

A higher-order ice-flow model applied to the dynamics of Greenland's outlet glaciers

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Motivation for improved Ice Sheet Models

Current models do not capture observed ice sheet behaviour¹, because:

- (1) fundamental physics are lacking (e.g. solving simplified equations, negating realistic simulation of outlet glaciers and ice streams)
- (2) processes of fundamental importance are not accounted for (e.g. simplified, static treatment of basal boundary conditions, ignoring atmos. and ocean coupling, etc.)

Accurate predictions of future sea-level rise (SLR) will require advanced models which can demonstrate skill at reproducing and explaining recent dramatic ice sheet behaviors.

¹ IPCC (2007)

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- (1) accounting for both vertical AND horizontal stresses (e.g. HO models)
- (2) improved basal and lateral BCs (e.g. plastic bed, floating ice at margins)
- (3) detailed treatment of grounding line¹ ... e.g. Nick/Vieli and others² model mimicks Helheim glacier over last ~10 yrs (and, arguably, provides explanation for observed behaviour)

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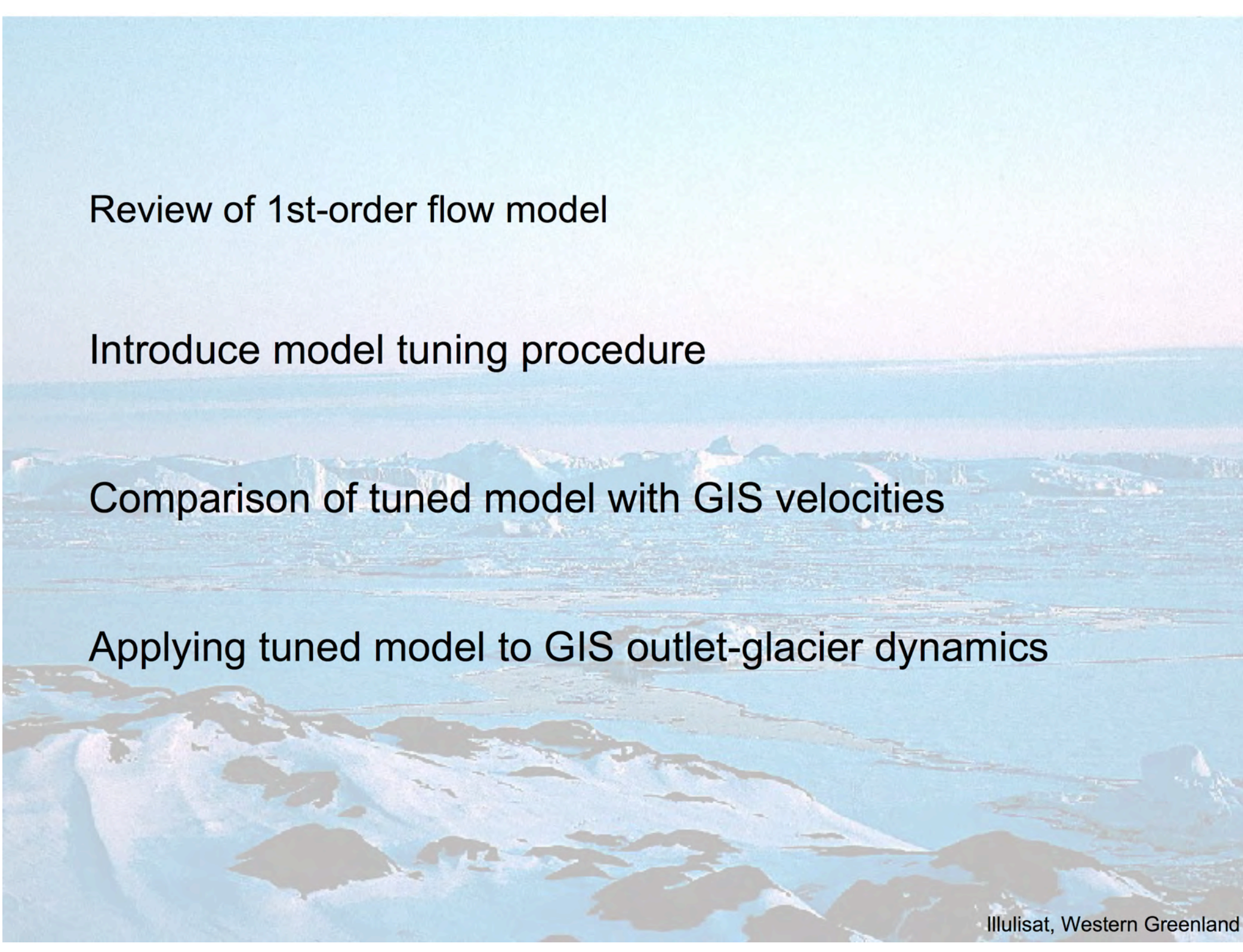
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Today: Introduce an “easy” tuning method that may allow us to fill the gap and make some useful predictions using improved, stand-alone models (along with some simple, perturbation experiments).

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Introduce model tuning procedure

Comparison of tuned model with GIS velocities

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Equations of Stress Equilibrium in Cartesian Coordinates (Stokes Flow)

Assume static balance of forces by ignoring acceleration

$$x: \frac{\partial \tau_{xx}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} - \frac{\partial P}{\partial x} = 0$$

$$y: \frac{\partial \tau_{yy}}{\partial y} + \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \tau_{yz}}{\partial z} - \frac{\partial P}{\partial y} = 0$$

$$z: \frac{\partial \tau_{zz}}{\partial z} + \frac{\partial \tau_{zy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial x} - \frac{\partial P}{\partial z} = \rho g$$

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... apply some scaling based on H/L ...

$$\dot{\epsilon}_{xz} = \frac{1}{2} \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right)$$



$$\frac{\partial \sigma_{xz}}{\partial x} = \frac{\partial \sigma_{yz}}{\partial y} \approx 0$$

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$$\dot{\epsilon}_{xz} = \frac{1}{2} \left(\frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \rightarrow \dot{\epsilon}_{xz} = \frac{1}{2} \left(\frac{\partial u}{\partial z} \right)$$

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“First-Order” Approximation

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... integrate w.r.t. z to get explicit expression for P ...

First Order Approximation (unscaled)

...use definition of deviatoric stress to eliminate vertical-normal stress deviator in horiz. equations ...

$$x: \frac{\partial \tau_{xx}}{\partial x} - \frac{\partial \tau_{zz}}{\partial x} + \frac{\partial \tau_{xy}}{\partial y} + \frac{\partial \tau_{xz}}{\partial z} = \rho g \frac{\partial s}{\partial x}$$

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$$\tau_{ij} = B \dot{\epsilon}_e^{\frac{1-n}{n}} \dot{\epsilon}_{ij}, \quad B = B(T) \quad (\text{Glen's law})$$

$$\dot{\epsilon}_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right) \quad (\text{strain rate tensor})$$

$$2\dot{\epsilon}_e = \dot{\epsilon}_{ij} \dot{\epsilon}_{ij} \quad (\text{effective strain rate})$$

$$\eta \equiv \frac{1}{2} B \dot{\epsilon}_e^{\frac{1-n}{n}} \quad (\text{effective viscosity})$$

$$\tau_{ij} = 2\eta \dot{\epsilon}_{ij} \quad (\text{constitutive relation})$$

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...use constitutive relation to write stresses in terms of strain rates and eff. visc., write strain rates in terms of vel. grads. ...

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$$x: \quad 2 \frac{\partial}{\partial x} (2\eta \dot{\epsilon}_{xx}) + \frac{\partial}{\partial x} (2\eta \dot{\epsilon}_{yy}) + \frac{\partial}{\partial y} (2\eta \dot{\epsilon}_{xy}) + \frac{\partial}{\partial z} (2\eta \dot{\epsilon}_{xz}) = \rho g \frac{\partial s}{\partial x}$$

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$$x: 4 \frac{\partial}{\partial x} \left(\eta \frac{\partial u}{\partial x} \right) + 2 \frac{\partial}{\partial x} \left(\eta \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial y} \left[\eta \left(\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} \right) \right] + \frac{\partial}{\partial z} \left(\eta \frac{\partial u}{\partial z} \right) = \rho g \frac{\partial s}{\partial x}$$

First Order Approximation (solution)

Solve for u by moving v terms to RHS, vice versa for v
(operator splitting)

Recover w through continuity

Equations discretized using Finite Difference Method

Vertical coordinates transformed to sigma coordinates

Iterate on effective viscosity using “unstable manifold correction¹”

Conservation of energy (heat balance model) similar to
GLIMMER model

Surface and basal boundary conditions are fully HO
(not 0-order approx.)

¹Hindmarsh and Payne (*Ann. Glac.*, 1996); Pattyn (*JGR*, 2003)

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How do model and target¹ velocity fields compare?

¹Here, GIS balance velocities from Bamber et al. (*J.Glac.*, **46**, 2000)

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 - to first order, both equally important to dynamics, but the latter is more poorly constrained
- (3) must accept the use of a “compensatory accumulation rate” (=ss SMB implied by tuned velocity field ... may or may not be a good representation of actual SMB). OK if main concern is the dynamic response?

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- Observed velocities (InSAR) are not in equilibrium with known or actual geometry.
- InSAR vels are at a much higher resolution than our current geometry data, and so resolve features our model cannot.
- Arguably, better to tune sliding to (more conservative) ss velocities than to (more erratic) transient observed velocities (e.g. accelerated fields from Jak., Kang., and Hel. glaciers).

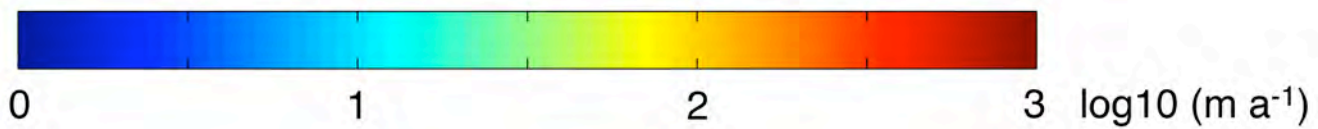
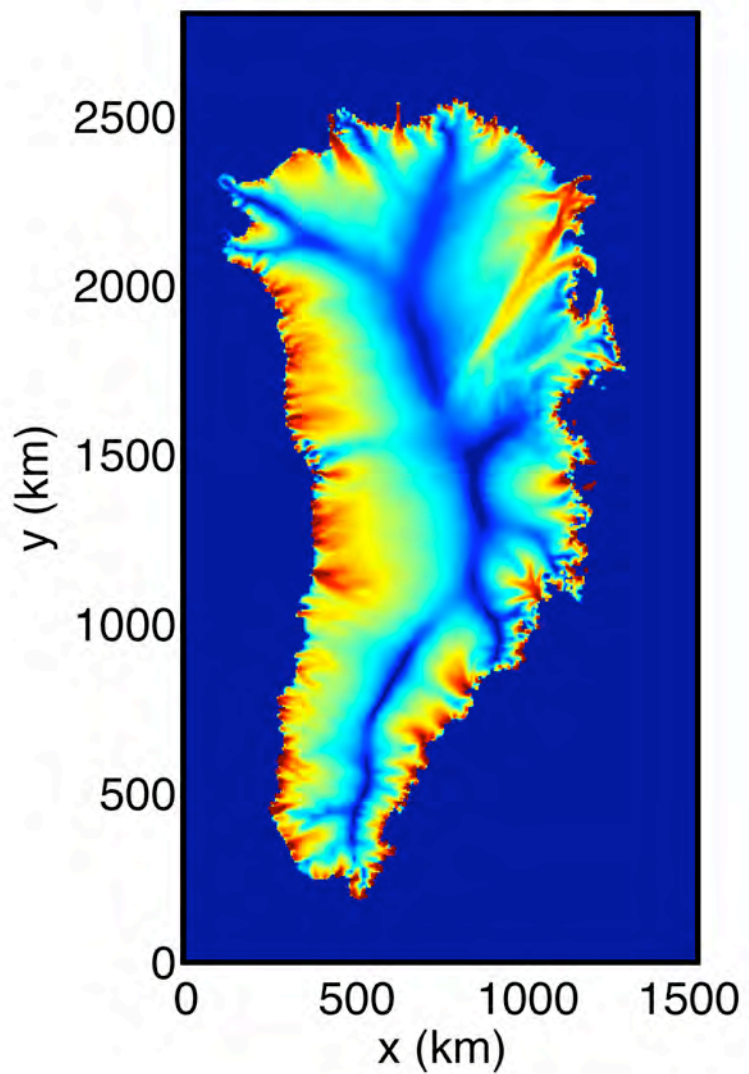
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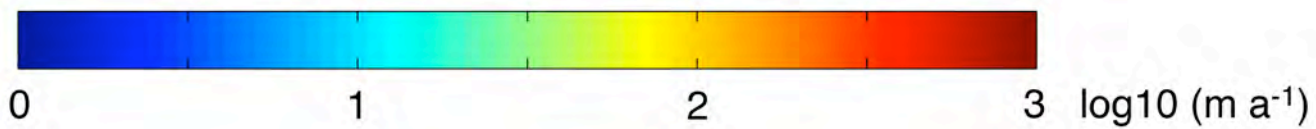
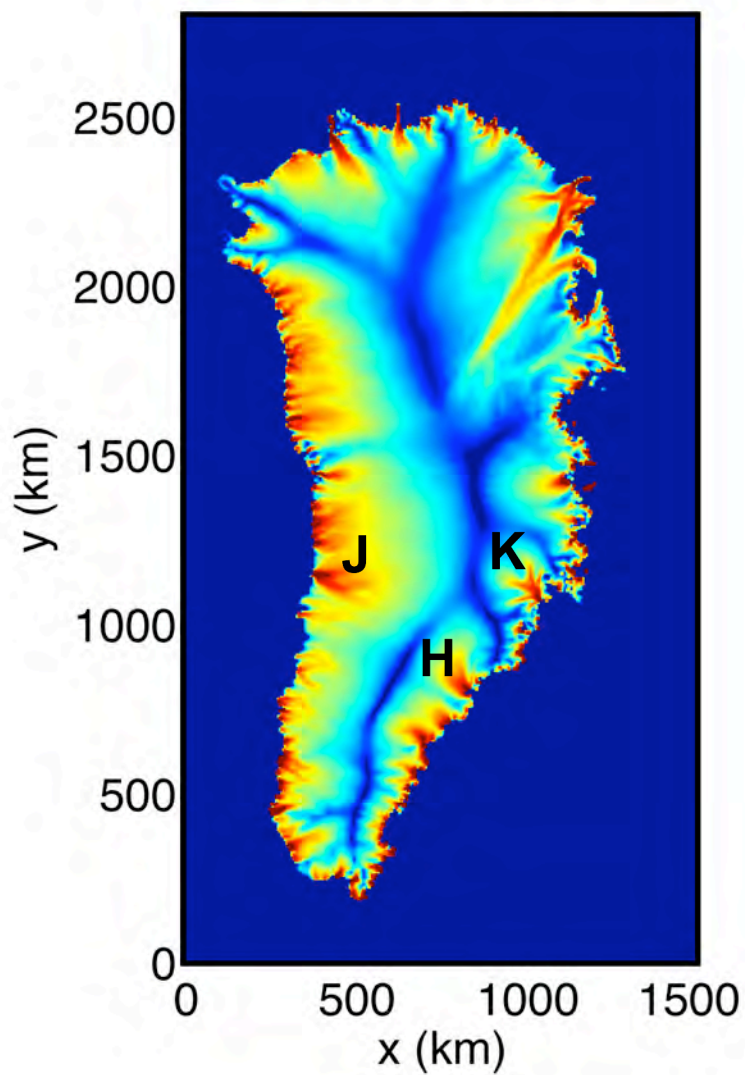
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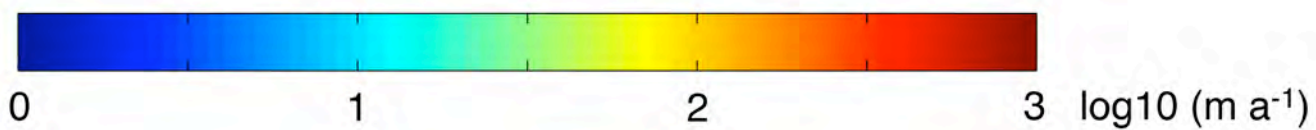
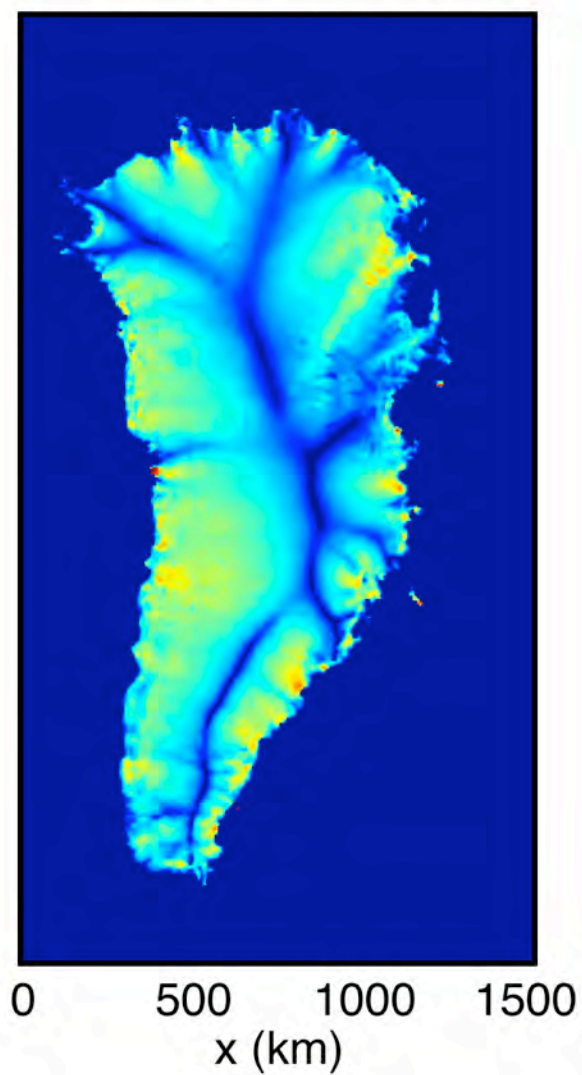
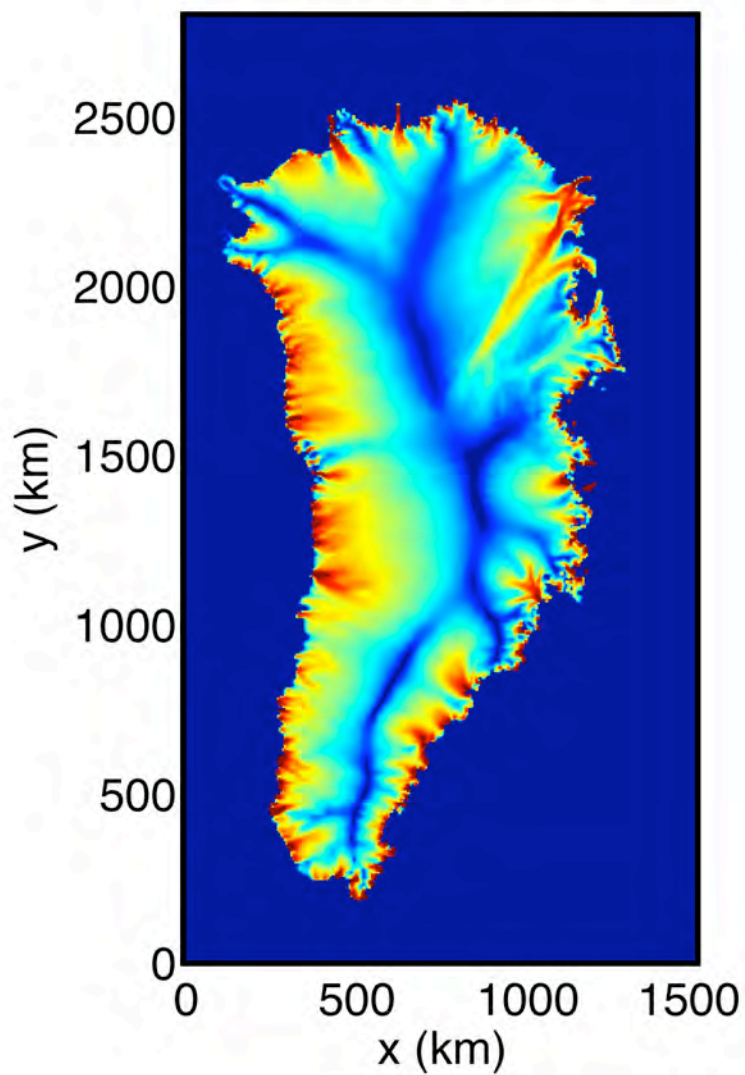


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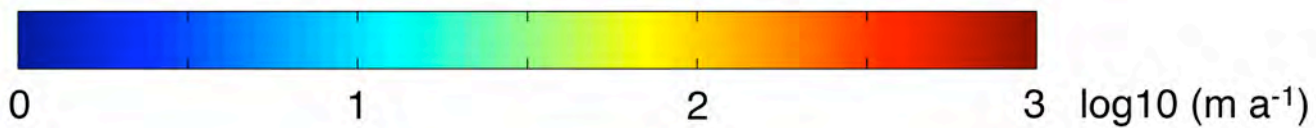
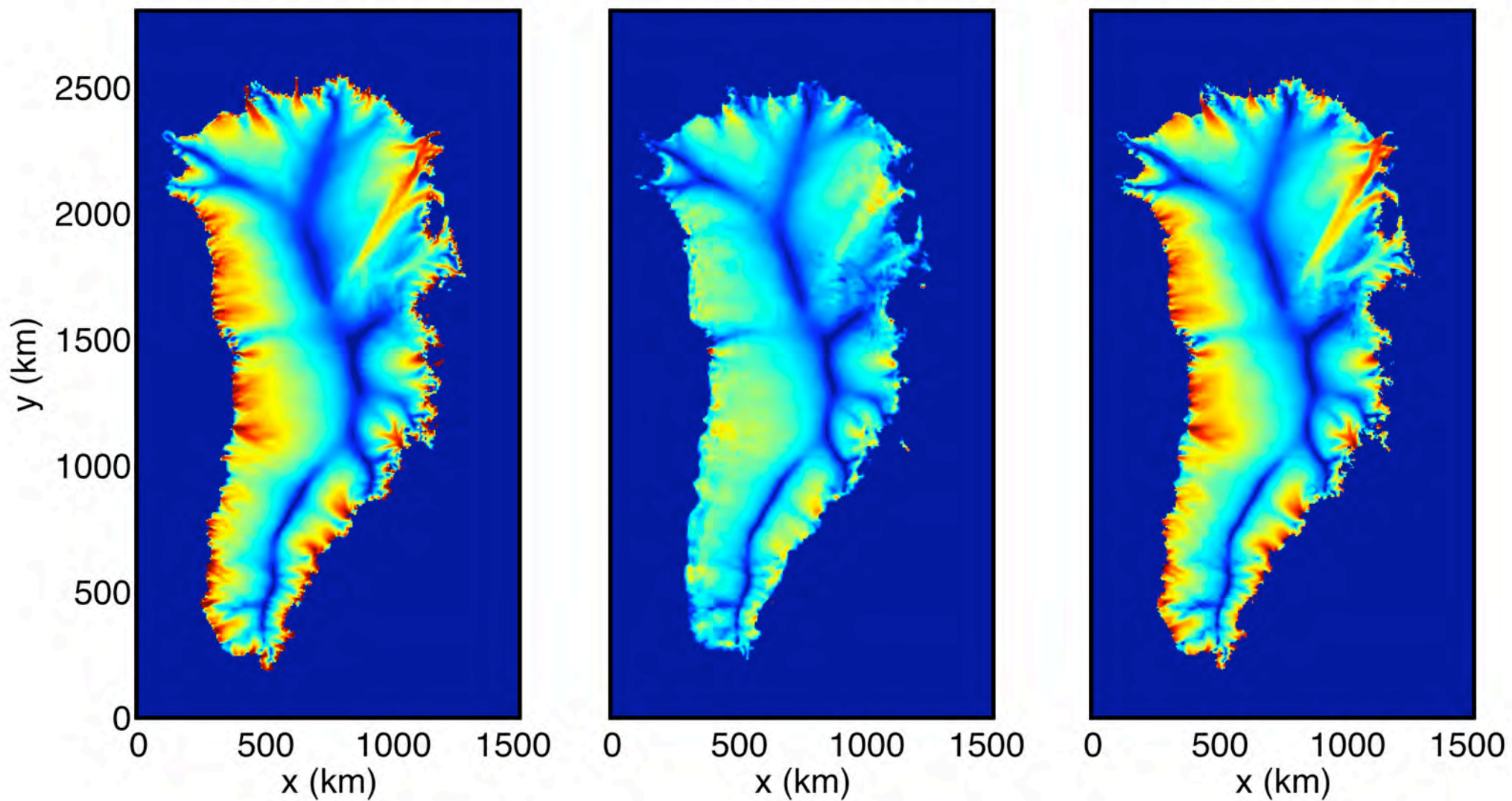
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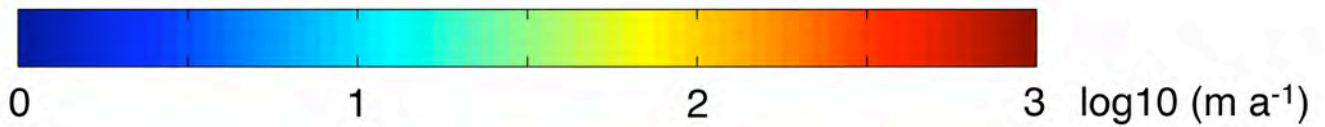
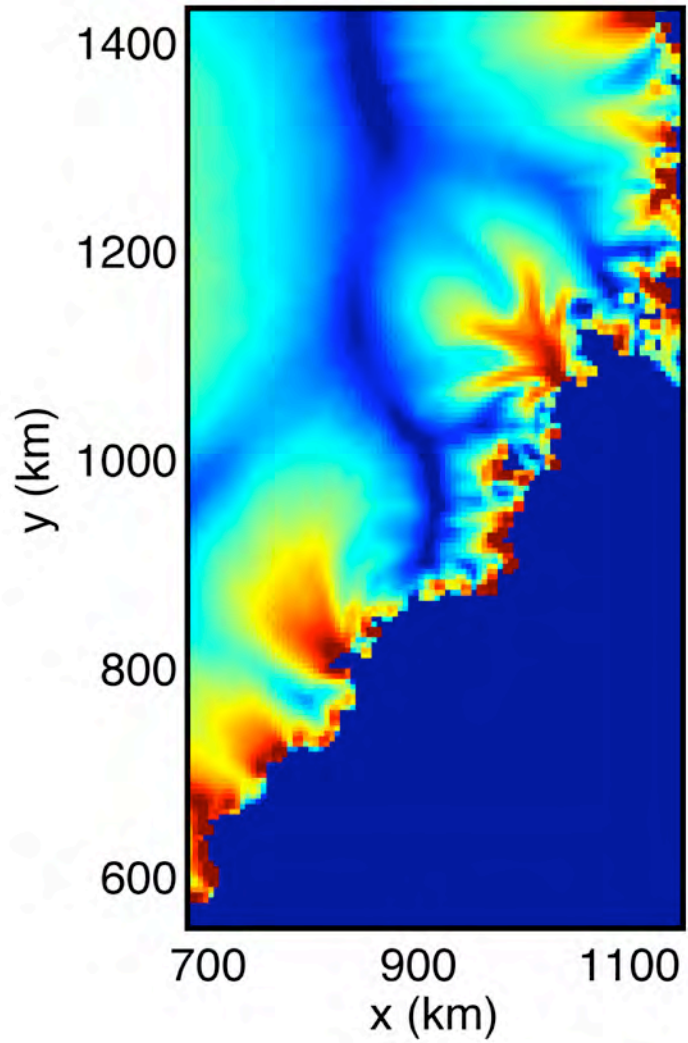
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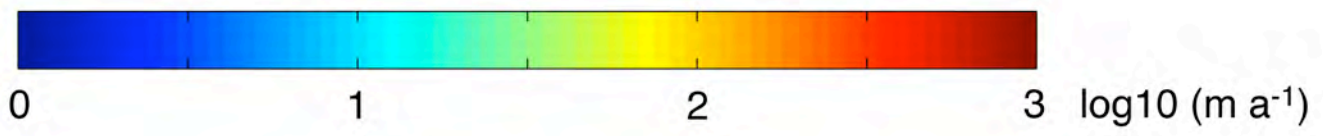
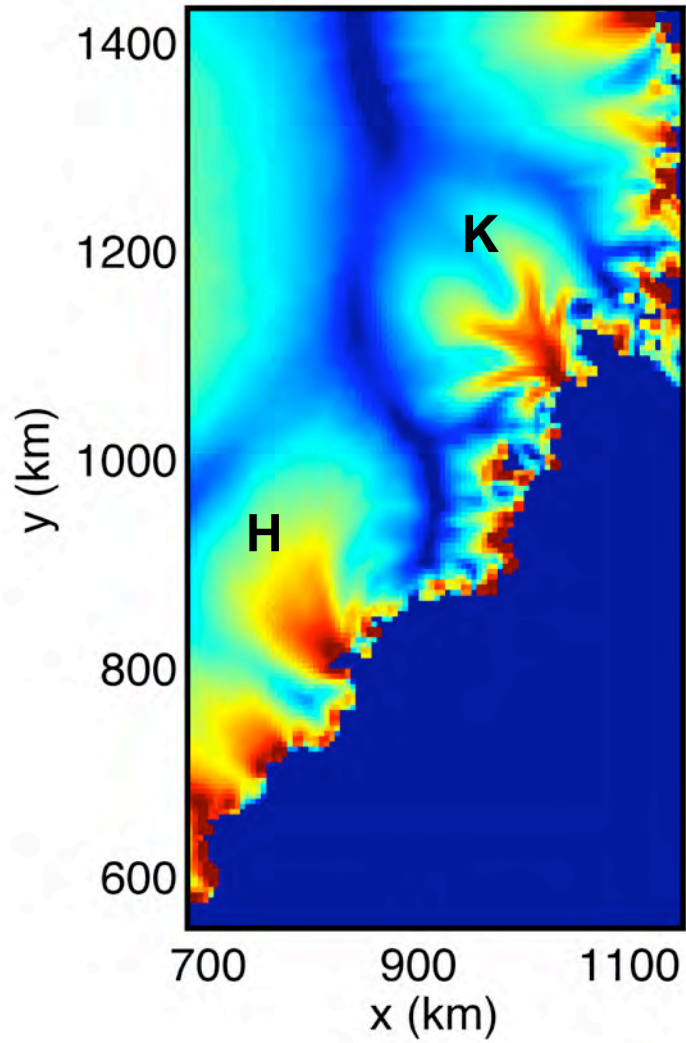
ss sliding velocities



balance velocities

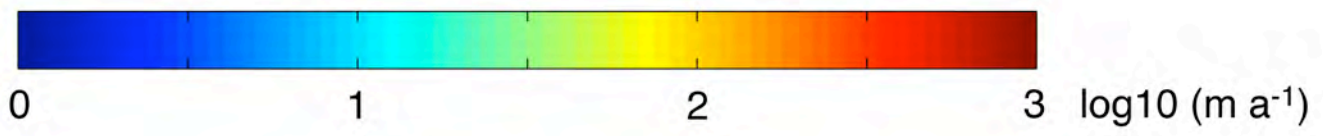
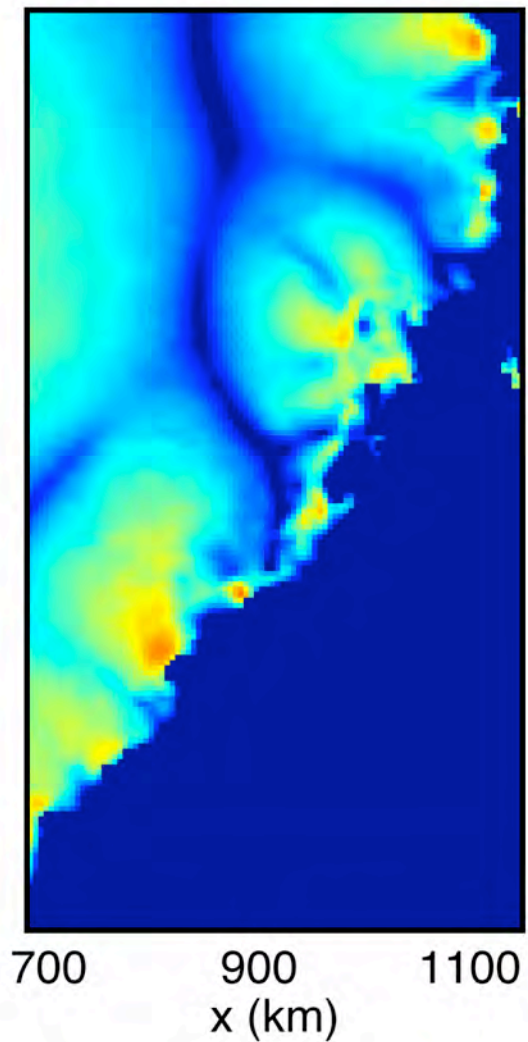
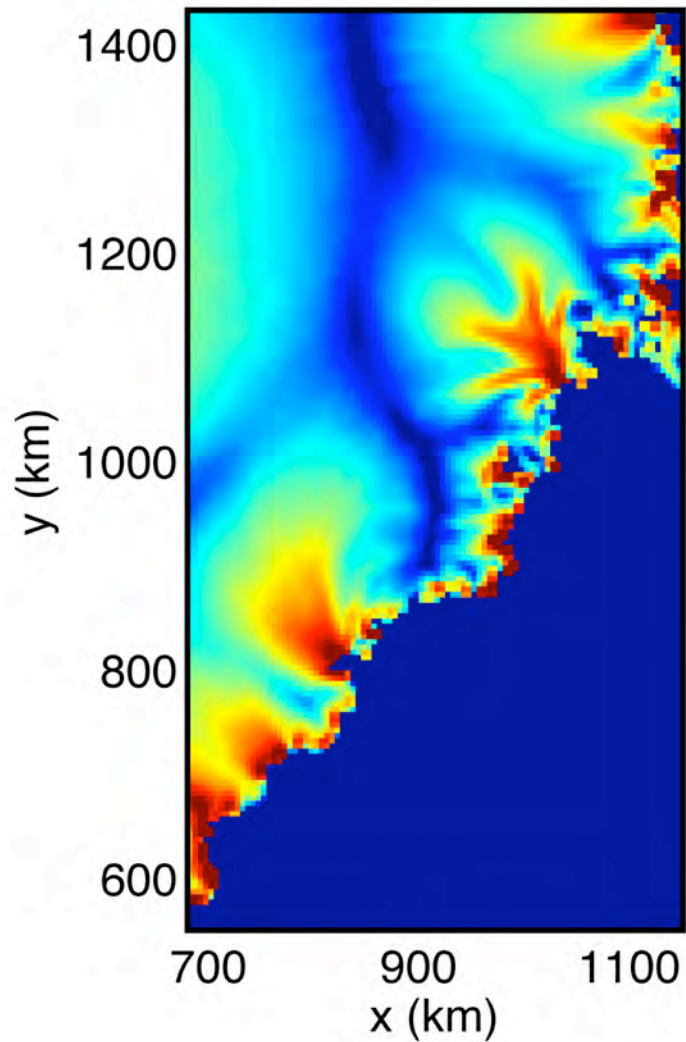


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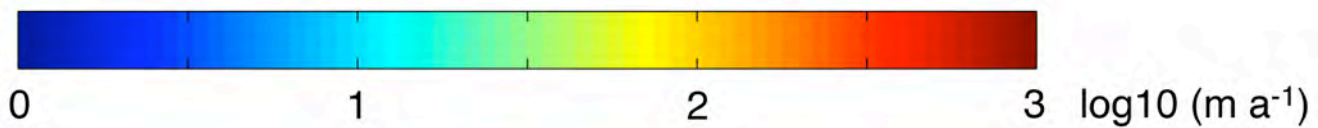
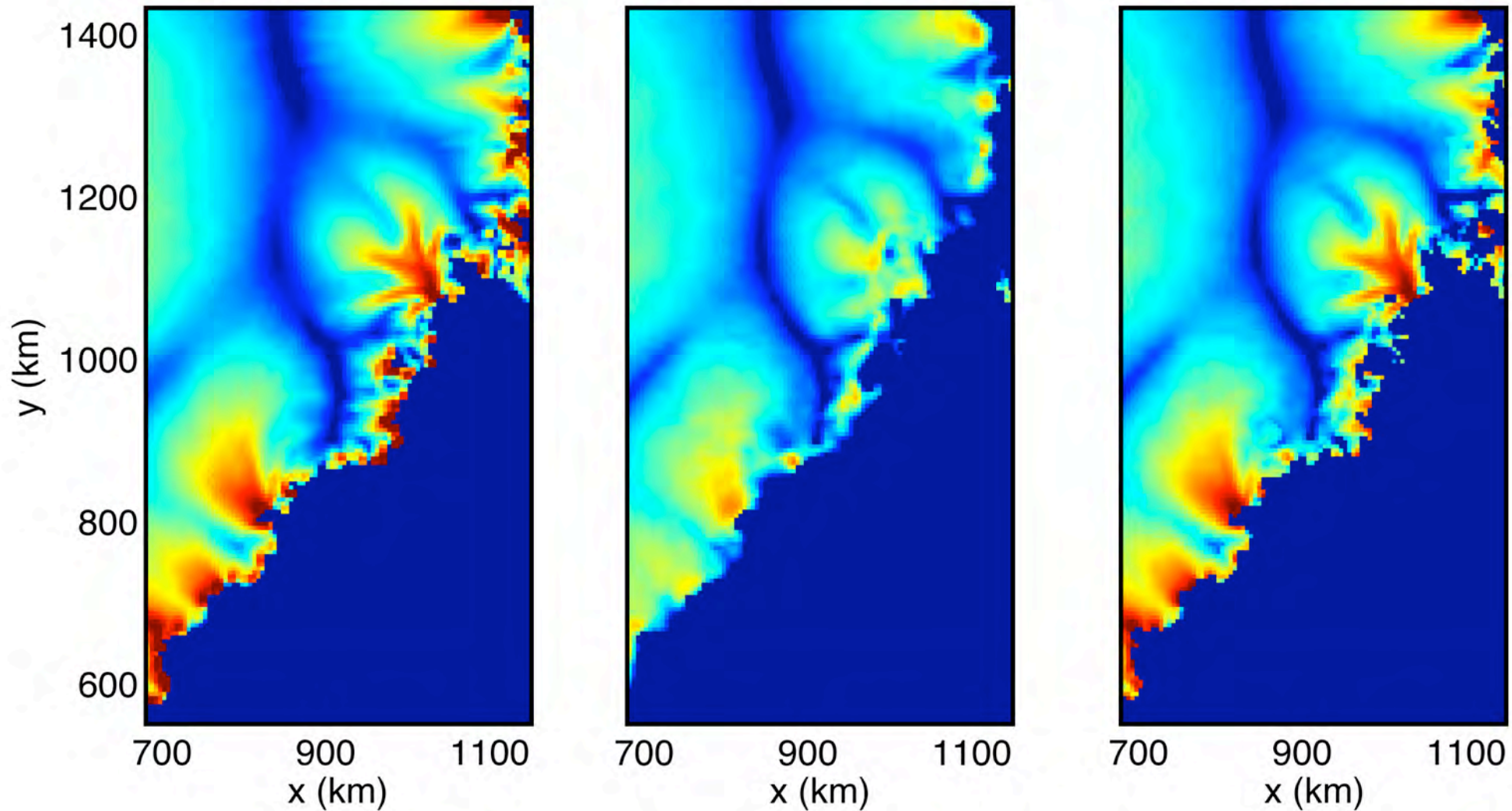
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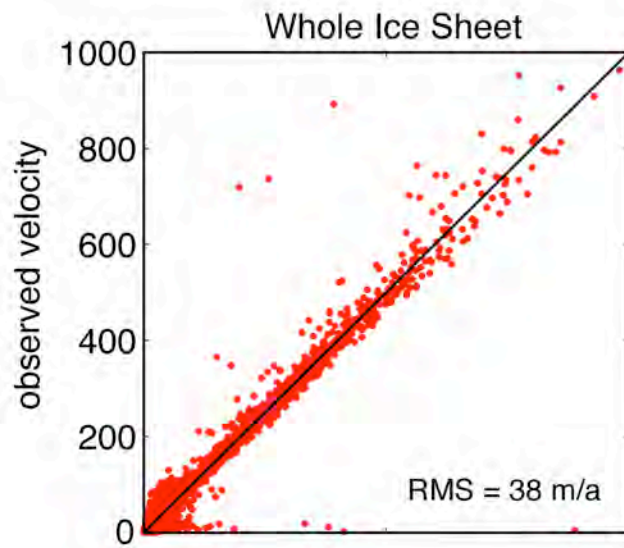
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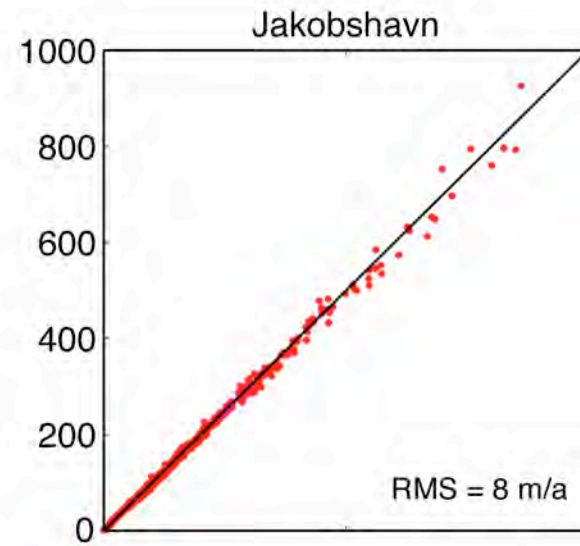
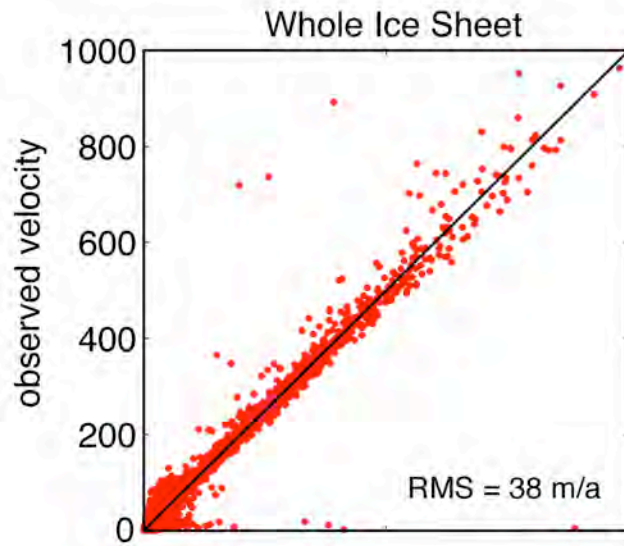


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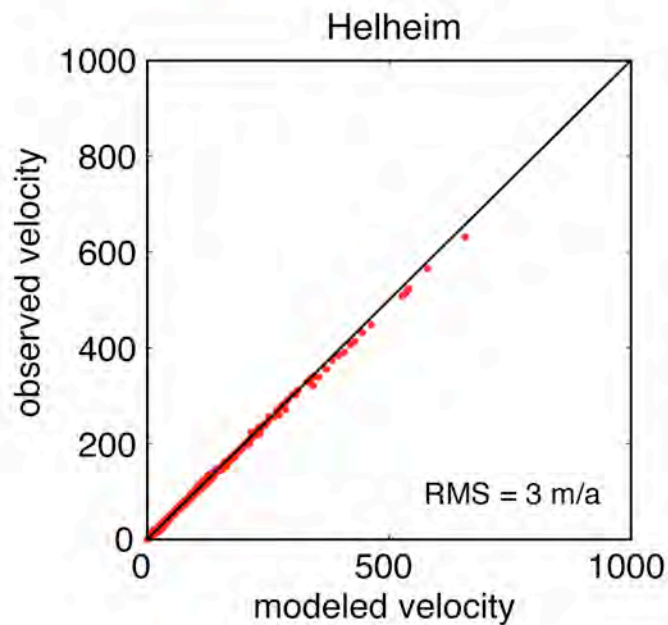
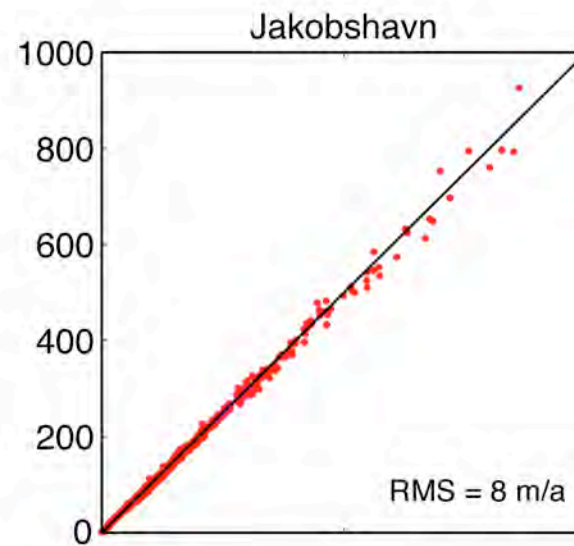
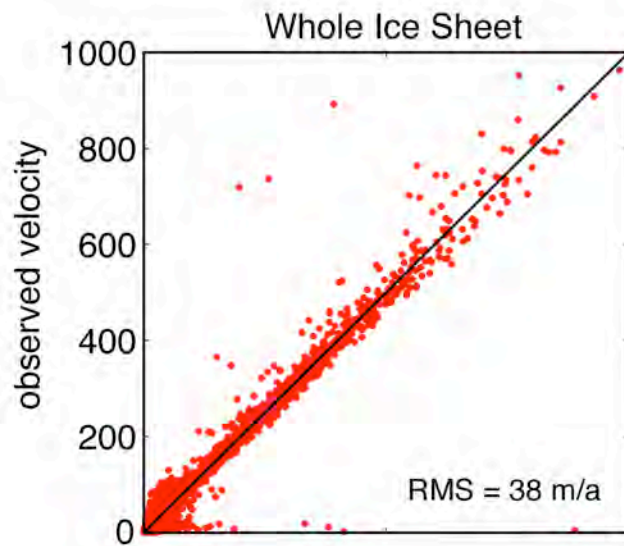
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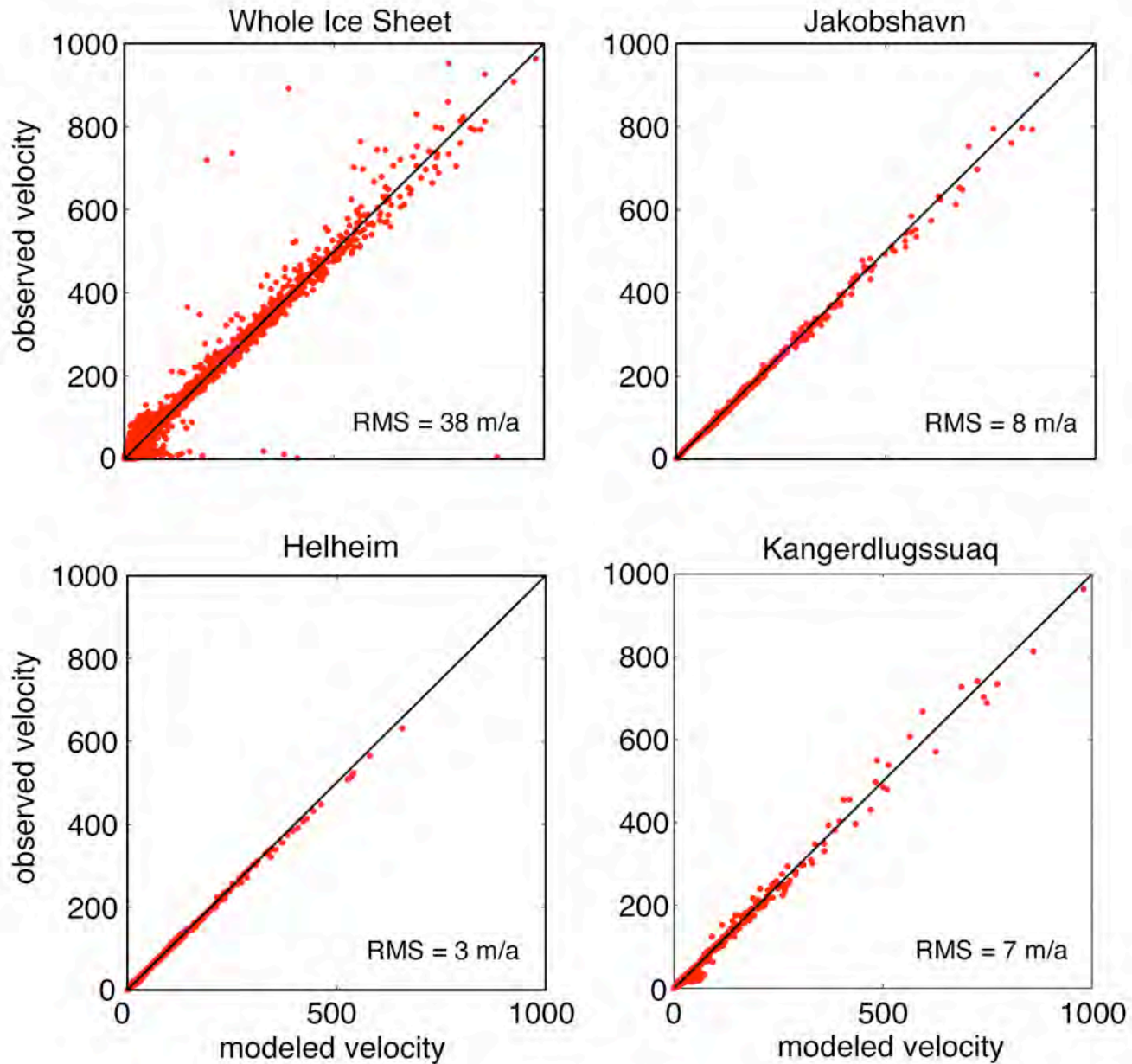
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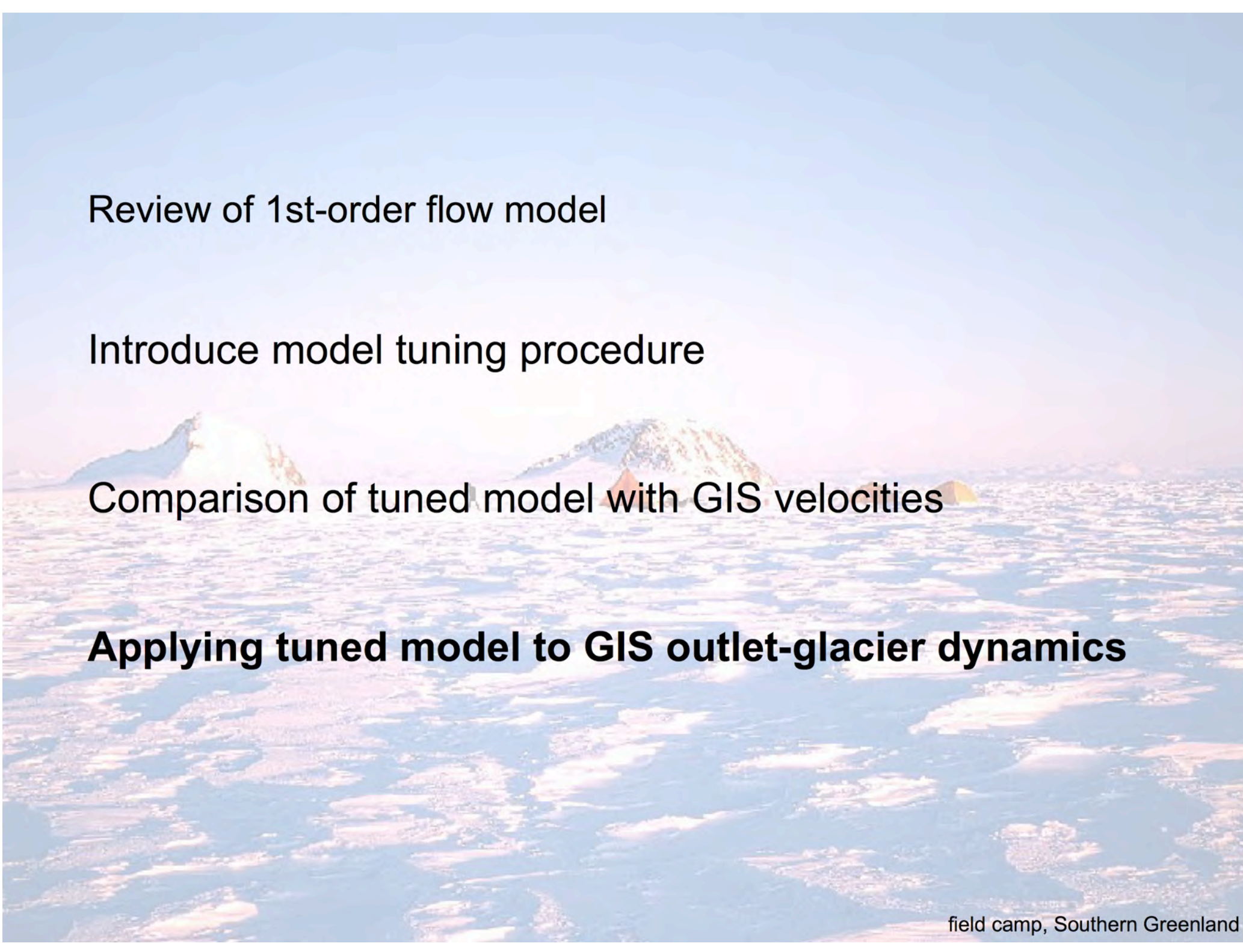


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The total represents a ***minimum*** estimate for SLR over that time period from GIS outlet glacier dynamics.

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Discharge ($\text{km}^3 \text{ a}^{-1}$) from select GIS basins:

<u>Basin</u>	<u>model (% error)</u>	<u>observations¹</u>
Jakobshavn:	22.8 (-3)	23.6
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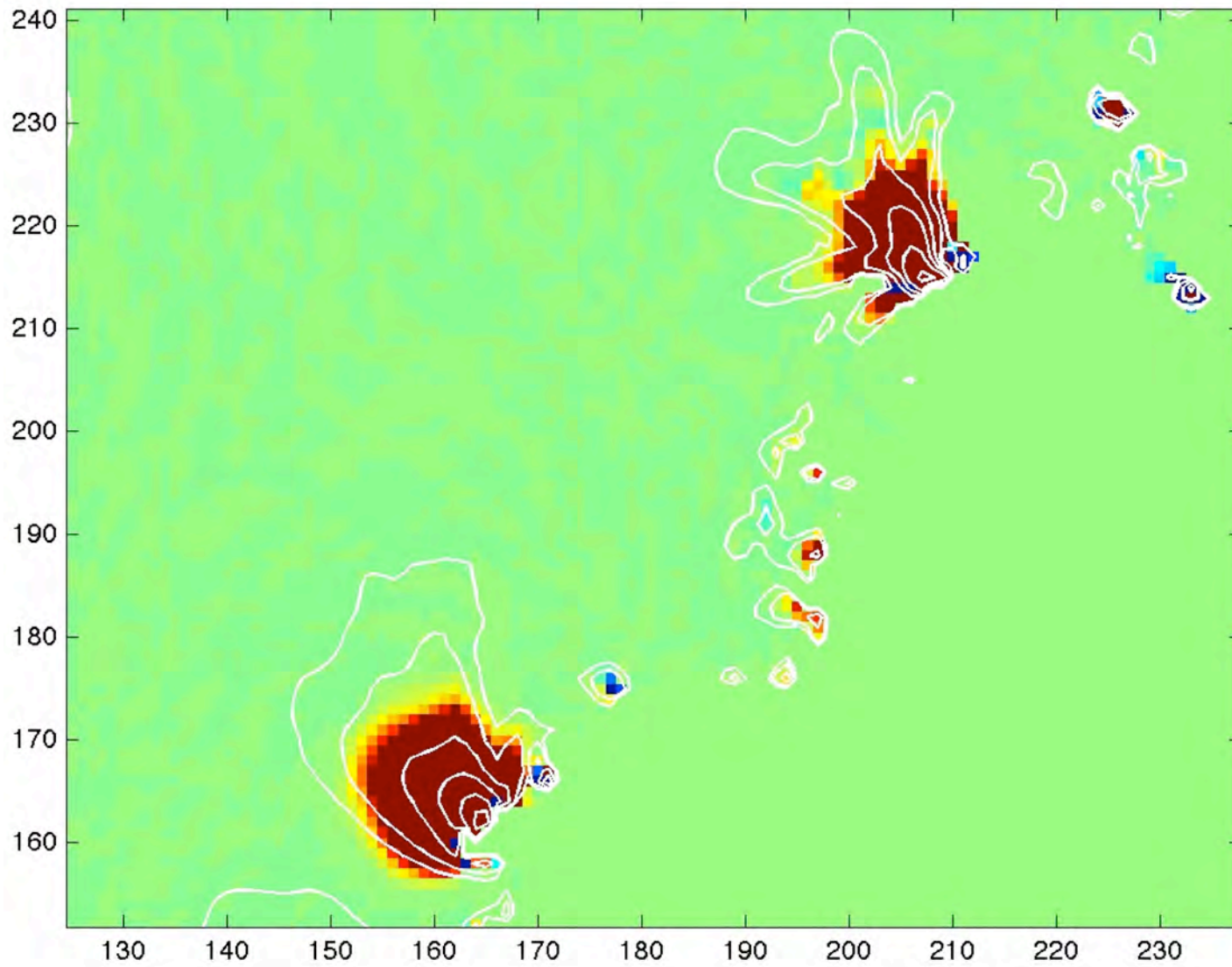
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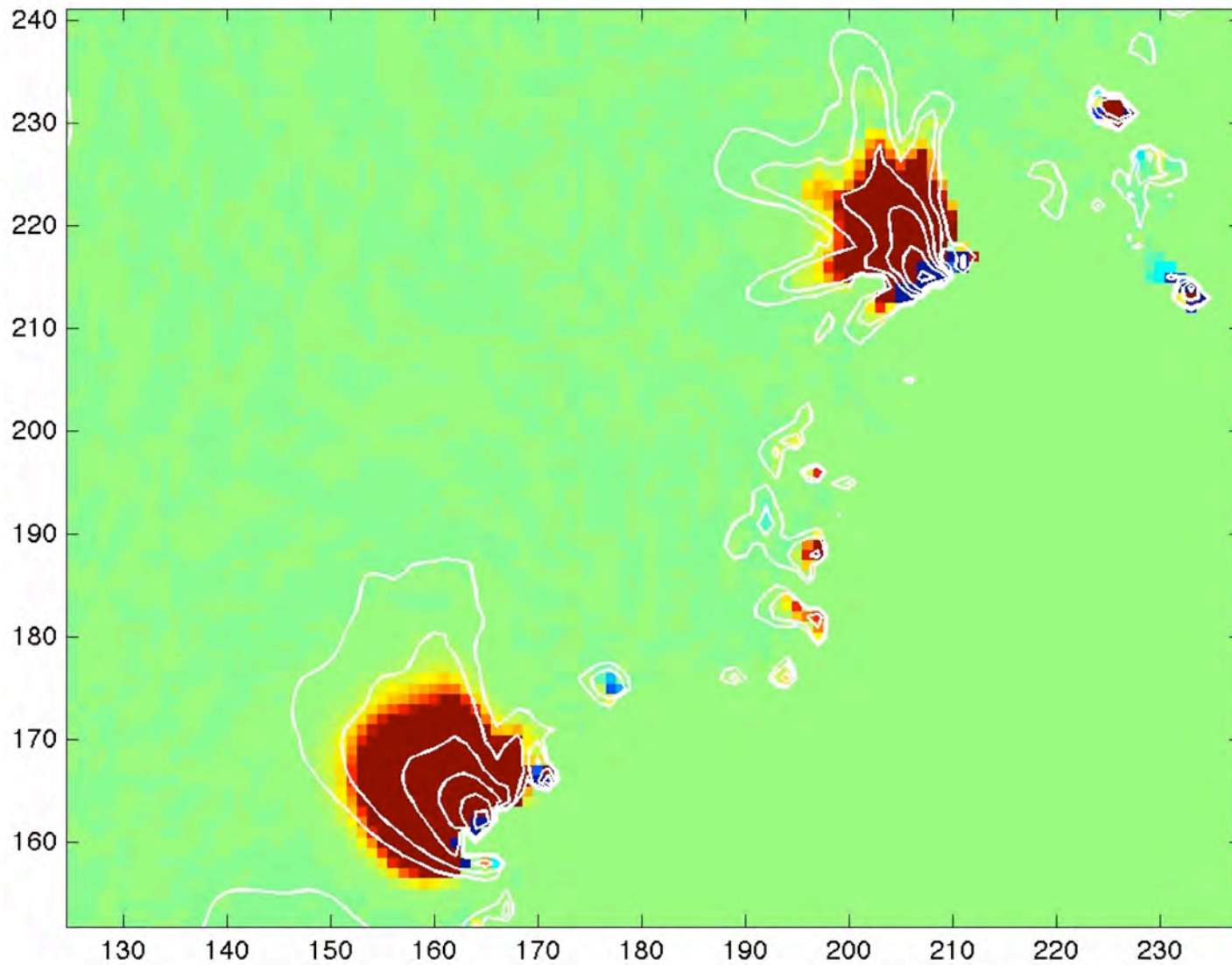
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... Next, time series showing change in speed over Kang. and Hel. Glacier catchments from year 1 to year 9 ...

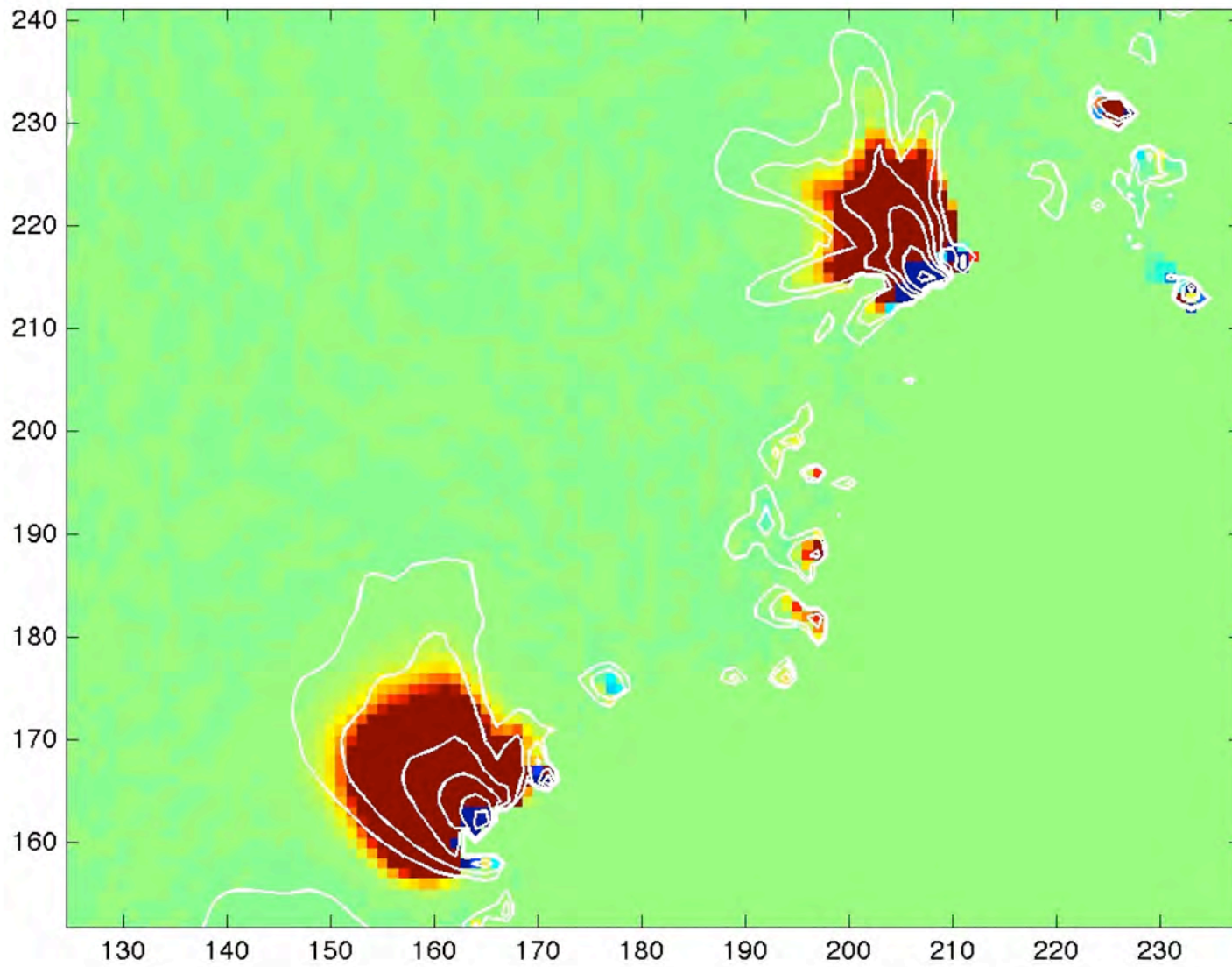
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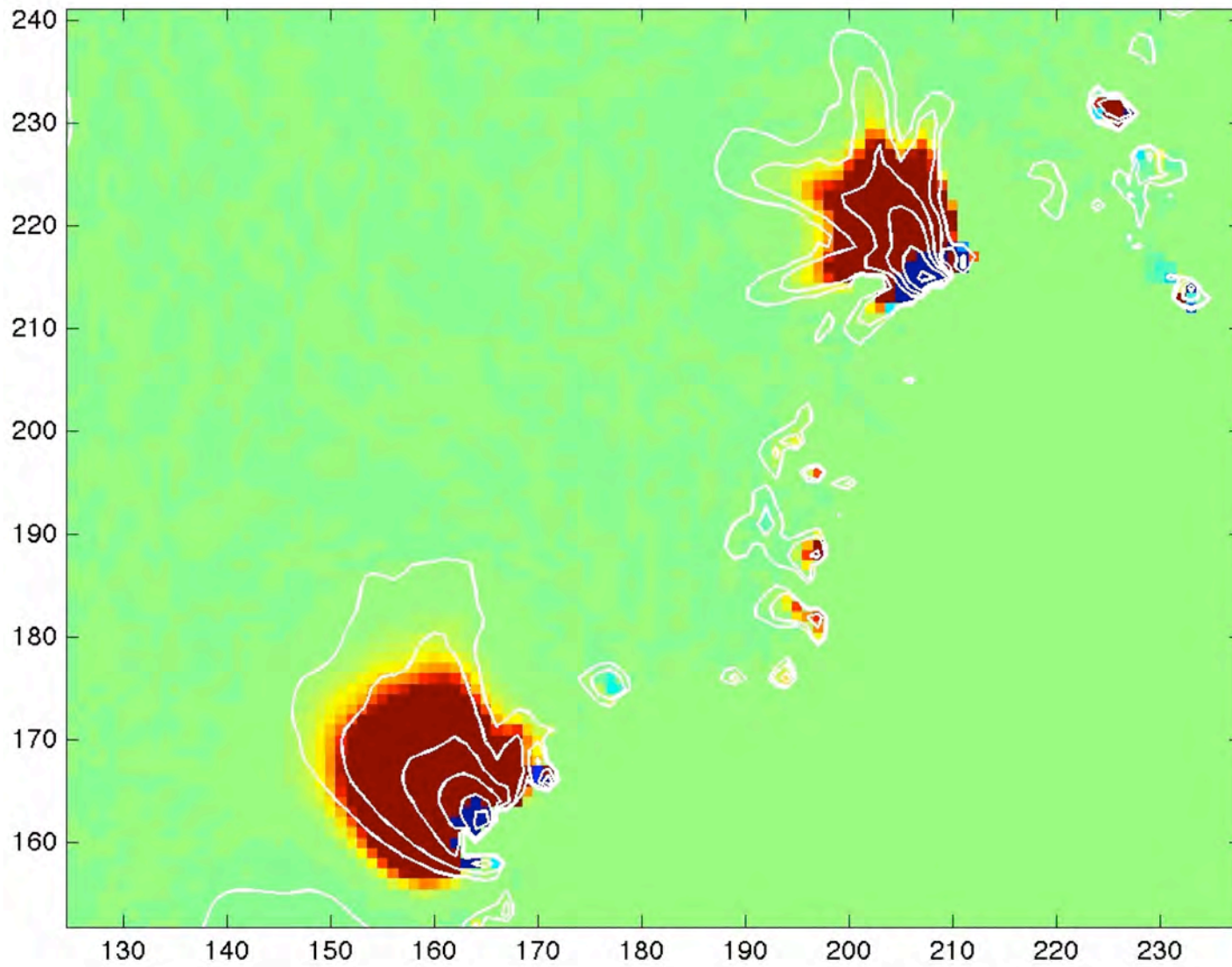
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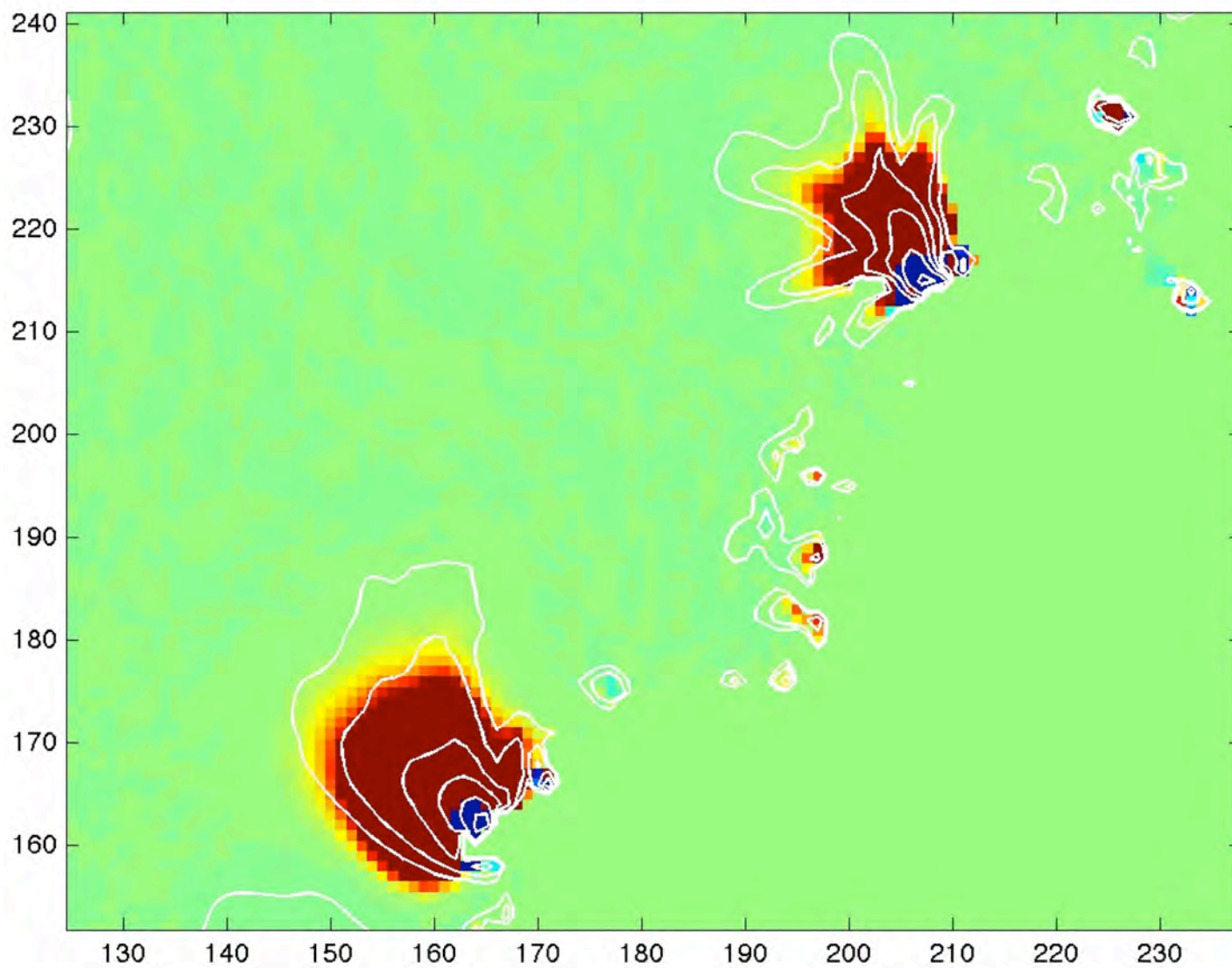
time = 3 yrs



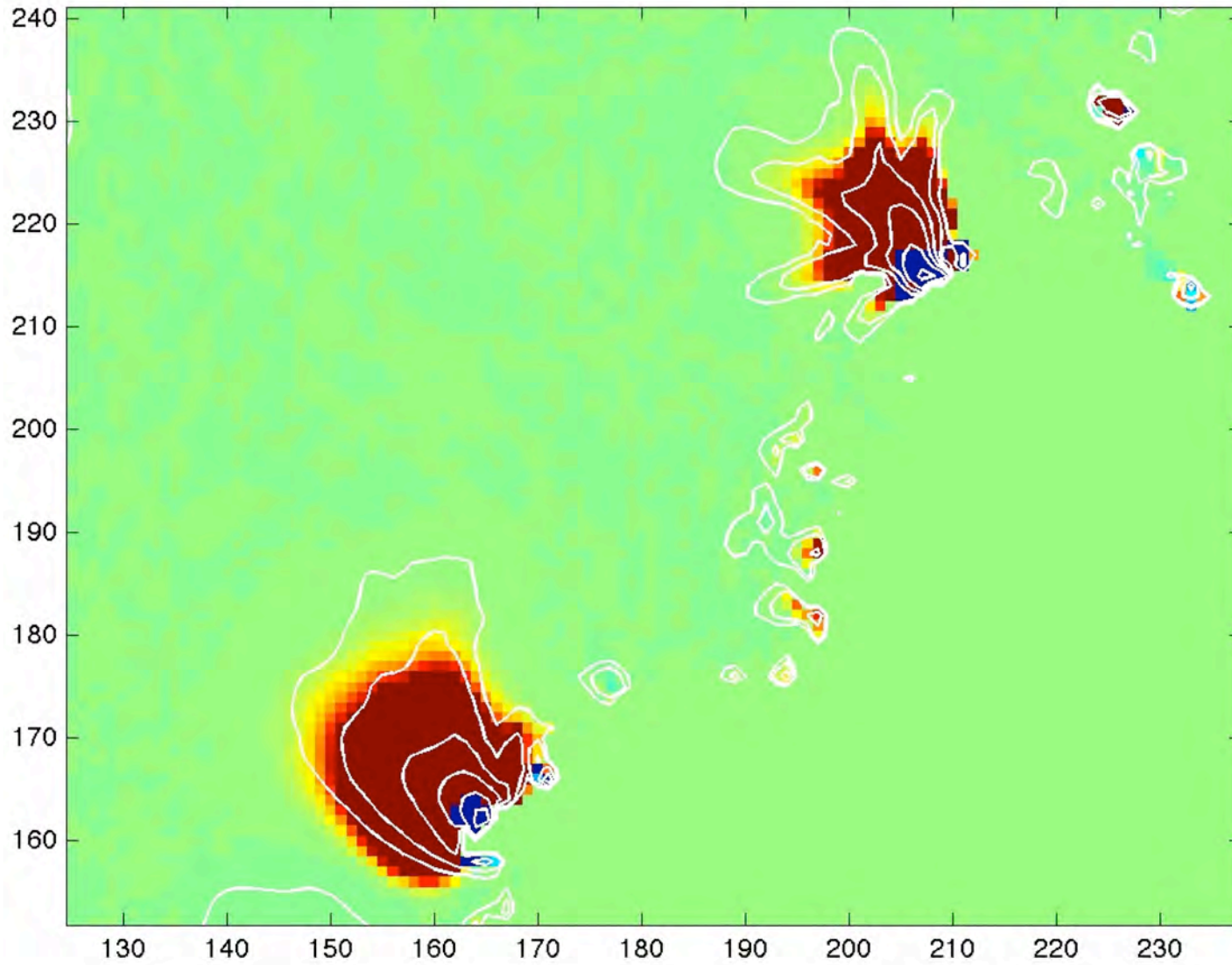
time = 4 yrs



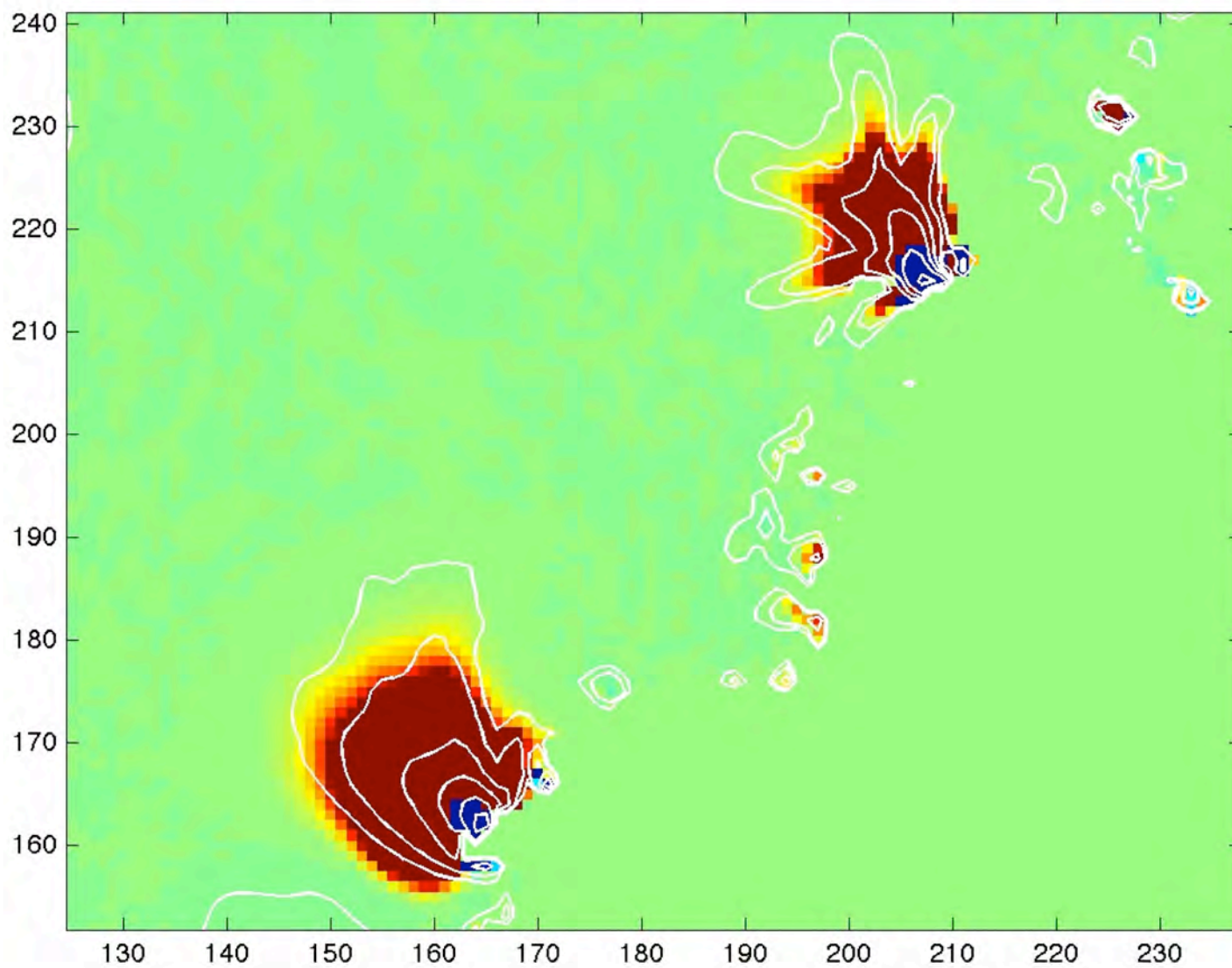
time = 5 yrs



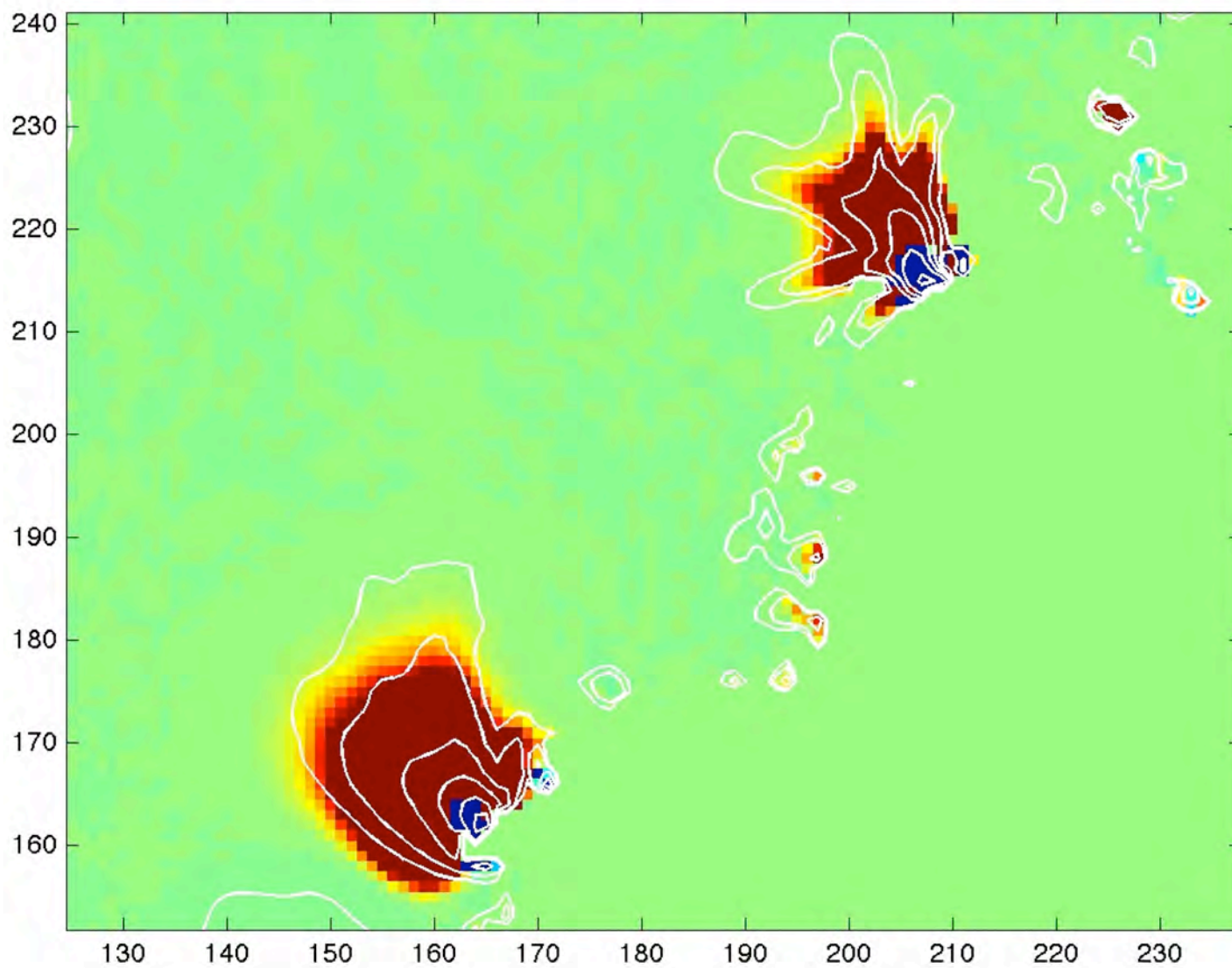
time = 6 yrs



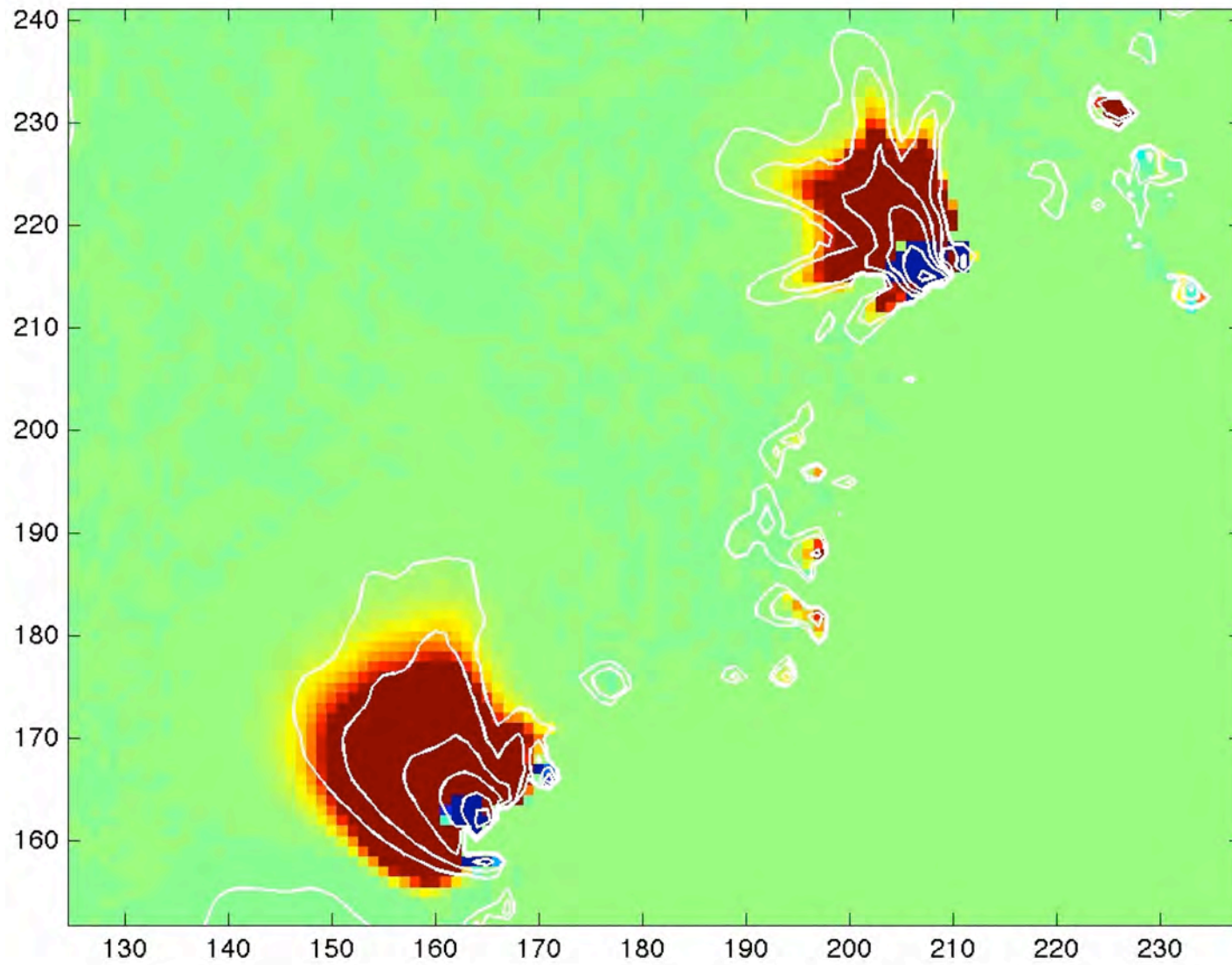
time = 7 yrs

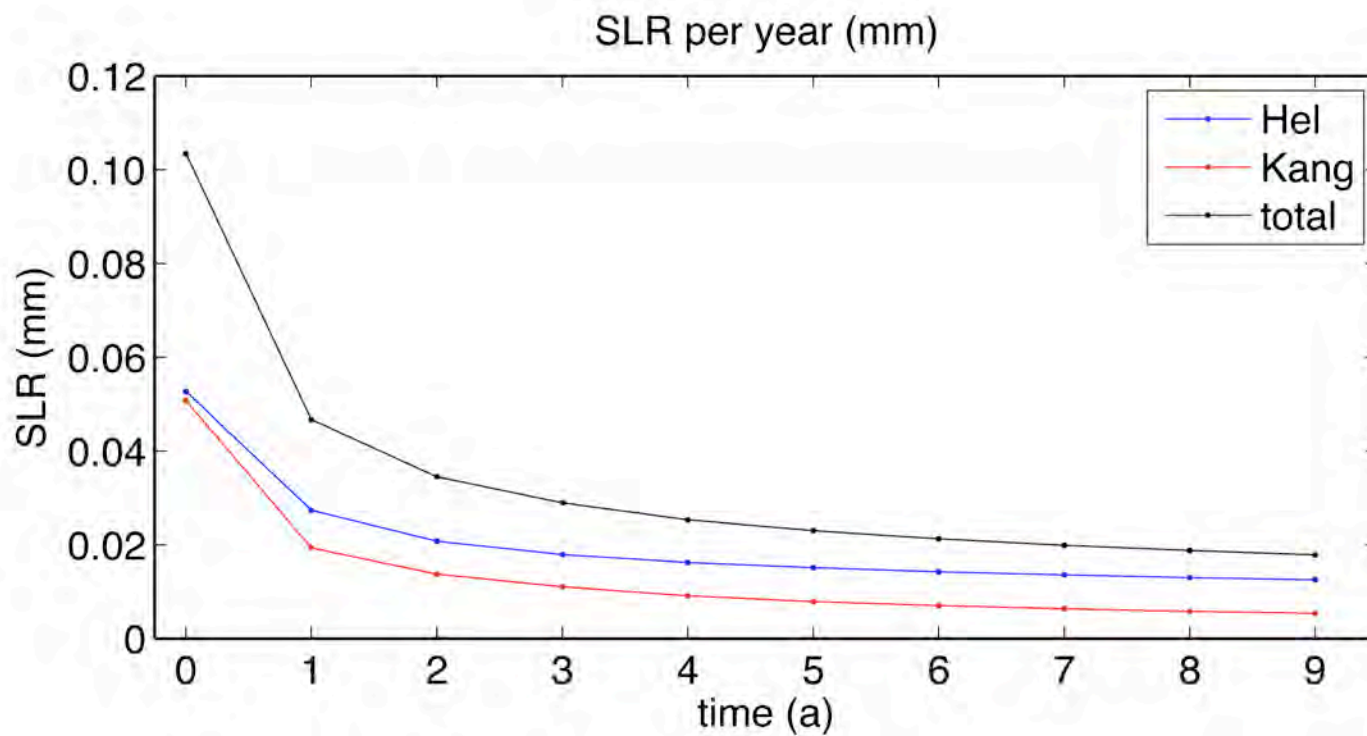
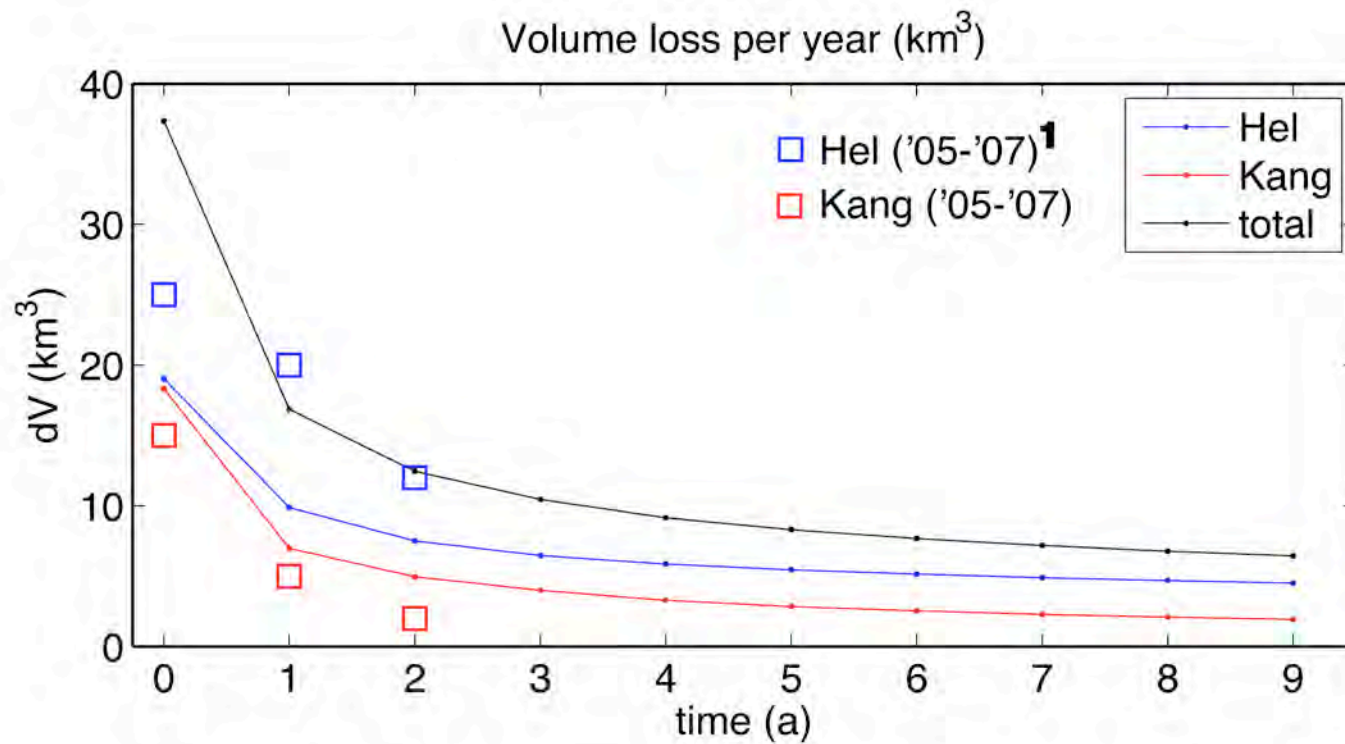


time = 8 yrs

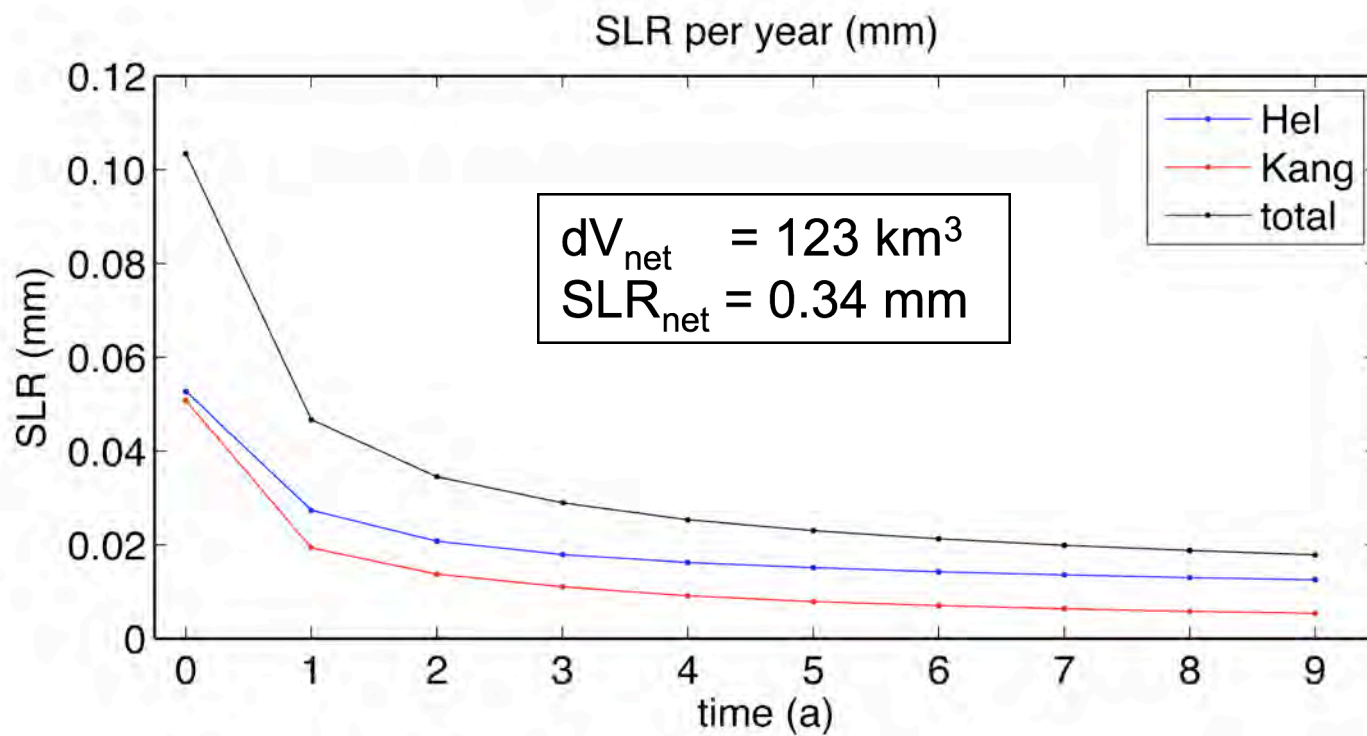
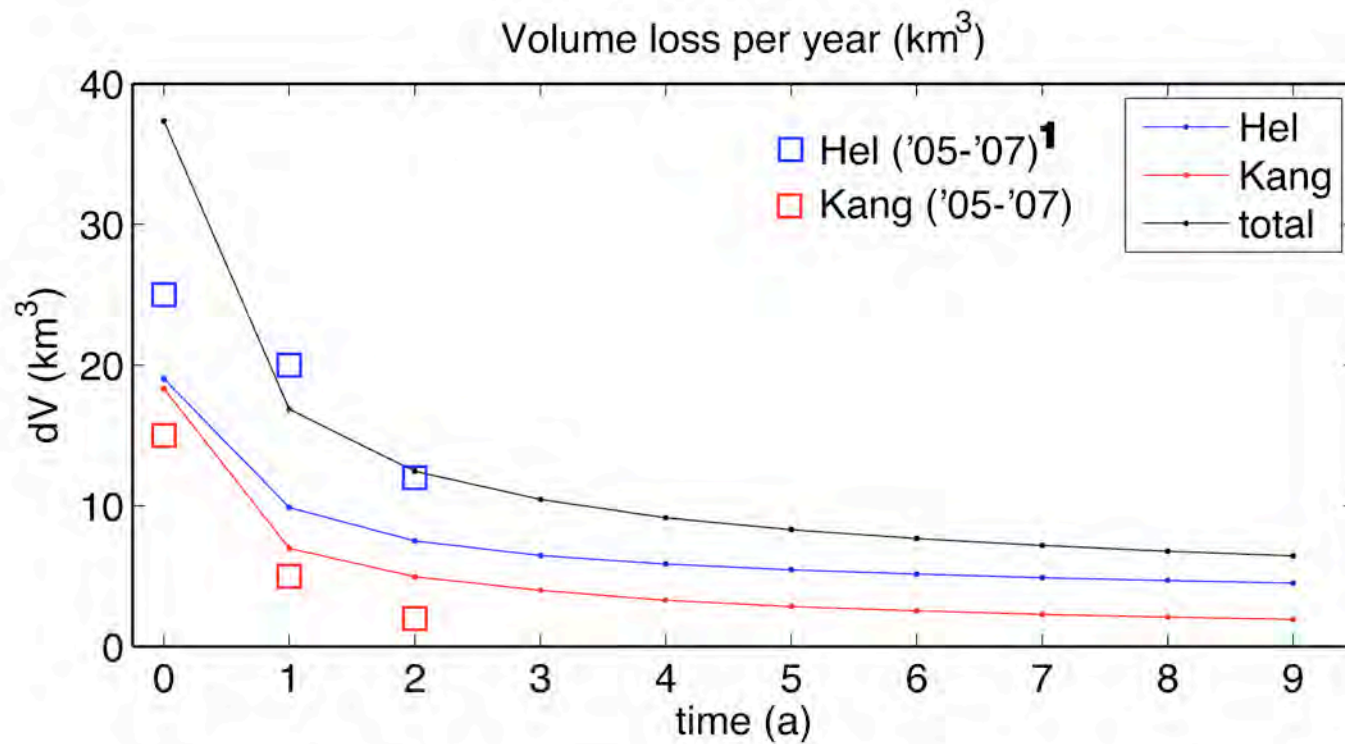


time = 9 yrs



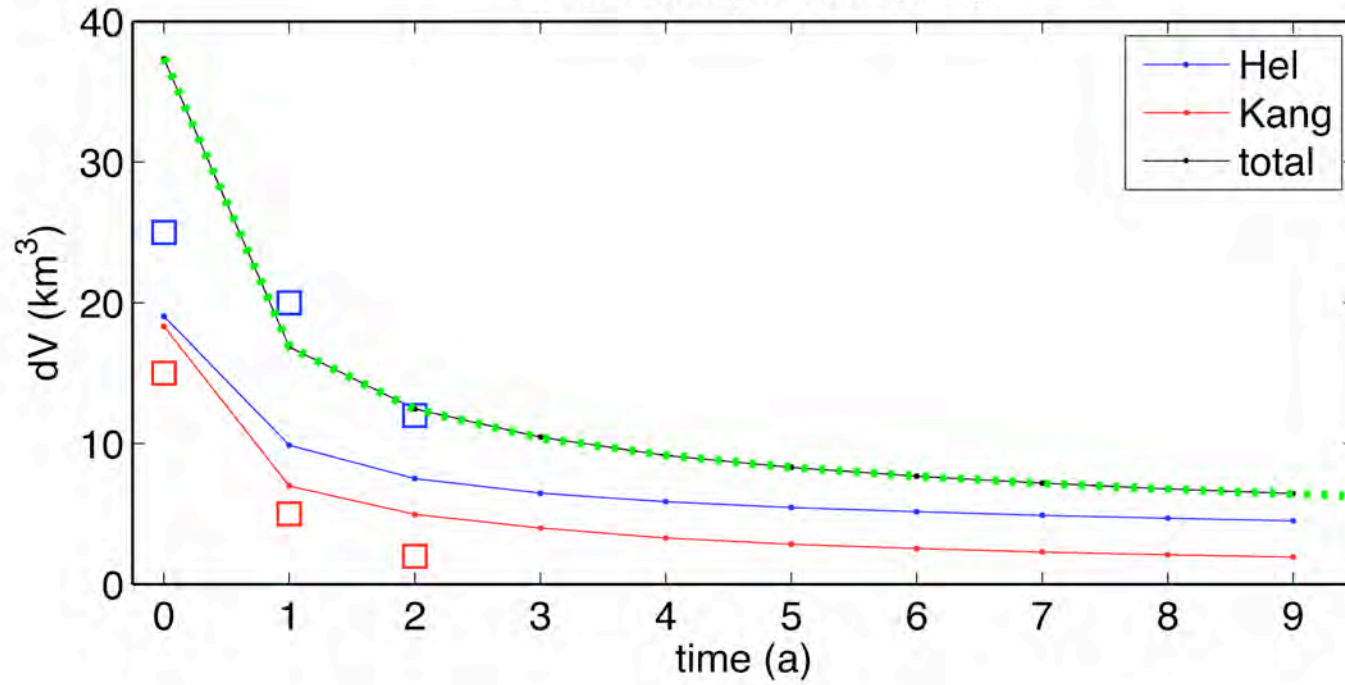


¹Howat et al.
(*Science*, **315**, 2007)

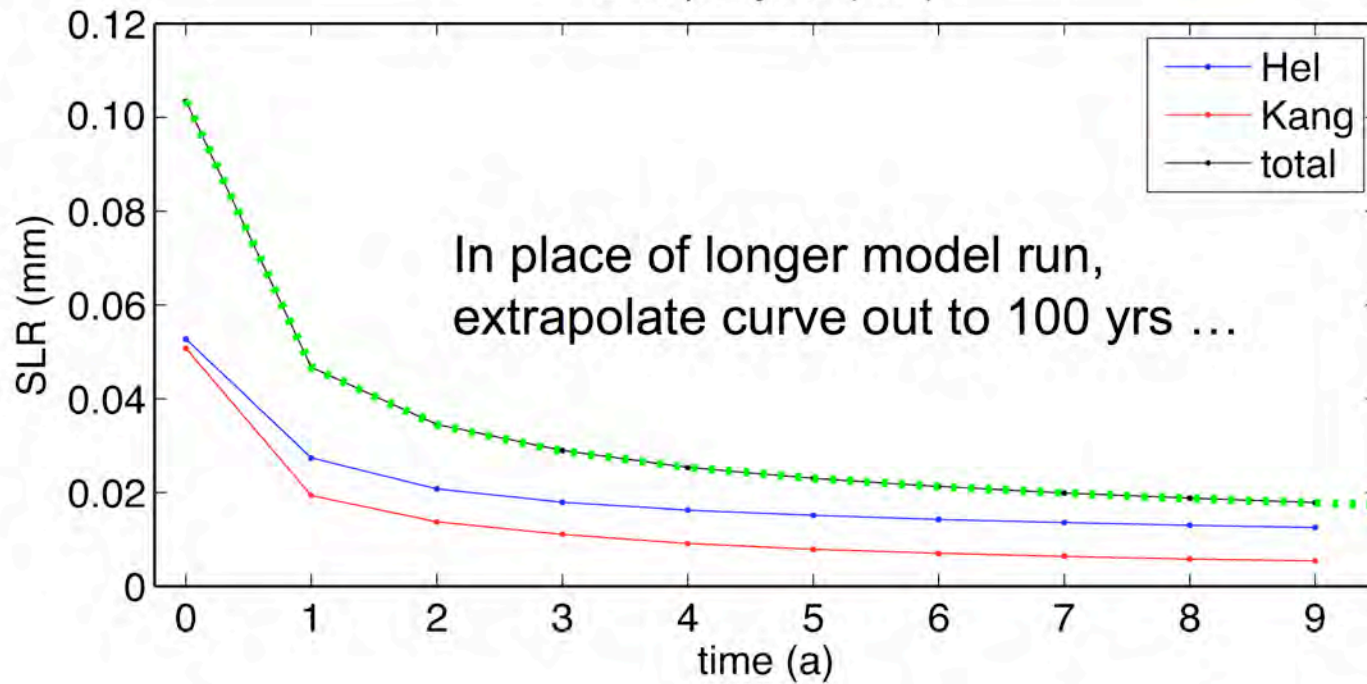


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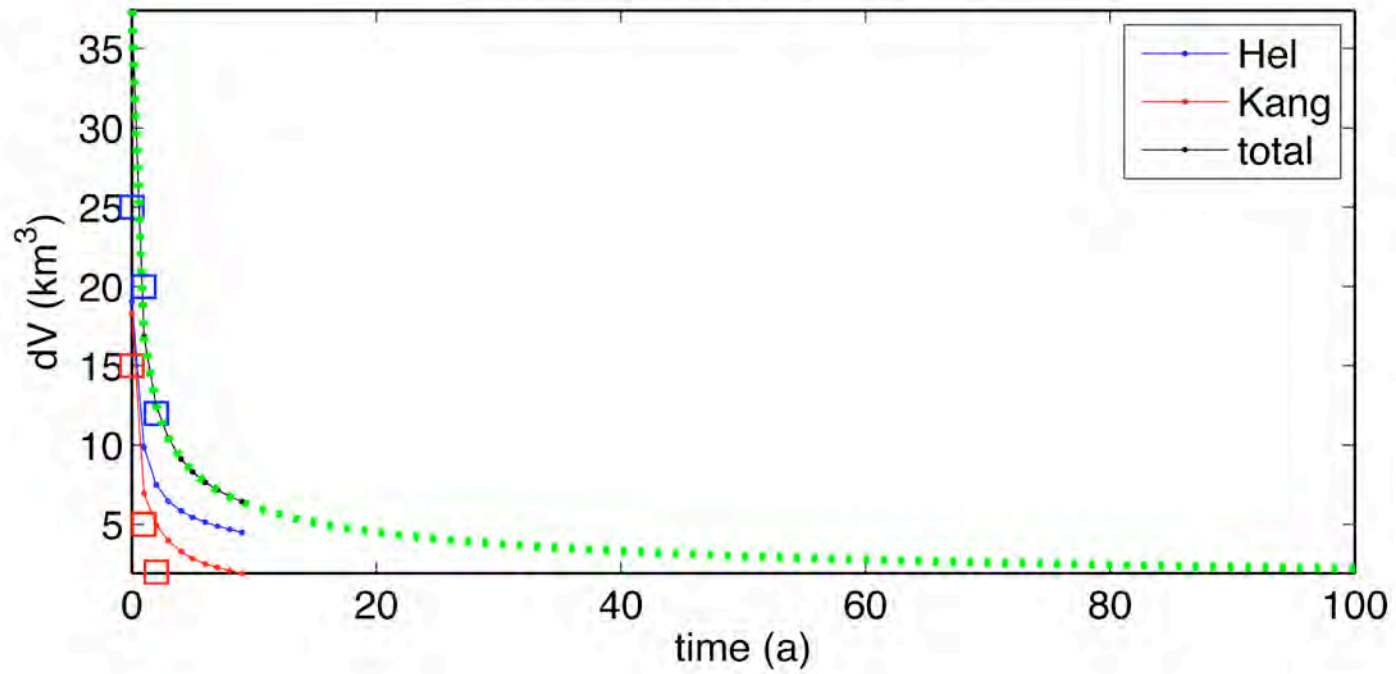
Volume loss per year (km^3)



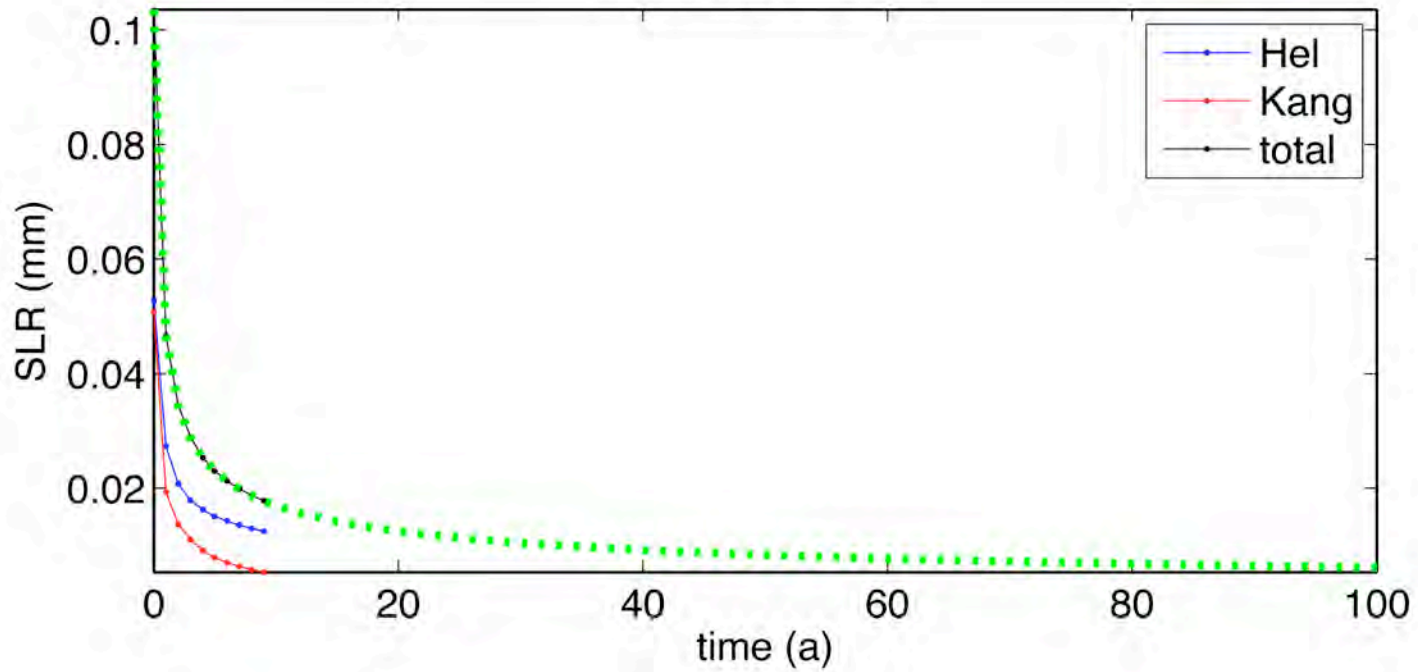
SLR per year (mm)



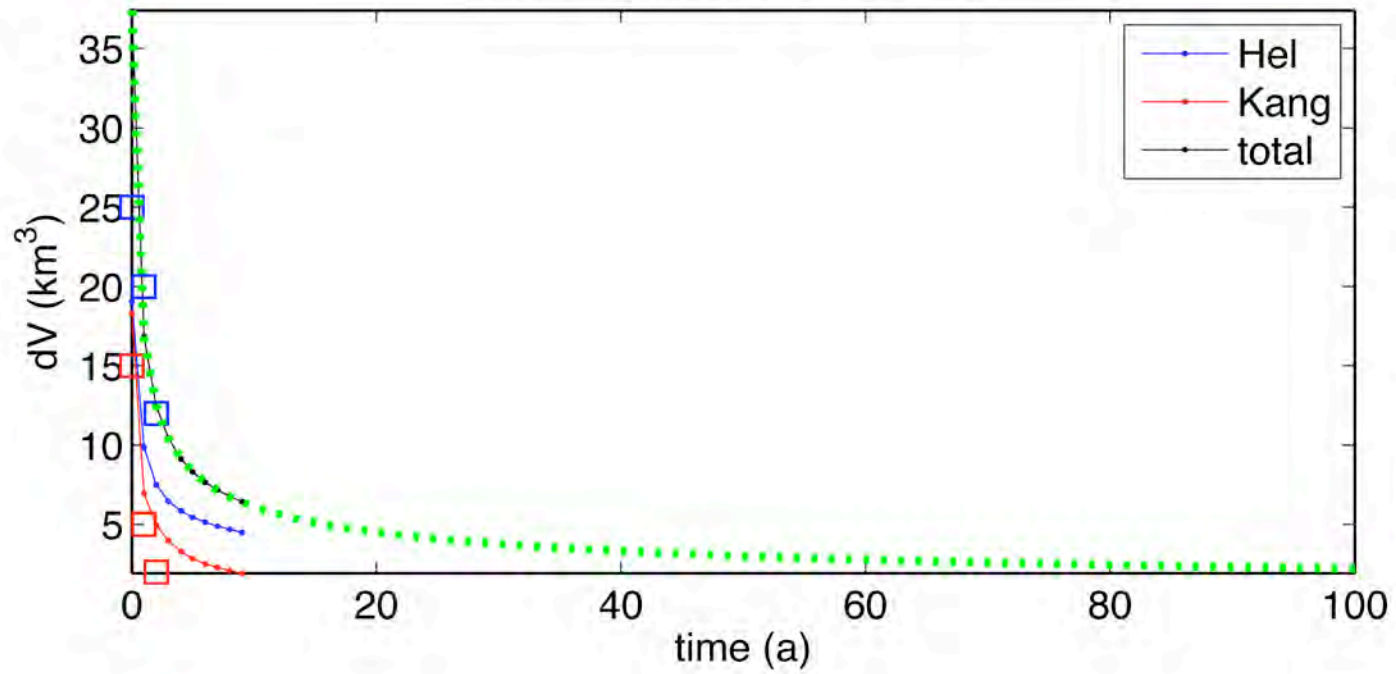
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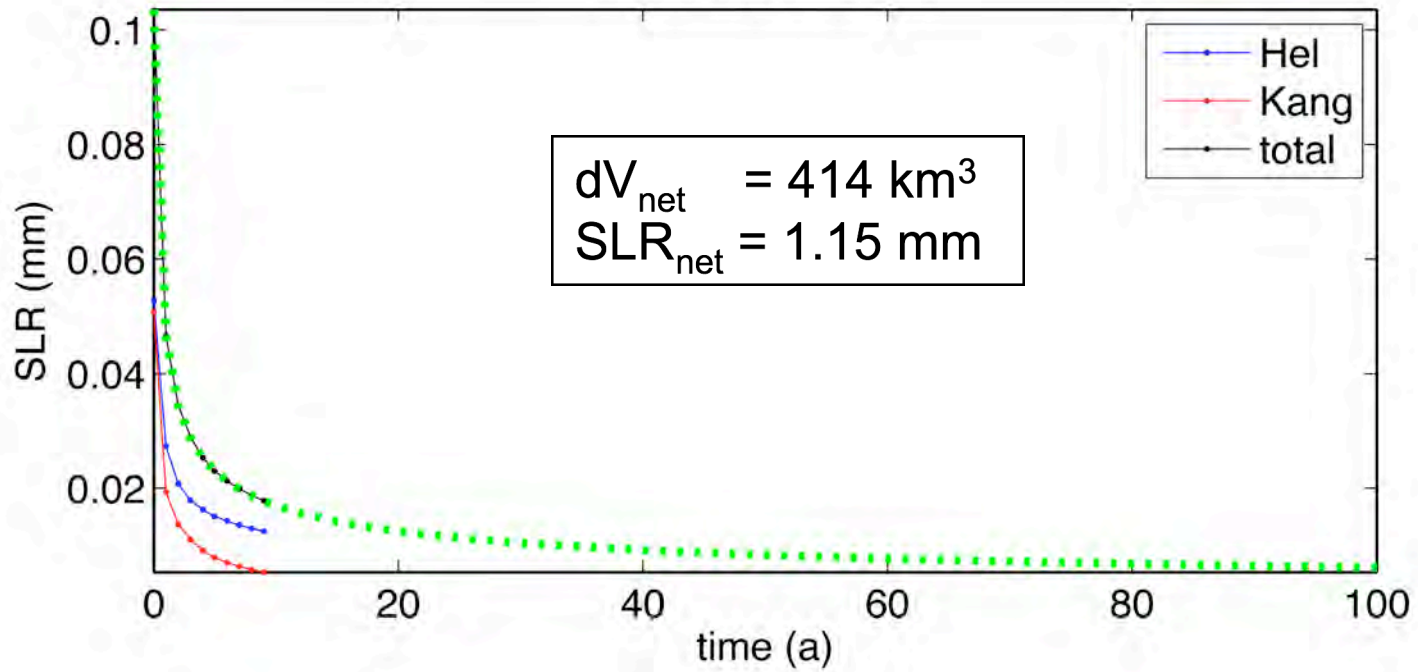
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Jakobshavn, in a continual state of retreat since the late ‘90’s, would raise these numbers significantly (as would the contribution of numerous other marine outlets that have retreated and thinned during the last decade).

END



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surface:	free surface
bed:	$u=v=0$
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Calculation:

... hold geometry, T_{surf} , Q_{geo} steady ...

... allow $B(T)$, \mathbf{u} , and η_{eff} to evolve to steady state ...



¹Blatter (*J. Glac.*, **41**, 1995)

Illulisat, Western Greenland