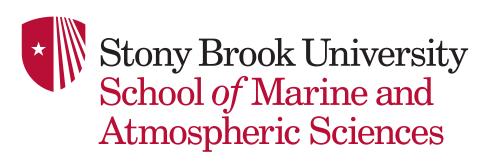


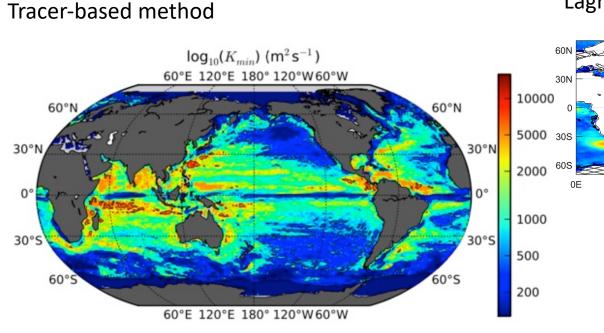
## Vertical structure of tracer diffusivity in an idealized basin circulation model

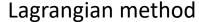
Wenda Zhang and Christopher L.P. Wolfe School of Marine and Atmospheric Sciences Stony Brook University

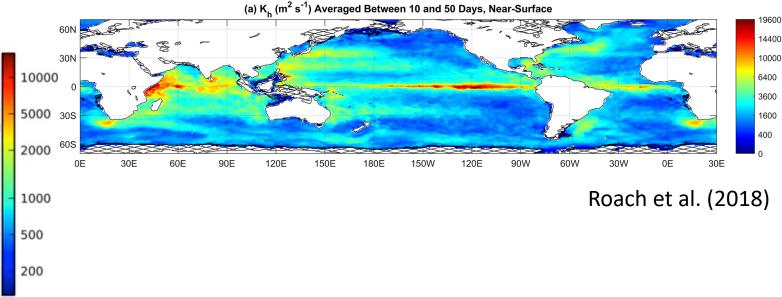


### Observations of mesoscale mixing rate

Parameterization for eddy fluxes:  $\overline{u'\tau'} = -\mathbf{K}\nabla\overline{\tau}$ 







Abernathey et al. (2013)

Groeskamp et al. (2020)

# Vertical structure of diffusivity

Based on the suppressed mixing length theory (Ferrari and Nikurashin, 2010), Groeskamp et al., 2020 proposed an estimate of full-depth diffusivity,

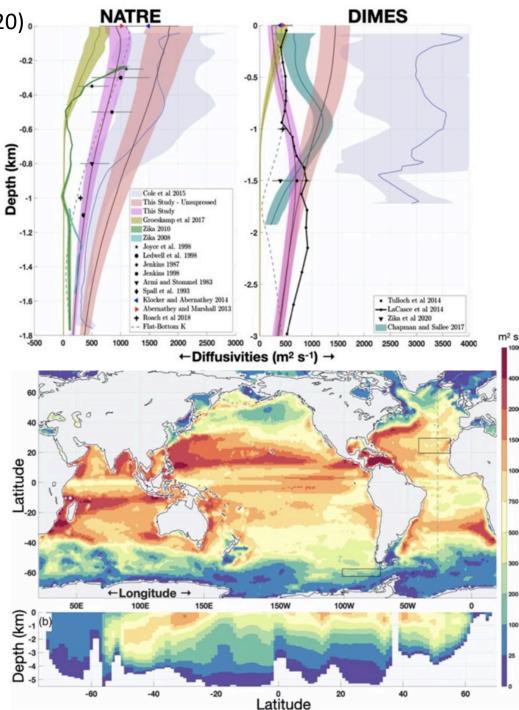
 $K(x, y, z) = \Gamma \phi(z) \sqrt{2 \text{EKE}_0} L_d \times \min(S^x, S^y),$ 

where  $\phi(z)$  is a **vertical structure** function of the first "**surface mode**" (LaCasce 2017), EKE<sub>0</sub> is the EKE at surface,  $L_d$  is the deformation radius of the the first "surface mode".  $S^x$  and  $S^y$  are the **suppression factors** accounting for eddy propagation relative to mean flow, and  $\Gamma$ is the mixing efficiency.

#### **Questions:**

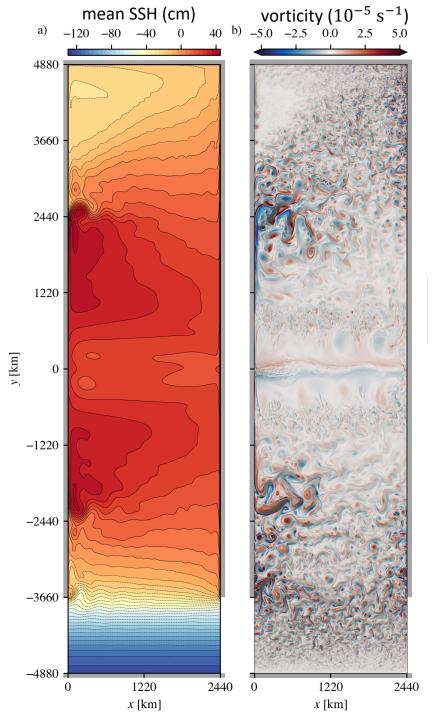
Does mixing length theory work over full water column? Can the vertical structure of EKE be reconstructed from hydrography?

How does this theory apply to the entire diffusivity tensor?



### Numerical Model

- This model is an idealized configuration of the MITgcm (Marshall et al. 1997a,b; Campin et al. 2020) used by previous studies (e.g., Wolfe et al. 2008; Wolfe and Cessi 2009, 2010, 2011)
- A two-hemisphere basin on an equatorial β-plane with a flat bottom
- Extent of the domain: 2440 km in zonal direction, 9880 km in meridional direction and a uniform depth of 2440 m
- The southernmost eighth of the domain is **zonally reentrant**
- 5.4 km horizontal resolution; 20 vertical levels
- Forcing: zonally uniform **zonal winds** and a relaxation to a zonally uniform **surface temperature distribution**

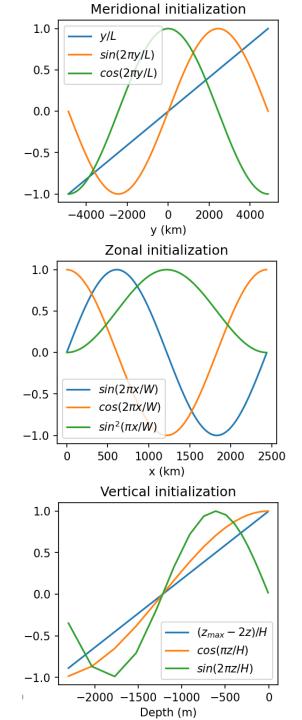


## Diagnosing ${\bf K}$

- 9 components of diffusivity tensor **K** are diagnosed using the tracer inversion method of Bachman et al. (2015)
- Advect a total of 27 passive tracers with 9 initial conditions, each set relaxed to their initial conditions with 3 different relaxation rates (1 year, 3 years, and 6 years)
- Solve for diffusivity tensor **K** in a least-squares sense:

$$K_{ij} = -\overline{u_i'\tau_\alpha'} \big[\partial_j \bar{\tau}_\alpha\big]^{\dagger}$$

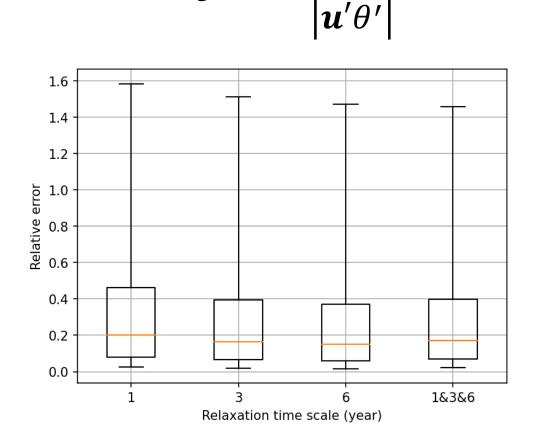
where (i, j) = (x, y, z),  $\alpha = 0, ..., 26$  is the tracer number,  $\overline{(\cdot)}$  is a 20-year and 152-km spatial average, and  $(\cdot)^{\dagger}$  indicates the Moore-Penrose pseudoinverse

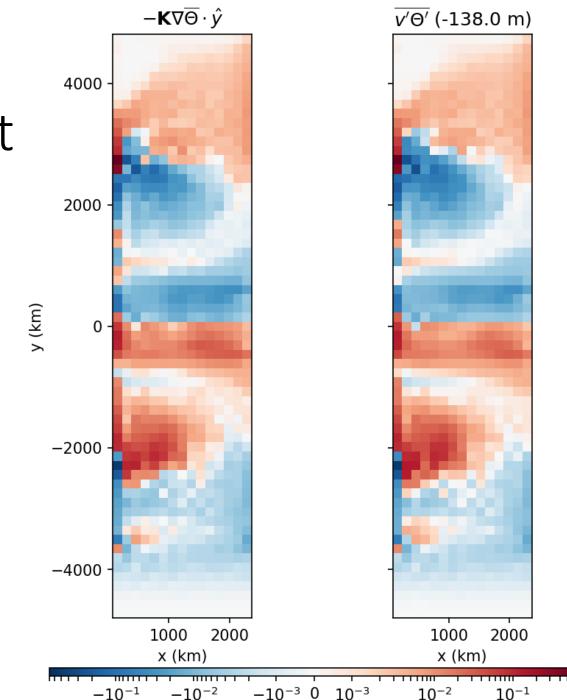


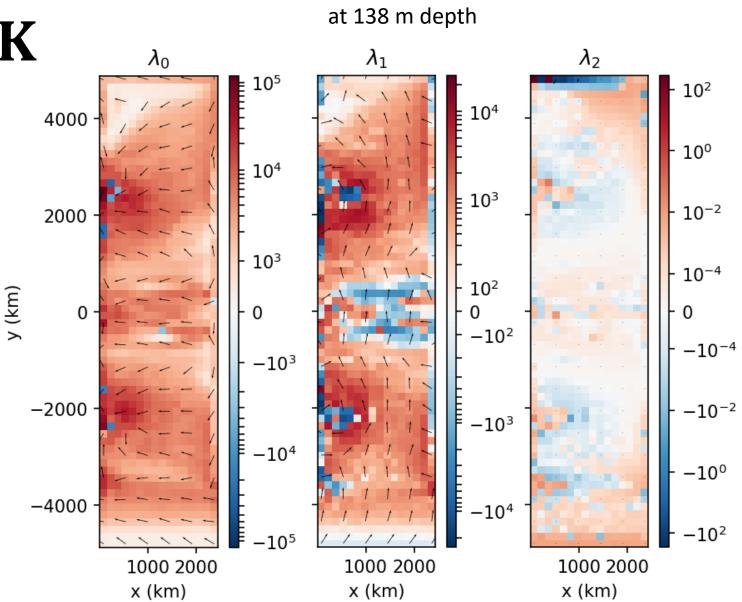
# Reconstruction of buoyancy flux is excellent

• Relative error of reconstructed flux:  $|\overline{u'\theta'} + \mathbf{K}\nabla\overline{\theta}|$ 

 $\mathcal{E} =$ 







Eigenvalues in m<sup>2</sup> s<sup>-1</sup> (colors) and eigenvectors (arrows)

# Eigenvalues of symmetric part of **K**

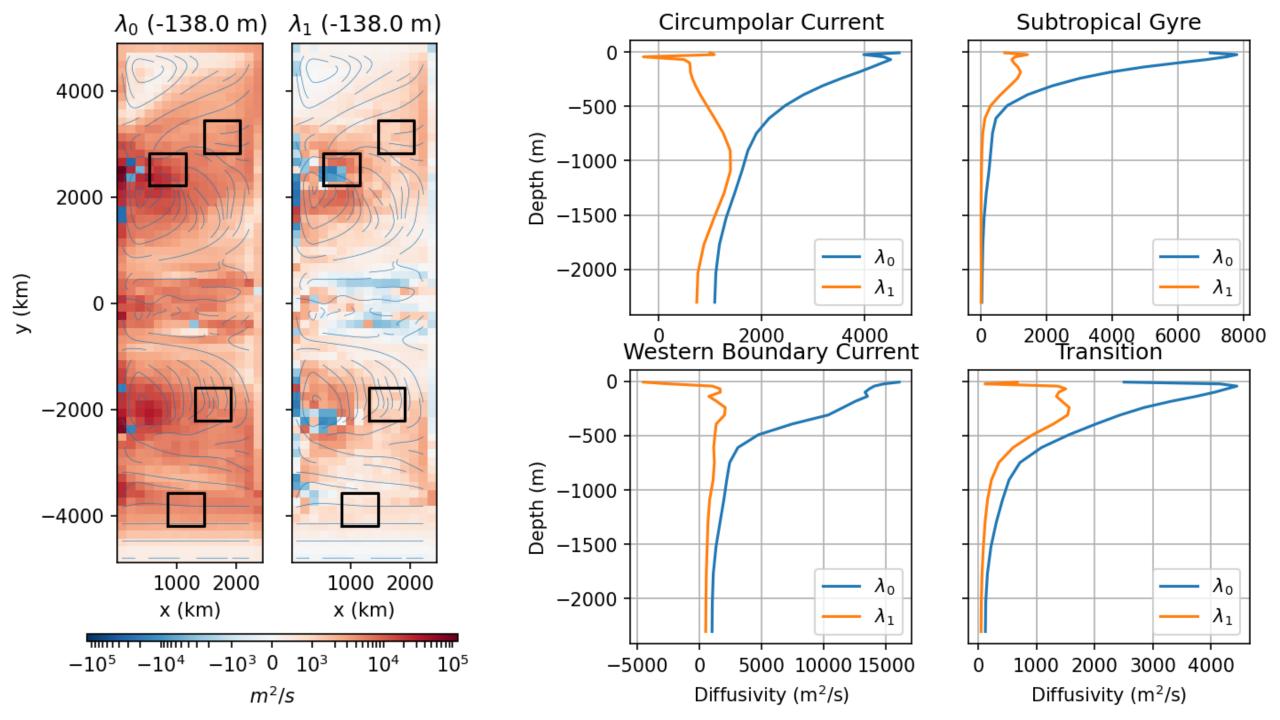
•  $\mathbf{K} = \mathbf{S} + \mathbf{A}$  where

$$\mathbf{S} = \frac{\mathbf{K} + \mathbf{K}^T}{2}$$
,  $\mathbf{A} = \frac{\mathbf{K} - \mathbf{K}^T}{2}$ 

• The **symmetric** tensor **S** is diffusive and can be decomposed as

$$\mathbf{S} \cdot \mathbf{v}_i = \lambda_i \cdot \mathbf{v}_i$$

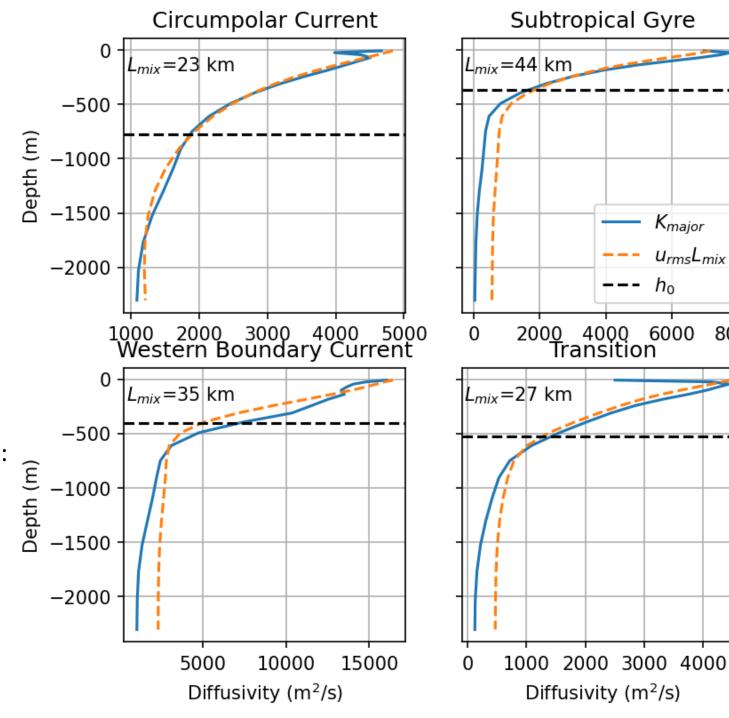
where  $\lambda_i$  (i = 0, 1, 2) are the eigenvalues of **S**, representing the diffusivity along eigenvectors  $\mathbf{v}_i$ 



## Major diffusivity $(\lambda_0)$

- Mixing length theory:
  - $K \sim u_{rms} L_{mix}$
- $u_{rms}$  is diagnosed from model, and  $L_{mix}$  is assumed to be constant over depth and obtained by least squares fitting.
- The scale height of stratification  $h_0$ :

$$h_{0} = -\frac{\int_{-H}^{0} z N^{2} dz}{\int_{-H}^{0} N^{2} dz}$$



*K<sub>major</sub>* 

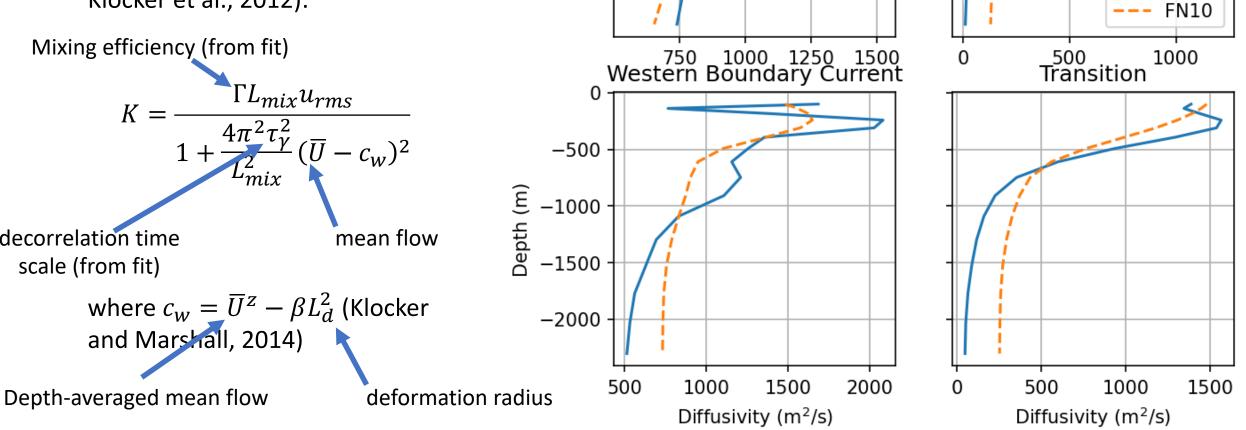
 $h_0$ 

6000

*u*<sub>rms</sub>L<sub>mix</sub>

8000

#### Circumpolar Current Subtropical Gyre 0 Minor diffusivity -500 $(\lambda_1)$ Depth (m) -1000 Suppressed mixing length theory -1500 (Ferrari and Nikurashin, 2010; -2000 Klocker et al., 2012): Mixing efficiency (from fit) 750 1000 1250 1500 Western Boundary Current 0 $\frac{\Gamma L_{mix} u_{rms}}{4\pi^2 \tau_{\nu}^2 \tau_{\overline{\nu}}^2}$ 0 K = $(\overline{U} - c_w)^2$ -500 -mix -1000decorrelation time



K<sub>minor</sub>

# "Surface mode" used by Groeskamp et al. (2020)

• Assuming a wave-like solution to the linear QG PV equation yields an equation for the vertical structure

$$\frac{d}{dz}\left(\frac{f_0^2}{N^2}\frac{d\phi}{dz}\right) + \frac{\phi}{L_d^2} = 0$$

 Traditional solution assumes a rigid lid and flat bottom:

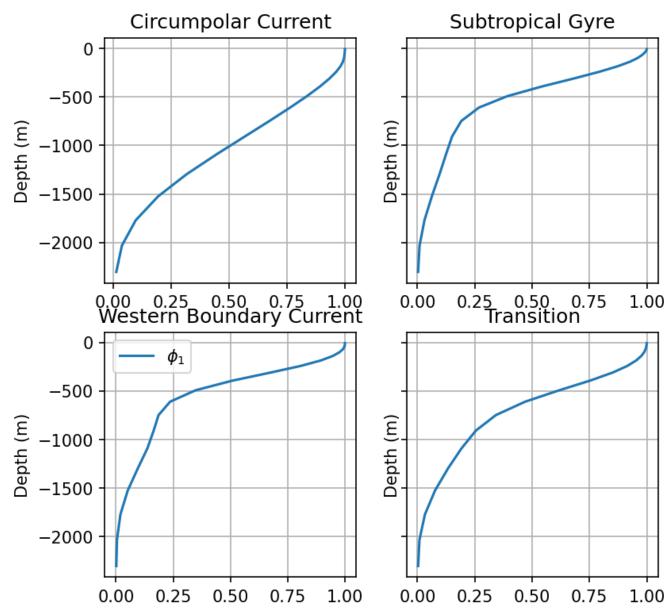
$$\frac{d\phi}{dz} = 0 \text{ at } z = 0, -H.$$

• "Surface mode" solution (de La Lama et al. 2016; LaCasce 2017) assumes a **rough bottom**:

$$\frac{d\phi}{dz} = 0 \text{ at } z = 0, \qquad \phi = 0 \text{ at } z = -H$$

 Vertical structure of current meter data is found to be similar as the first "surface mode" (de La Lama et al. 2016).

#### Surface mode solutions

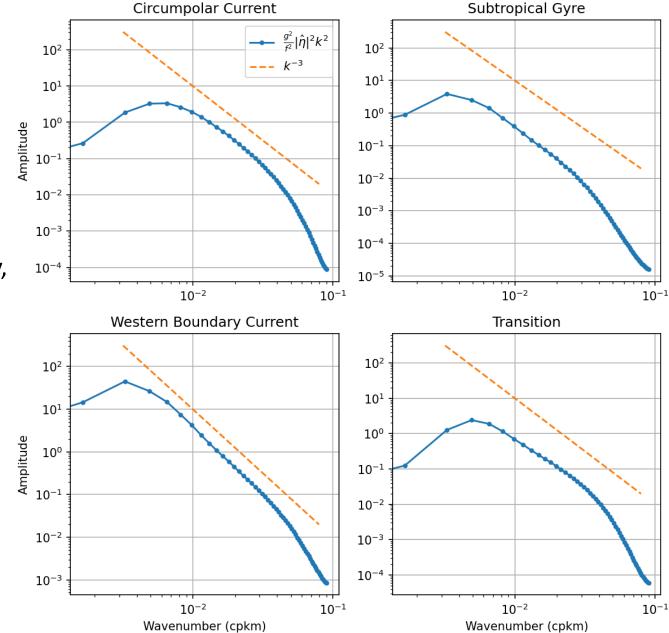


#### EKE spectrum from SSH spectrum ( $600 \text{ km} \times 600 \text{ km}$ window)

## SQG solution

- Eddy streamfunction  $\psi$  can be decomposed into an **interior solution**,  $\psi_{int}$ , with zero buoyancy anomaly on the boundaries and a **surface solution**,  $\psi_{sur}$ , with zero interior PV anomaly (Lapeyre and Klein, 2006)
- $\psi_{sur}$  is driven by the surface buoyancy anomaly, described by the **surface quasigeostrophic** (SQG) dynamics (Held et al., 1995)
- $\psi_{sur}$  is solved numerically in Fourier space,

$$\begin{split} & \frac{\partial}{\partial z} \left( \frac{f_0^2}{N^2} \frac{\partial \hat{\psi}_{sur}}{\partial z} \right) - \kappa^2 \, \hat{\psi}_{sur} = 0, \\ & \hat{\psi} \Big|_{z=0} = \frac{g}{f_0} \hat{\eta}, \qquad \frac{\partial \hat{\psi}}{\partial z} \Big|_{z=-H} = 0, \\ & \text{where } \kappa^2 = k^2 + l^2. \end{split}$$



## Reconstruction of EKE

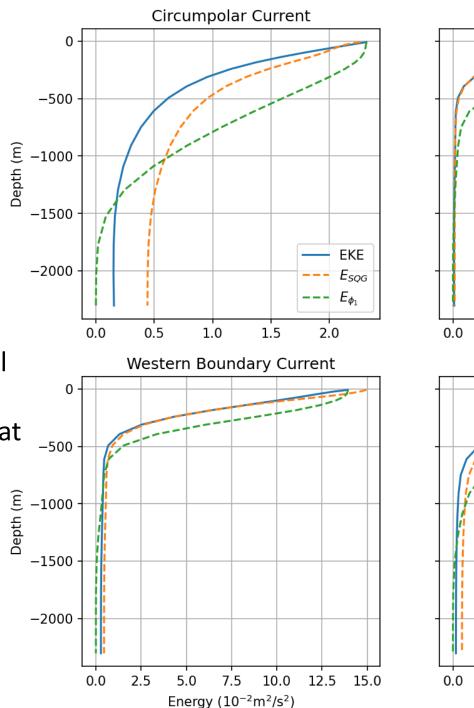
• EKE from the first "surface mode" (Groeskamp et al., 2020):

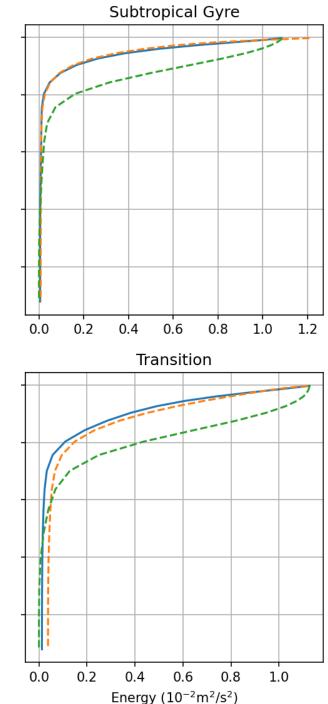
 $E_{\phi_1} = \phi_1(z)^2 E K E_0,$ 

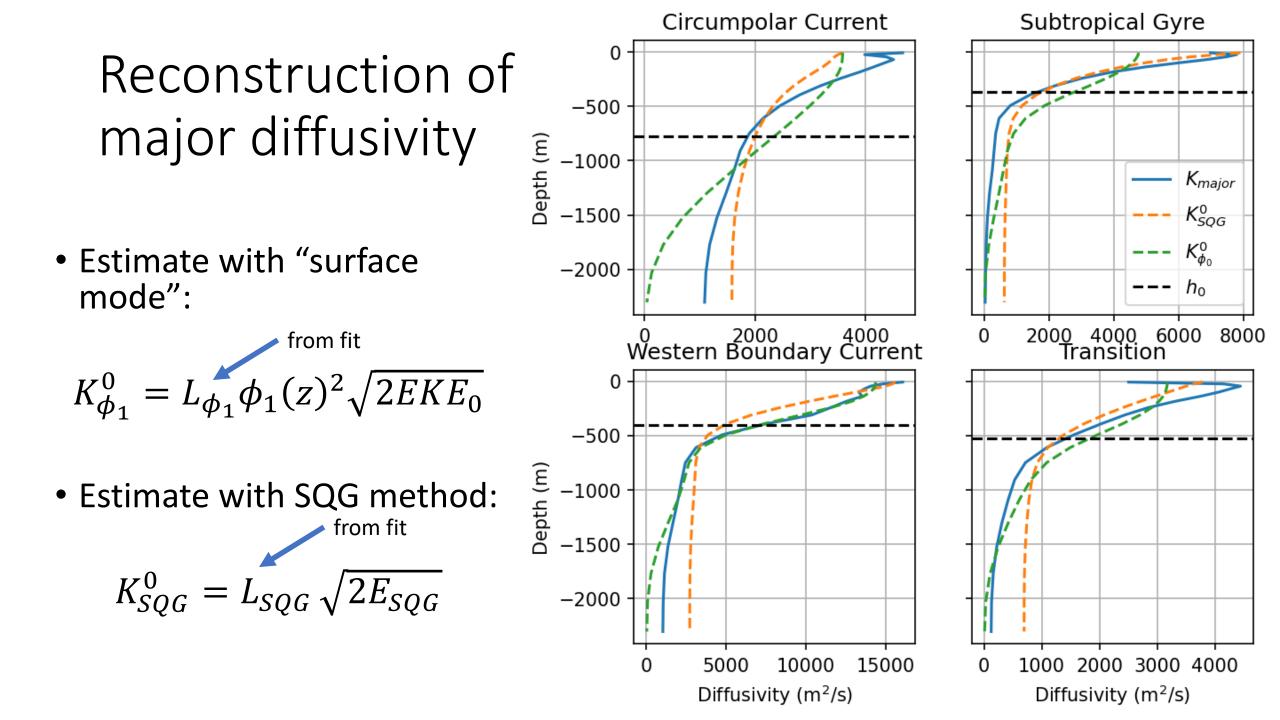
where  $\phi_1$  is the normalized vertical structure function of the first "surface mode",  $EKE_0$  is the EKE at surface.

• EKE from SQG method:

$$E_{SQG} = \frac{1}{2} \sum_{k,l} \kappa^2 \left| \hat{\psi}(\boldsymbol{k}, \boldsymbol{z}) \right|^2$$





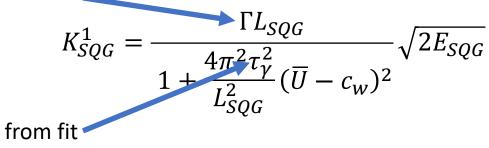


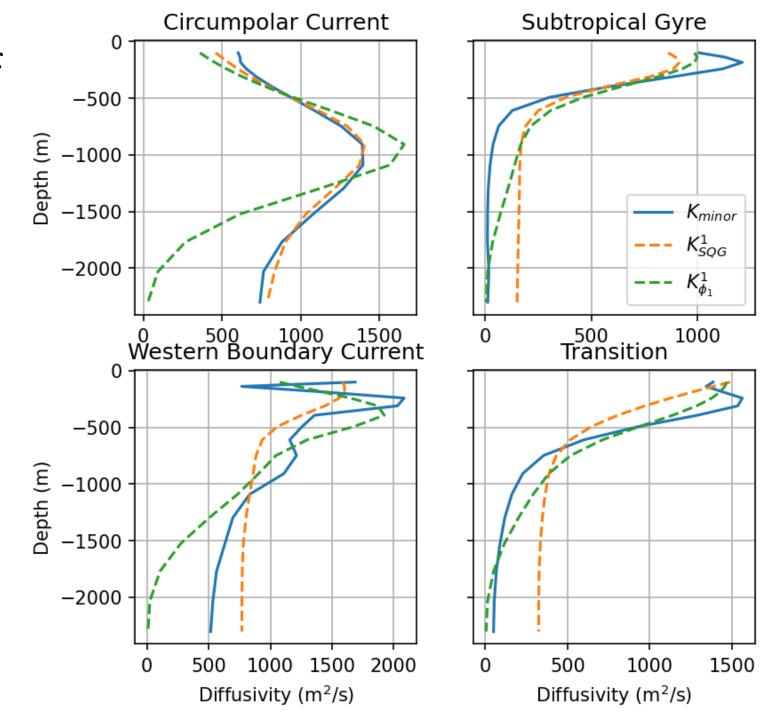
# Reconstruction of minor diffusivity

• Estimate with "surface mode":

from fit  $K_{\phi_1}^1 = \frac{\Gamma L_{\phi_1} \phi_1(z)}{1 + \frac{4\pi^2 \tau_\gamma^2}{L_{\phi_1}^2} (\overline{U} - c_w)^2} \sqrt{2EKE_0}$ from fit

 Estimate with SQG method: from fit





## Summary

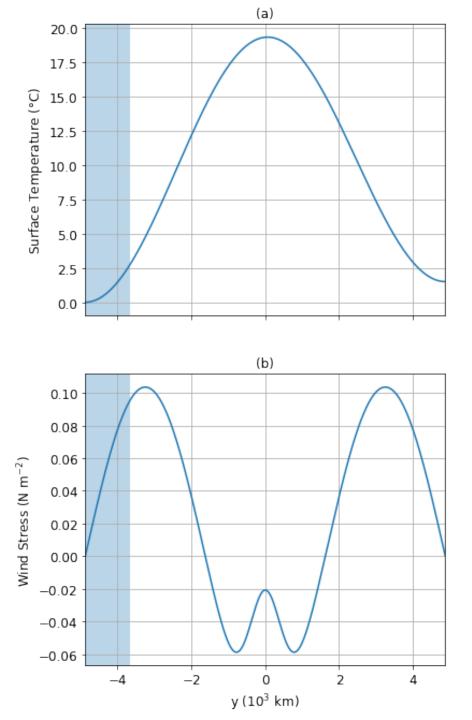
- Mixing length theory assuming **constant mixing length over depth** works well for the **major diffusivity** in the upper levels, but often overestimates the diffusivity at depth away from the channel.
- The **minor diffusivity** is consistent with the suppressed mixing length theory, in which the diffusivity is **modulated by eddy propagation relative to the mean flow**.
- Vertical structure of **EKE** is better reconstructed by the **SQG method** than the first "surface mode".
- SQG method performs better than the first "surface mode" in reconstructing the **diffusivities** in upper levels; the former (latter) tends to overestimate (underestimate) the diffusivities in the deep levels.

#### **Future work**

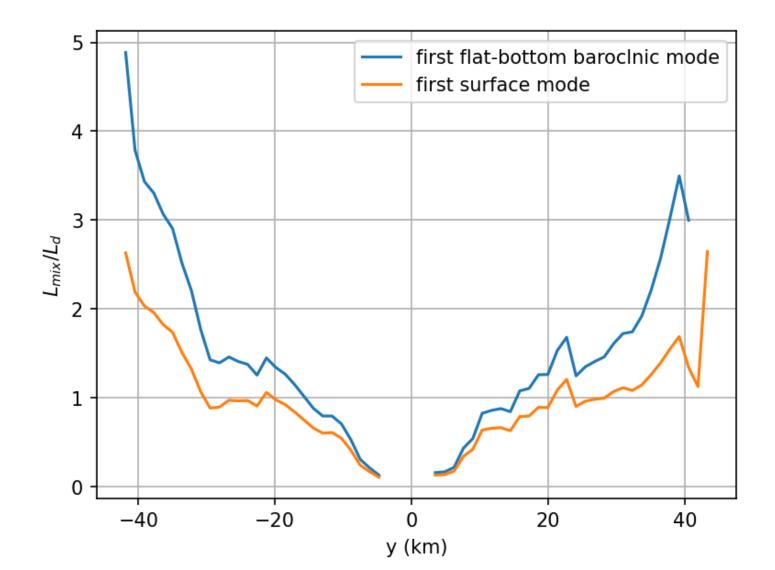
- The current model has a flat bottom. We are planning to add topography to further compare the reconstruction with the "surface mode" and SQG method.
- Determination of the **mixing length** and **decorrelation time scale** is still an open question.

## Extra slides

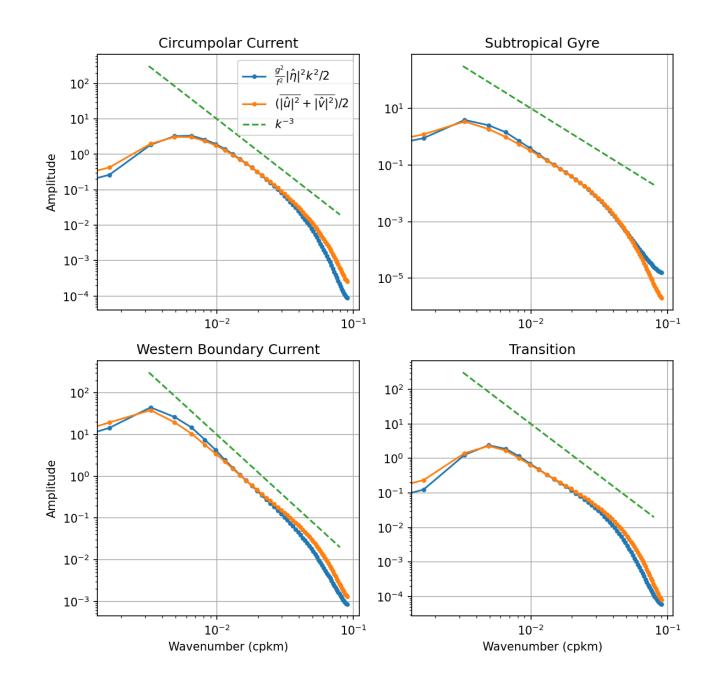
### Forcings (extra slide)



### Mixing length (extra slide)

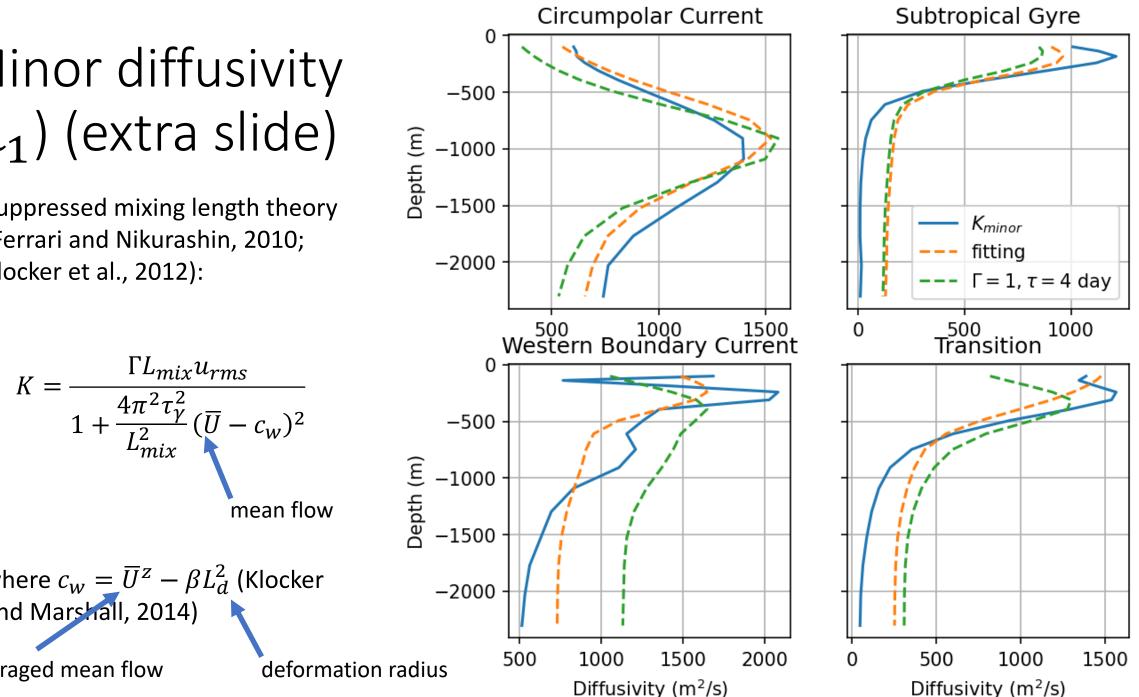


## Energy spectrum (extra slide)



## Minor diffusivity $(\lambda_1)$ (extra slide)

 Suppressed mixing length theory (Ferrari and Nikurashin, 2010; Klocker et al., 2012):



mean flow where  $c_w = \overline{U}^z - \beta L_d^2$  (Klocker and Marshall, 2014)

Depth-averaged mean flow

Diffusivity  $(m^2/s)$