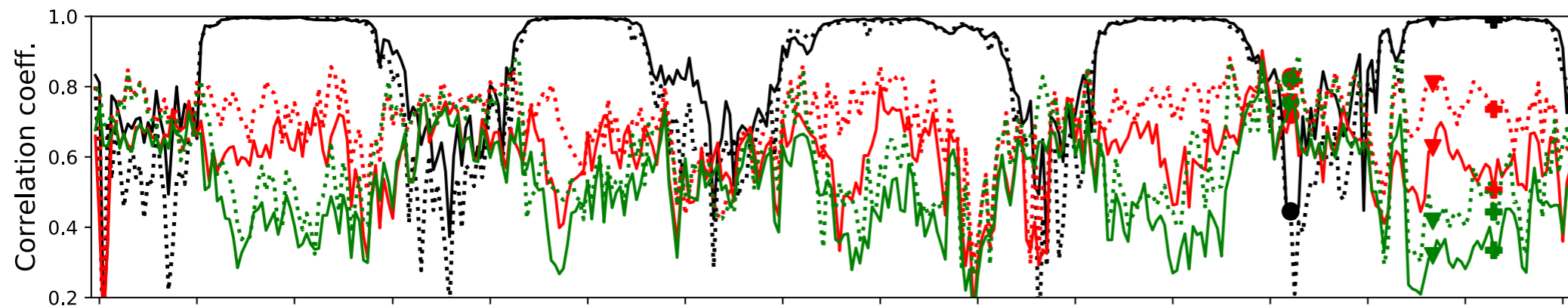
(Extra slide) Temperature parametrization



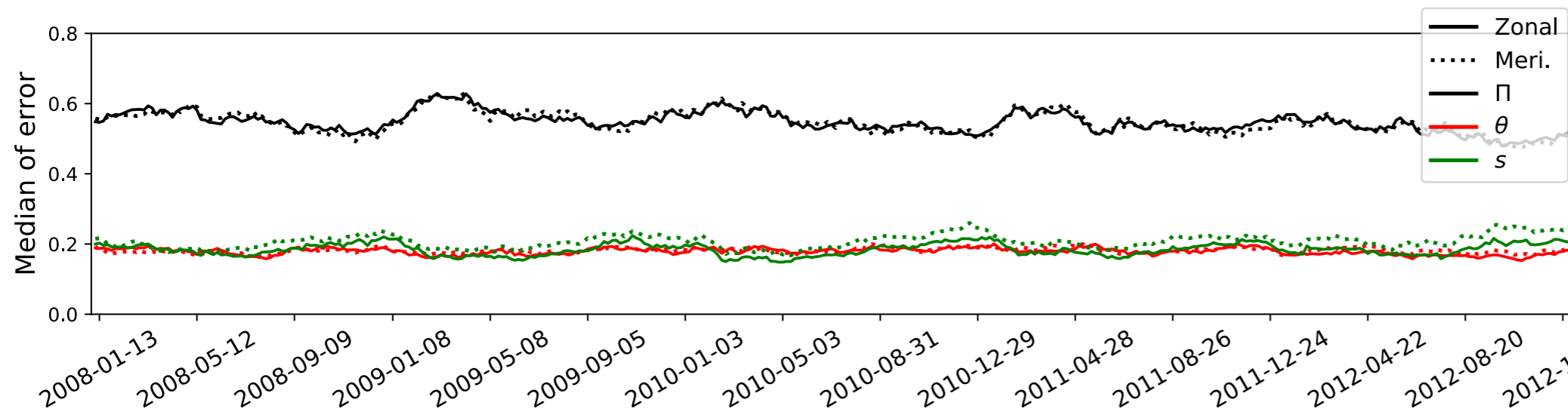
(Extra slide) Salinity parametrization



(Extra slide) Correlation and error of parametrization



$$\frac{\sum [(F_C - \langle F_C \rangle)(F_C^{\text{param}} - \langle F_C^{\text{param}} \rangle)]}{\sqrt{\sum (F_C - \langle F_C \rangle)^2} \sqrt{\sum (F_C^{\text{param}} - \langle F_C^{\text{param}} \rangle)^2}}$$



$$\frac{|F_C - F_C^{\text{param}}|}{|F_C|}; F_C^{\text{param}} = G_C \cdot K_C$$

Jan. 18, 2012

Jul. 11, 2012

Sep. 24, 2012

